

# O<sub>2</sub>scl\_eos - Equation of State Sub-Library for O<sub>2</sub>scl

Version 0.910

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## 1 O2scl Equation of State Sub-Library User's Guide

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### 1.1 Feature Overview

O2scl\_eos is sub-library for working with particles and nuclei designed to work with O2scl\_part and O2scl . It includes

- Several classes for computing the hadronic and quark equations of state of matter near and above the nuclear saturation density
  - Classes to solve the TOV equations of neutron star structure
  - One class to perform a rudimentary computation of the structure of closed-shell nuclei in the Hartree approximation.
- 

### 1.2 Quick Reference to User's Guide

- [Equations of State of Hadronic Matter](#)
- [Equations of State of Quark Matter](#)
- [Solution of the Tolman-Oppenheimer-Volkov equations](#)
- [Cold Neutron Stars](#)
- [Nuclear structure in the Hartree approximation](#)
- [Equations of State for Core-Collapse Supernovae](#)
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## 2 Equations of State of Hadronic Matter

The hadronic equations of state are all inherited from [hadronic\\_eos](#): [schematic\\_eos](#), [skyrme\\_eos](#), [rmf\\_eos](#), [apr\\_eos](#), and [gen-potential\\_eos](#).

[hadronic\\_eos](#) includes several methods that can be used to calculate the saturation properties of nuclear matter. These methods are sometimes overloaded in descendants when exact formulas are available.

There is also a set of classes to modify the quartic term of the symmetry energy: [rmf4\\_eos](#), [apr4\\_eos](#), [skyrme4\\_eos](#), and [mdi4\\_eos](#) all based on [sym4\\_eos\\_base](#) which can be used in [sym4\\_eos](#).

---

## 3 Equations of State of Quark Matter

The equations of state of quark matter are all inherited from `quark_eos`: `bag_eos` is a simple bag model, `nambu_jl_eos` is the Nambu-Jona-Lasinio model. The CFL and 2SC phases are described by `cfl_njl_eos`, and `cfl6_eos` adds the color-superconducting 't Hooft interaction.

## 4 Solution of the Tolman-Oppenheimer-Volkov equations

The class `tov_solve` provide a solution to the Tolman-Oppenheimer-Volkov (TOV) equations given an equation of state. This is particularly useful for static neutron star structure: given any equation of state one can calculate the mass vs. radius curve and the properties of any star of a given mass. An adaptive integration is employed and calculates the gravitational mass, the baryonic mass (if the baryon density is supplied), and the gravitational potential. The remaining columns is the equation of state are also interpolated into the solution, e.g. if a chemical potential is given, then the radial dependence of the chemical potential for a 1.4 solar mass star can be automatically computed. The equation of state may be specified in arbitrary units so long as an appropriate conversion factor is supplied. An equation of state for low densities (baryon density  $< 0.08 \text{ fm}^{-3}$ ) is provided and can be automatically appended to the user-defined equation of state.

This is still experimental.

## 5 Cold Neutron Stars

The class `cold_nstar` computes the structure of zero-temperature spherically-symmetric neutron stars. It uses `tov_solve` to compute the structure, given a hadronic equation of state (of type `hadronic_eos`). It also computes the adiabatic index, the speed of sound, and determines the possibility of the direct Urca process as a function of density or radius.

This class is still experimental.

## 6 Nuclear structure in the Hartree approximation

See class `rmf_nucleus`.

## 7 Example Source Code

### 7.1 Example list

- [Cold neutron star example](#)

### 7.2 Cold neutron star example

```
/* Example: ex_cold_nstar.cpp
-----
This example solves the TOV equations using class cold_nstar using a
relativistic mean-field EOS from class rmf_eos.
*/

#include <o2scl/cold_nstar.h>
#include <o2scl/rmf_eos.h>
#include <o2scl/hdf_file.h>
#include <o2scl/hdf_io.h>
#include <o2scl/test_mgr.h>
```

```

using namespace std;
using namespace o2scl;
using namespace o2scl_hdf;

// For hc_mev_fm
using namespace o2scl_const;

// A simple function to load the NL3 model
int load_nl3(rmf_eos &rmf);

int main(void) {

    cout.setf(ios::scientific);

    test_mgr t;
    t.set_output_level(1);

    cold_nstar nst;

    // Initialize EOS
    rmf_eos rmf;
    load_nl3(rmf);

    rmf.saturation();
    cout << "Saturation density: " << rmf.n0 << endl;
    cout << "Binding energy: " << rmf.eoa*hc_mev_fm << endl;
    cout << "Effective mass: " << rmf.msom << endl;
    cout << "Symmetry energy: " << rmf.esym*hc_mev_fm << endl;
    cout << "Compressibility: " << rmf.comp*hc_mev_fm << endl;

    // Compute EOS, include muons
    nst.include_muons=true;
    nst.set_eos(rmf);
    nst.calc_eos();
    o2_shared_ptr<table_units>::type te=nst.get_eos_results();

    // Compute mass vs. radius
    nst.calc_nstar();
    o2_shared_ptr<table_units>::type tr=nst.get_tov_results();
    cout << "Maximum mass: " << tr->max("gm") << endl;
    cout << "Radius of maximum mass star: "
        << tr->get("r",tr->lookup("gm",tr->max("gm"))) << endl;
    cout << "Central baryon density of maximum mass star: ";
    cout << tr->get("nb",tr->lookup("gm",tr->max("gm"))) << endl;

    // Output EOS and TOV results to files
    hdf_file hf;
    hf.open("ex_cold_nstar_eos.o2");
    hdf_output(hf,*te,"eos");
    hf.close();
    hf.open("ex_cold_nstar_tov.o2");
    hdf_output(hf,*tr,"tov");
    hf.close();

    t.report();
    return 0;
}
// End of example

```

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## 9 Other Todos

**Idea for Future** Right now, the equation of state classes depend on the user to input the correct value of `non_interacting` for the particle inputs. This is not very graceful...

## 10 Equations of State for Core-Collapse Supernovae

There are several classes designed to provide a consistent interface to several supernova EOS tables. The abstract base class is `gen_sn_eos`. The child classes correspond to different EOS table formats:

- `ls_eos` - The Lattimer-Swesty EOS tables from Jim's webpage ([Lattimer91](#))
- `stos_eos` - The H. Shen et al. EOS tables ([Shen98](#) and [Shen98b](#))
- `sht_eos` - The G. Shen et al. EOS tables ([Shen10a](#), [Shen10b](#), [Shen11](#))
- `hfs1_eos` - The M. Hempel et al. EOS tables ([Hempel10](#) and [Hempel11](#))
- `oo_eos` - The Lattimer-Swesty and H. Shen et al. tables reformatted by O'Connor and Ott ([OConnor10](#))

The `O2scl` distribution does not contain the tables themselves, as they are quite large and most are freely available. `O2scl` includes code which parses these tables and puts them in `tensor_grid3` objects for analysis by the user.

All of these classes are experimental.

## 11 Ideas for Future Development

### Class `apr_eos`

There might be room to improve the testing of the finite temperature part a bit.

There is some repetition between `calc_e()` and `calc_temp_e()` that possibly could be removed.

### Global `apr_eos::parent_method`

This function is probably unnecessary, as the syntax

### Class `bps_eos`

Can the pressure be made to match more closely?

Convert to a `hadronic_eos` object and offer an associated interface?

### Class `cfl_njl_eos`

This class internally mixes `ovector`, `omatrix`, `gsl_vector` and `gsl_matrix` objects in a confusing and non-optimal way. Fix this.

Allow user to change derivative object? This isn't possible right now because the stepsize parameter of the derivative object is used.

### Class `cold_nstar`

Warn if the EOS becomes pure neutron matter.

### Class `ddc_eos`

Implement the finite temperature EOS properly.

### Class `gen_potential_eos`

Calculate the chemical potentials analytically

### Class `gen_sn_eos`

Create a `table` object, possibly using `tensor_grid::vector_slice`.

Show how `matrix_slice` and `vector_slice` can be used with this object.

Could this be a child of `hadronic_eos_temp` and then directly used in `cold_nstar()`?

Add option to load and store a separate lepton/photon EOS

Add muons and/or pions

### Class `hadronic_eos`

Could write a function to compute the "symmetry free energy" or the "symmetry entropy"

### Global `hadronic_eos::saturation ()`

It would be great to provide numerical uncertainties in the saturation properties.

### Class `ldrop_mass_skin`

Add translational energy?

Remove excluded volume correction and compute nuclear mass relative to the gas rather than relative to the vacuum.

In principle,  $T_c$  should be self-consistently determined from the EOS.

Does this work if the nucleus is "inside-out"?

### Class `nse_eos`

Right now `calc_density()` needs a very good guess. This could be fixed, probably by solving for the  $\log(\mu/T)$  instead of  $\mu$ .

### Page **Other Todos**

Right now, the equation of state classes depend on the user to input the correct value of `non_interacting` for the particle inputs. This is not very graceful...

### Class `rmf_delta_eos`

Finish the finite temperature EOS

**Class `rmf_eos`**

It might be nice to remove explicit reference to the meson masses in functions which only compute nuclear matter since they are unnecessary. This might, however, demand redefining some of the couplings.

- Fix `calc_p()` to be better at guessing
- The number of couplings is getting large, maybe new organization is required.
- Overload `hadronic_eos::fcomp()` with an exact version

**Global `rmf_eos::calc_e` (fermion &ne, fermion &pr, thermo &lth)**

Improve the operation of this function when the proton density is zero.

**Global `rmf_eos::calc_e_fields` (fermion &ne, fermion &pr, thermo &lth, double &sig, double &ome, double &rho)**

Improve the operation of this function when the proton density is zero.

**Global `rmf_eos::calc_eq_p` (fermion &neu, fermion &p, double sig, double ome, double rho, double &f1, double &f2, double &f3, thermo &th)**

Probably best to have f1, f2, and f3 scaled in some sensible way, i.e. scaled to the fields?

**Class `rmf_nucleus`**

Sort energy levels at the end by energy

Improve the numerical methods

Make the neutron and proton orbitals more configurable

Generalize to  $m_n \neq m_p$ .

Allow more freedom in the integrations

Consider converting everything to inverse fermis.

Convert to zero-indexed arrays

Warn when the level ordering is wrong, and unoccupied levels are lower energy than occupied levels

**Class `skyrme_eos`**

There is some code duplication between `calc_e()` and `calc_temp_e()` which could be simplified.

**Class `stos_eos`**

Add the T=0 and Ye=0 data to this class

**Class `tov_buchdahl_eos`**

Figure out what to do with the `buchfun()` function

## 12 Todo List

**Global `cfl_njl_eos::calc_eq_temp_p` (quark &u, quark &d, quark &s, double &qq1, double &qq2, double &qq3, double &gap1, double &gap2, double &gap3, double mu3, double mu8, double &n3, double &n8, thermo &qb, double temper)**

It surprises me that n3 is not -res[11]. Is there a sign error in the color densities?

**Global `cfl_njl_eos::gapped_eigenvalues` (double m1, double m2, double lmom, double mu1, double mu2, double tdelta, double lam[4], double dldmu1[4], double dldmu2[4], double dldm1[4], double dldm2[4], double dldg[4])**

In the code, the equal mass case seems to be commented out. Why?

**Global `ddc_eos::calc_eq_e` (fermion &neu, fermion &p, double sig, double ome, double rho, double &f1, double &f2, double &f3, thermo &th)**

Is the thermodynamic identity is satisfied even when the field equations are not solved? Check this.

**Class `gen_sn_eos`**

Ensure all chemical potentials are based on the same rest masses?

Allow logarithmic grids for any of nb, Ye, or T.

**Class `ldrop_mass_skin`**

This is based on LPRL, but it's a little different in Lattimer and Swesty. I should document what the difference is.

The testing could be updated.

**Class `ls_eos`**

There are still a few points for which the electron/photon EOS seems to be off.

**Class `rmf_eos`**

The functions `fcomp_fields()`, `fkprime_fields()`, and `fesym_fields()` are not quite correct if the neutron and proton masses are different. For this reason, they are currently unused by `saturation()`.

- The `fix_saturation()` and `calc_cr()` functions use `mnuc`, and should be modified to allow different neutron and proton masses.
- Check the formulas in the "Background" section
- There are two `calc_e()` functions that solve. One is specially designed to work without a good initial guess. Possibly the other `calc_e()` function should be similarly designed?
- Make sure that this class properly handles particles for which `inc_rest_mass` is true/false
- The error handler is called sometimes when `calc_e()` is used to compute pure neutron matter. This should be fixed.
- Decide whether to throw an error at [Ref. 1].
- Put the `err_nonconv` system into `calc_p()`, `calc_temp_e()` and `fix_saturation()`, etc.

**Global `rmf_eos::check_naturalness (rmf_eos &re)`**

I may have ignored some signs in the above, which are unimportant for this application, but it would be good to fix them for posterity.

**Global `rmf_eos::fix_saturation (double guess_cs=4.0, double guess_cw=3.0, double guess_b=0.001, double guess_c=-0.001)`**

Fix this for `zm_mode=true`

- Ensure solver is more robust

**Global `rmf_eos::fkprime_fields (double sig, double ome, double nb, double &k, double &kprime)`**

Does this work? Fix `fkprime_fields()` if it does not.

**Global `rmf_eos::n_charge`**

Should use `hadronic_eos::proton_frac` instead?

**Class `rmf_nucleus`**

Better documentation

Convert `energies()` to use EOS and possibly replace `sigma_rhs()` and related functions by the associated field equation method of `rmf_eos`.

Document `hw=3.923+23.265/cbrt(atot)`;

**Global `schematic_eos::set_a_from_mstar (double u_msom, double mnuc)`**

This was computed in `schematic_sym.nb`, which might be added to the documentation?

**Class `sht_eos`**

More documentation

**Class `skyrme_eos`**

Make sure that this class properly handles particles for which `inc_rest_mass` is true/false

- What about the spin-orbit units?
- Need to write a function that calculates saturation density?
- Remove use of `mnuc` in `calparfun()`?
- The compressibility could probably use some simplification
- Make sure the finite-temperature part is properly tested
- The testing code doesn't work if `err_mode` is 2, probably because of problems in `load()`.

- Document load() file format.
- Update reference list.

Global `skyrme_eos::calpar` (double `gt0=-10.0`, double `gt3=70.0`, double `galpha=0.2`, double `gt1=2.0`, double `gt2=-1.0`)

Does this work for both 'a' and 'b' non-zero?

Compare to similar formulae from [Margueron02](#)

Global `skyrme_eos::landau_neutron` (double `n0`, double `m`, double `&f0`, double `&g0`, double `&f1`, double `&g1`)

This needs to be checked

Global `skyrme_eos::landau_nuclear` (double `n0`, double `m`, double `&f0`, double `&g0`, double `&f0p`, double `&g0p`, double `&f1`, double `&g1`, double `&f1p`, double `&g1p`)

This needs to be checked.

Class `toy_eos`

Fix `read_table_file` and maybe `set_low_density_eos()`.

Class `toy_interp_eos`

Warn that the pressure in the low-density eos is not strictly increasing! (see at  $P=4.3e-10$ )

It might be useful to exit more gracefully when non-finite values are obtained in interpolation, analogous to the `err_nonconv` mechanism elsewhere.

Class `toy_solve`

baryon mass doesn't work for `fixed()` (This may be fixed. We should make sure it's tested.)

- Combine `maxoutsize` and `kmax`?
- Document column naming issues
- Document surface gravity and redshift
- Standardize `xmev_kg`, etc.
- Use `convert_units`?
- Double check that `fixed()` doesn't give a solution on the unstable branch
- Ensure that this class copies over the units of the user-specified columns to the table

## 13 Bug List

Class `gen_potential_eos`

The BGBD EOS doesn't work and the effective mass for the GBD EOS doesn't work

Class `sym4_eos_base`

Testing was disabled in HDF conversion. Fix this.

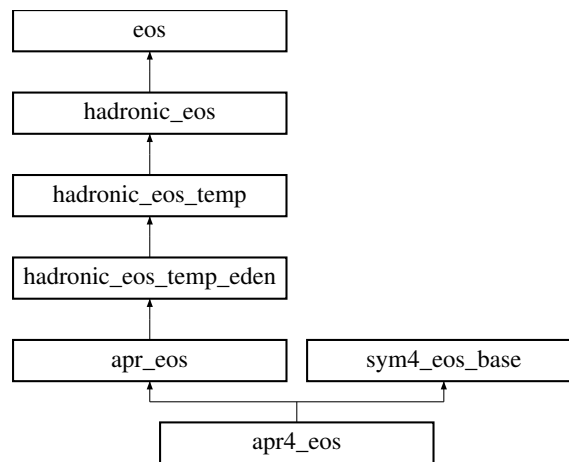
## 14 Data Structure Documentation

### 14.1 `apr4_eos` Class Reference

A version of `apr_eos` to separate potential and kinetic contributions.

```
#include <sym4_eos.h>
```

Inheritance diagram for `apr4_eos`:



### 14.1.1 Detailed Description

#### References:

Created for [Steiner06](#).

Definition at line 125 of file `sym4_eos.h`.

#### Public Member Functions

- virtual int [calc\\_e\\_sep](#) (**fermion** &ne, **fermion** &pr, double &ed\_kin, double &ed\_pot, double &mu\_n\_kin, double &mu\_p\_kin, double &mu\_n\_pot, double &mu\_p\_pot)  
*Compute the potential and kinetic parts separately.*

The documentation for this class was generated from the following file:

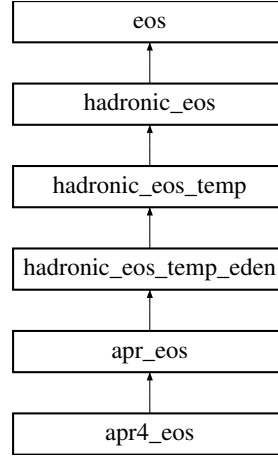
- `sym4_eos.h`

## 14.2 apr\_eos Class Reference

EOS from Akmal, Pandharipande, and Ravenhall.

```
#include <apr_eos.h>
```

Inheritance diagram for `apr_eos`:



### 14.2.1 Detailed Description

The EOS of Akmal, Pandharipande, and Ravenhall, from [Akmal98](#) (APR).

The Hamiltonian is:

$$\mathcal{H}_{APR} = \mathcal{H}_{kin} + \mathcal{H}_{pot}$$

$$\mathcal{H}_{kin} = \left( \frac{\hbar^2}{2m} + (p_3 + (1-x)p_5)ne^{-p_4n} \right) \tau_n + \left( \frac{\hbar^2}{2m} + (p_3 + xp_5)ne^{-p_4n} \right) \tau_p$$

$$\mathcal{H}_{pot} = g_1 \left( 1 - (1-2x)^2 \right) + g_2 (1-2x)^2$$

The following are definitions for  $g_i$  in the low-density phase (LDP) or the high-density phase (HDP):

$$g_{1,LDP} = -n^2 \left( p_1 + p_2n + p_6n^2 + (p_{10} + p_{11}n)e^{-p_9^2n^2} \right)$$

$$g_{2,LDP} = -n^2 \left( p_{12}/n + p_7 + p_8n + p_{13}e^{-p_9^2n^2} \right)$$

$$g_{1,HDP} = g_{1,LDP} - n^2 \left( p_{17}(n - p_{19}) + p_{21}(n - p_{19})^2 e^{p_{18}(n - p_{19})} \right)$$

$$g_{2,HDP} = g_{2,LDP} - n^2 \left( p_{15}(n - p_{20}) + p_{14}(n - p_{20})^2 e^{p_{16}(n - p_{20})} \right)$$

The chemical potentials include the rest mass energy and the energy density includes the rest mass energy density.

#### Note

APR seems to have been designed to be used with non-relativistic neutrons and protons with equal masses of 939 MeV. This gives a saturation density very close to 0.16.

The variables  $v_n$  and  $v_p$  contain the expressions  $(-\mu_n + V_n)/T$  and  $(-\mu_p + V_p)/T$  respectively, where  $V$  is the potential part of the single particle energy for particle  $i$  (i.e. the derivative of the Hamiltonian w.r.t. density while energy density held constant). Equivalently,  $v_n$  is just  $-k_{Fn}^2/2m^*$ .

The selection between the LDP and HDP is controlled by [pion](#). The default is to use the LDP at densities below  $0.16 \text{ fm}^{-3}$ , and for larger densities to just use whichever minimizes the energy.

The finite temperature approximations from [Prakash97](#) are used in testing.

## Note

Since this EOS uses the effective masses and chemical potentials in the fermion class, the values of **part::non\_interacting** for neutrons and protons are set to false in many of the functions.

The parameter array is unit indexed, so that `par[0]` is unused.

**Idea for Future** There might be room to improve the testing of the finite temperature part a bit.

There is some repetition between `calc_e()` and `calc_temp_e()` that possibly could be removed.

Definition at line 123 of file `apr_eos.h`.

## Public Member Functions

- `apr_eos()`  
*Create an EOS object with the default parameter set ( $A18 + UIX^* + \delta v$ ).*
- virtual int `calc_e(fermion &n, fermion &p, thermo &th)`  
*Equation of state as a function of density.*
- virtual int `calc_temp_e(fermion &n, fermion &p, double temper, thermo &th)`  
*Equation of state as a function of densities.*
- double `fcomp(double nb)`  
*Compute the compressibility.*
- double `fesym_diff(double nb)`  
*Calculate symmetry energy of matter as energy of neutron matter minus the energy of nuclear matter.*
- void `select(int model_index)`  
*Select model.*
- int `gradient_qij2(double nn, double np, double &qnn, double &qnp, double &qpp, double &dqnnndnn, double &dqnpndnp, double &dqnpdnn, double &dqppdnn, double &dqppdnp)`  
*Calculate  $Q$ 's for semi-infinite nuclear matter.*
- double `get_par(int n)`  
*Get the value of one of the parameters.*
- int `set_par(int n, double x)`  
*Set the value of one of the parameters.*
- virtual const char \* `type()`  
*Return string denoting type ("apr\_eos")*

## Data Fields

- bool `parent_method`  
*If true, use the methods from `hadronic_eos` for `fcomp()`*

## Protected Attributes

- double \* `par`  
*Storage for the parameters.*
- int `lp`  
*An integer to indicate which phase was used in `calc_e()`*
- `nonrel_fermion nrf`  
*Desc.*
- int `choice`  
*The variable indicating which parameter set is to be used.*



## Choice of phase

- static const int `best` = 0  
use LDP for densities less than 0.16 and for higher densities, use the phase which minimizes energy (default)
- static const int `ldp` = 1  
LDP (no pion condensation)
- static const int `hdp` = 2  
HDP (pion condensation)
- int `pion`  
Choice of phase (default `best`)
- int `last_phase` ()  
Return the phase of the most recent call to `calc_e()`

## 14.2.2 Member Function Documentation

## 14.2.2.1 double apr\_eos::fcomp ( double nb )

See general notes at `hadronic_eos::fcomp()`. This computes the compressibility (at fixed proton fraction = 0.5) exactly, unless `parent_method` is true in which case the derivative is taken numerically in `hadronic_eos::fcomp()`.

## 14.2.2.2 double apr\_eos::fesym\_diff ( double nb ) [virtual]

This function returns the energy per baryon of neutron matter minus the energy per baryon of nuclear matter. This will deviate significantly from the results from `fesym()` only if the dependence of the symmetry energy on  $\delta$  is not quadratic.

Reimplemented from `hadronic_eos`.

## 14.2.2.3 void apr\_eos::select ( int model\_index )

Valid values for `model_index` are:

- 1 - A18+UIX\*+deltav (preferred by Akmal, et. al. - this is the default)
- 2 - A18+UIX\*
- 3 - A18+deltav
- 4 - A18

If any other integer is given, A18+UIX\*+deltav is assumed.

## 14.2.2.4 int apr\_eos::gradient\_qij2 ( double nn, double np, double &amp; qnn, double &amp; qnp, double &amp; qpp, double &amp; dqnnn, double &amp; dqnpdp, double &amp; dqnpdnn, double &amp; dqnpdnp, double &amp; dqppdnn, double &amp; dqppdnp )

For general discussion, see the documentation to `hadronic_eos::qs()`.

For APR, we set  $x_1 = x_2 = 0$  so that  $Q_i = P_i/2$  and then

$$\begin{aligned} P_1 &= \left( \frac{1}{2} p_3 - p_5 \right) e^{-p_4 n} \\ P_2 &= \left( \frac{1}{2} p_3 + p_5 \right) e^{-p_4 n} \end{aligned}$$

This gives

$$\begin{aligned} Q_{nn} &= \frac{1}{4} e^{-p_4 \rho} [-6p_5 - p_4(p_3 - 2p_5)(n_n + 2n_p)] \\ Q_{np} &= \frac{1}{8} e^{-p_4 \rho} [4(p_3 - 4p_5) - 3p_4(p_3 - 2p_5)(n_n + n_p)] \\ Q_{pp} &= \frac{1}{4} e^{-p_4 \rho} [-6p_5 - p_4(p_3 - 2p_5)(n_p + 2n_n)] \end{aligned}$$

See the Mathematica notebook

doc/o2scl/extras/apr\_eos.nb  
doc/o2scl/extras/apr\_eos.ps

### 14.2.3 Field Documentation

#### 14.2.3.1 bool apr\_eos::parent\_method

This can be set to true to check the difference in the compressibility between the exact expressions and the numerical values from class [hadronic\\_eos](#).

**Idea for Future** This function is probably unnecessary, as the syntax

```
apr_eos apr;
ccout << apr.hadronic_eos::fcomp(0.16) << endl;
```

works just as well.

Definition at line 282 of file apr\_eos.h.

The documentation for this class was generated from the following file:

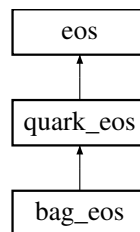
- apr\_eos.h

## 14.3 bag\_eos Class Reference

Simple bag model.

```
#include <bag_eos.h>
```

Inheritance diagram for bag\_eos:



### 14.3.1 Detailed Description

An equation of state with  $P = -B + P_{u,FG} + P_{d,FG} + P_{s,FG}$  where  $P_{i,FG}$  is the Fermi gas contribution from particle  $i$  and  $B$  is a density- and temperature-independent bag constant.

The finite temperature functions run the zero temperature code if the temperature is less than or equal to 0.

Definition at line 45 of file bag\_eos.h.

#### Public Member Functions

- virtual int [calc\\_p](#)(**quark** &u, **quark** &d, **quark** &s, **thermo** &th)  
*Calculate equation of state as a function of chemical potentials.*

- virtual int `calc_e` (`quark &u`, `quark &d`, `quark &s`, `thermo &th`)  
*Calculate equation of state as a function of density.*
- virtual int `calc_temp_p` (`quark &u`, `quark &d`, `quark &s`, double `temper`, `thermo &th`)  
*Calculate equation of state as a function of the chemical potentials.*
- virtual int `calc_temp_e` (`quark &u`, `quark &d`, `quark &s`, double `temper`, `thermo &th`)  
*Calculate equation of state as a function of the densities.*
- virtual const char \* `type` ()  
*Return string denoting type ("bag\_eos")*

#### Data Fields

- double `bag_constant`  
*The bag constant in  $\text{fm}^{-4}$  (default  $200/(\hbar c)$ ).*

#### 14.3.2 Member Function Documentation

14.3.2.1 virtual int `bag_eos::calc_temp_p` ( `quark &u`, `quark &d`, `quark &s`, double `temper`, `thermo &th` ) [virtual]

This function returns zero (success) unless the call to `quark::pair_mu()` fails.

Reimplemented from `quark_eos`.

14.3.2.2 virtual int `bag_eos::calc_temp_e` ( `quark &u`, `quark &d`, `quark &s`, double `temper`, `thermo &th` ) [virtual]

This function returns zero (success) unless the call to `quark::pair_density()` fails.

Reimplemented from `quark_eos`.

The documentation for this class was generated from the following file:

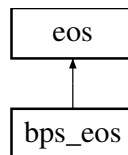
- `bag_eos.h`

## 14.4 bps\_eos Class Reference

Baym-Pethick-Sutherland equation of state.

```
#include <bps_eos.h>
```

Inheritance diagram for `bps_eos`:



#### 14.4.1 Detailed Description

This calculates the equation of state of electrons and nuclei using the approach of [Baym71](#) (based on the discussion in [Shapiro83](#)) between about  $8 \times 10^6 \text{ g/cm}^3$  and  $4.3 \times 10^{11} \text{ g/cm}^3$ . Below these densities, more complex Coulomb corrections need to be considered, and above these densities, neutron drip is important.

The default mass formula is semi-empirical

$$M(A, Z) = (A - Z)m_n + Z(m_p + m_e) - 15.76A - 17.81A^{2/3} - 0.71Z^2/A^{1/3} - 94.8/A(A/2 - Z)^2 + E_{\text{pair}}$$

where

$$E_{\text{pair}} = \pm 39/A^{3/4}$$

if the nucleus is odd-odd (plus sign) or even-even (minus sign) and  $E_{\text{pair}}$  is zero for odd-even and even-odd nuclei. The nuclei are assumed not to contribute to the pressure. The electronic contribution to the pressure is assumed to be equal to the Fermi gas contribution plus a "lattice" contribution

$$\varepsilon_L = -1.444Z^{2/3}e^2n_e^{4/3}$$

This is Eq. 2.7.2 in [Shapiro83](#). The rest mass energy of the nucleons is included in the energy density.

The original results from [Baym71](#) are stored as a **table** in file `data/o2scl/bps.eos`. The testing code for this class compares the calculations to the table and matches to within .2 percent for the energy density and 9 percent for the pressure (for a fixed baryon number density).

**Idea for Future** Can the pressure be made to match more closely?

Convert to a [hadronic\\_eos](#) object and offer an associated interface?

Definition at line 81 of file `bps_eos.h`.

#### Public Member Functions

- virtual int [calc\\_density](#) (double barn, **thermo** &th, int &Z, int &A)  
*Calculate the equation of state as a function of the baryon number density barn.*
- virtual int [calc\\_pressure](#) (**thermo** &th, double &barn, int &Z, int &A)  
*Calculate the equation of state as a function of the pressure.*
- virtual double [lattice\\_energy](#) (int Z)  
*The electron lattice energy.*
- virtual const **fermion** & [get\\_electron](#) ()  
*Get a pointer to the electron.*
- virtual double [mass\\_formula](#) (int Z, int A)  
*The mass formula.*
- virtual const char \* [type](#) ()  
*Return string denoting type ("bps\_eos")*
- int [set\\_mass\\_formula](#) (**nuclear\_mass** &nm)  
*Set the nuclear mass formula to be used.*
- int [calc\\_density\\_fixedA](#) (double barn, **thermo** &th, int &Z, int A)  
*Compute the ground state assuming a fixed atomic number.*

#### Data Fields

- **semi\_empirical\_mass** [def\\_mass](#)  
*Default mass formula.*
- **fermion** [e](#)  
*The electron thermodynamics.*

#### Protected Member Functions

- virtual int [eq274](#) (size\_t nv, const **ovector\_base** &nx, **ovector\_base** &ny, int &Zt)  
*Solve Equation 2.7.4 for a given pressure.*
- double [gibbs](#) (int Z, int A)  
*The Gibbs free energy.*
- double [energy](#) (double barn, int Z, int A)  
*The energy density.*

## Protected Attributes

- `fermion_zerot` [fzt](#)  
*Desc.*
- `gsl_mroot_hybrids` < `mm_funct` <> > [gs](#)  
*A solver to solve Eq. 2.7.4.*
- `nuclear_mass` \* [nmp](#)  
*The nuclear mass formula.*

## 14.4.2 Member Function Documentation

14.4.2.1 `virtual int bps_eos::calc_density ( double barn, thermo & th, int & Z, int & A )` `[virtual]`

This calculates the equation of state as a function of the baryon number density in  $\text{fm}^{-3}$ , returning the representative nucleus with proton number *Z* and atomic number *A*. The pressure and energy density are returned in *th* in  $\text{fm}^{-4}$ .

14.4.2.2 `virtual int bps_eos::calc_pressure ( thermo & th, double & barn, int & Z, int & A )` `[virtual]`

This calculates the equation of state as a function of the pressure, returning the representative nucleus with proton number *Z* and atomic number *A* and the baryon number density *barn* in  $\text{fm}^{-3}$ . The energy density is also returned in  $\text{fm}^{-4}$  in *th*.

14.4.2.3 `virtual double bps_eos::mass_formula ( int Z, int A )` `[virtual]`

The nuclear mass without the contribution of the rest mass of the electrons. The electron rest mass energy is included in the electron thermodynamics elsewhere.

## 14.4.3 Field Documentation

14.4.3.1 `fermion bps_eos::e`

## Note

The electron rest mass is included by default in the energy density and the chemical potential

Definition at line 145 of file `bps_eos.h`.

The documentation for this class was generated from the following file:

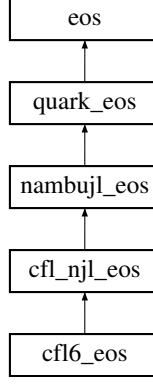
- `bps_eos.h`

## 14.5 cfl6\_eos Class Reference

An EOS like [cfl\\_njl\\_eos](#) but with a color-superconducting 't Hooft interaction.

```
#include <cfl6_eos.h>
```

Inheritance diagram for `cfl6_eos`:



### 14.5.1 Detailed Description

Beginning with the Lagrangian:

$$\mathcal{L} = \mathcal{L}_{Dirac} + \mathcal{L}_{NJL} + \mathcal{L}_{tHooft} + \mathcal{L}_{SC} + \mathcal{L}_{SC6}$$

$$\mathcal{L}_{Dirac} = \bar{q} (i\partial - m - \mu\gamma^0) q$$

$$\mathcal{L}_{NJL} = G_S \sum_{a=0}^8 \left[ (\bar{q} \lambda^a q)^2 - (\bar{q} \lambda^a \gamma^5 q)^2 \right]$$

$$\mathcal{L}_{tHooft} = G_D \left[ \det_f \bar{q} (1 - \gamma^5) q + \det_f \bar{q} (1 + \gamma^5) q \right]$$

$$\mathcal{L}_{SC} = G_{DIQ} \left( \bar{q}_{i\alpha} i\gamma^5 \epsilon^{ijk} \epsilon^{\alpha\beta\gamma} q_{j\beta}^C \right) \left( \bar{q}_{\ell\delta} i\gamma^5 \epsilon^{\ell mk} \epsilon^{\delta\epsilon\gamma} q_{m\epsilon}^C \right)$$

$$\mathcal{L}_{SC6} = K_D \left( \bar{q}_{i\alpha} i\gamma^5 \epsilon^{ijk} \epsilon^{\alpha\beta\gamma} q_{j\beta}^C \right) \left( \bar{q}_{\ell\delta} i\gamma^5 \epsilon^{\ell mn} \epsilon^{\delta\epsilon\eta} q_{m\epsilon}^C \right) (\bar{q}_{k\gamma} q_{n\eta})$$

We can simplify the relevant terms in  $\mathcal{L}_{NJL}$ :

$$\mathcal{L}_{NJL} = G_S \left[ (\bar{u}u)^2 + (\bar{d}d)^2 + (\bar{s}s)^2 \right]$$

and in  $\mathcal{L}_{tHooft}$ :

$$\mathcal{L}_{NJL} = G_D (\bar{u}u \bar{d}d \bar{s}s)$$

Using the definition:

$$\Delta^{k\gamma} = \langle \bar{q} i\gamma^5 \epsilon \epsilon q^C \rangle$$

and the ansatzes:

$$(\bar{q}_1 q_2)(\bar{q}_3 q_4) \rightarrow \bar{q}_1 q_2 \langle \bar{q}_3 q_4 \rangle + \bar{q}_3 q_4 \langle \bar{q}_1 q_2 \rangle - \langle \bar{q}_1 q_2 \rangle \langle \bar{q}_3 q_4 \rangle$$

$$(\bar{q}_1 q_2)(\bar{q}_3 q_4)(\bar{q}_5 q_6) \rightarrow \bar{q}_1 q_2 \langle \bar{q}_3 q_4 \rangle \langle \bar{q}_5 q_6 \rangle + \bar{q}_3 q_4 \langle \bar{q}_1 q_2 \rangle \langle \bar{q}_5 q_6 \rangle + \bar{q}_5 q_6 \langle \bar{q}_1 q_2 \rangle \langle \bar{q}_3 q_4 \rangle - 2 \langle \bar{q}_1 q_2 \rangle \langle \bar{q}_3 q_4 \rangle \langle \bar{q}_5 q_6 \rangle$$

for the mean field approximation, we can rewrite the Lagrangian

$$\mathcal{L}_{NJL} = 2G_S \left[ (\bar{u}u) \langle \bar{u}u \rangle + (\bar{d}d) \langle \bar{d}d \rangle + (\bar{s}s) \langle \bar{s}s \rangle - \langle \bar{u}u \rangle^2 - \langle \bar{d}d \rangle^2 - \langle \bar{s}s \rangle^2 \right]$$

$$\mathcal{L}_{tHooft} = -2G_D \left[ (\bar{u}u) \langle \bar{u}u \rangle \langle \bar{s}s \rangle + (\bar{d}d) \langle \bar{u}u \rangle \langle \bar{s}s \rangle + (\bar{s}s) \langle \bar{u}u \rangle \langle \bar{d}d \rangle - 2 \langle \bar{u}u \rangle \langle \bar{d}d \rangle \langle \bar{s}s \rangle \right]$$

$$\mathcal{L}_{SC} = G_{DIQ} \left[ \Delta^{k\gamma} \left( \bar{q}_{\ell\delta} i\gamma^5 \epsilon^{\ell mk} \epsilon^{\delta\epsilon\gamma} q_{m\epsilon}^C \right) + \left( \bar{q}_{i\alpha} i\gamma^5 \epsilon^{ijk} \epsilon^{\alpha\beta\gamma} q_{j\beta}^C \right) \Delta^{k\gamma\dagger} - \Delta^{k\gamma} \Delta^{k\gamma\dagger} \right]$$

$$\mathcal{L}_{SC6} = K_D \left[ (\bar{q}_{m\epsilon} q_{n\eta}) \Delta^{k\gamma} \Delta^{m\epsilon\dagger} + \left( \bar{q}_{i\alpha} i\gamma^5 \epsilon^{ijk} \epsilon^{\alpha\beta\gamma} q_{j\beta}^C \right) \Delta^{m\epsilon\dagger} \langle \bar{q}_{m\epsilon} q_{n\eta} \rangle \right]$$

$$+ K_D \left[ \Delta^{k\gamma} \left( \bar{q}_{\ell\delta} i\gamma^5 \epsilon^{\ell mn} \epsilon^{\delta\epsilon\eta} q_{m\epsilon}^C \right) \langle \bar{q}_{m\epsilon} q_{n\eta} \rangle - 2 \Delta^{k\gamma} \Delta^{m\epsilon\dagger} \langle \bar{q}_{m\epsilon} q_{n\eta} \rangle \right]$$

If we make the definition  $\tilde{\Delta} = 2G_{D1Q}\Delta$

---

### References:

Created for [Steiner05](#).

Definition at line 187 of file cfl6\_eos.h.

### Public Member Functions

- virtual int [calc\\_eq\\_temp\\_p](#) (**quark** &u, **quark** &d, **quark** &s, double &qq1, double &qq2, double &qq3, double &gap1, double &gap2, double &gap3, double mu3, double mu8, double &n3, double &n8, **thermo** &qb, double [temper](#))  
*Calculate the EOS.*
- virtual int [integrands](#) (double p, double res[])  
*The momentum integrands.*
- virtual int [test\\_derivatives](#) (double lmom, double mu3, double mu8, **test\_mgr** &t)  
*Check the derivatives specified by [eigenvalues\(\)](#)*
- virtual int [eigenvalues6](#) (double lmom, double mu3, double mu8, double egv[36], double dedmuu[36], double dedmud[36], double dedmus[36], double dedmu[36], double dedmd[36], double dedms[36], double dedu[36], double dedd[36], double deds[36], double dedmu3[36], double dedmu8[36])  
*Calculate the energy eigenvalues and their derivatives.*
- virtual int [make\\_matrices](#) (double lmom, double mu3, double mu8, double egv[36], double dedmuu[36], double dedmud[36], double dedmus[36], double dedmu[36], double dedmd[36], double dedms[36], double dedu[36], double dedd[36], double deds[36], double dedmu3[36], double dedmu8[36])  
*Construct the matrices, but don't solve the eigenvalue problem.*
- virtual const char \* [type](#) ()  
*Return string denoting type ("cfl6\_eos")*

### Data Fields

- double [KD](#)  
*The color superconducting 't Hooft coupling (default 0)*
- double [kdlimit](#)  
*The absolute value below which the CSC 't Hooft coupling is ignored(default  $10^{-6}$ )*

### Protected Member Functions

- int [set\\_masses](#) ()  
*Set the quark effective masses from the gaps and the condensates.*

### Protected Attributes

- **omatrix\_cx** [iprop6](#)  
*Storage for the inverse propagator.*
  - **omatrix\_cx** [eivec6](#)  
*The eigenvectors.*
  - **omatrix\_cx** [dipdgapu](#)  
*The derivative wrt the ds gap.*
  - **omatrix\_cx** [dipdgapd](#)  
*The derivative wrt the us gap.*
  - **omatrix\_cx** [dipdgaps](#)  
*The derivative wrt the ud gap.*
  - **omatrix\_cx** [dipdqqu](#)  
*The derivative wrt the up quark condensate.*
  - **omatrix\_cx** [dipdqqd](#)  
*The derivative wrt the down quark condensate.*
-

- **omatrix\_cx** [dipdqqs](#)  
*The derivative wrt the strange quark condensate.*
- **ovector** [eval6](#)  
*Storage for the eigenvalues.*
- **gsl\_eigen\_hermv\_workspace** \* [w6](#)  
*GSL workspace for the eigenvalue computation.*

#### Static Protected Attributes

- static const int [mat\\_size](#) = 36  
*The size of the matrix to be diagonalized.*

#### Private Member Functions

- **cfl6\_eos** (const [cfl6\\_eos](#) &)
- **cfl6\_eos** & **operator=** (const [cfl6\\_eos](#) &)

#### 14.5.2 Member Function Documentation

14.5.2.1 `virtual int cfl6_eos::calc_eq_temp_p ( quark & u, quark & d, quark & s, double & qq1, double & qq2, double & qq3, double & gap1, double & gap2, double & gap3, double mu3, double mu8, double & n3, double & n8, thermo & qb, double temper ) [virtual]`

Calculate the EOS from the quark condensates. Return the mass gap equations in qq1, qq2, qq3, and the normal gap equations in gap1, gap2, and gap3.

Using `fromqq=true` as in [nambu\\_jl\\_eos](#) and `nambu_jl_temp_eos` does not work here and will return an error.

If all of the gaps are less than `gap_limit`, then the `nambu_jl_temp_eos::calc_temp_p()` is used, and `gap1`, `gap2`, and `gap3` are set to equal `u.del`, `d.del`, and `s.del`, respectively.

Reimplemented from [cfl\\_njl\\_eos](#).

14.5.2.2 `virtual int cfl6_eos::eigenvalues6 ( double lmom, double mu3, double mu8, double egv[36], double dedmuu[36], double dedmud[36], double dedmus[36], double dedmu[36], double dedmd[36], double dedms[36], double dedu[36], double dedd[36], double deds[36], double dedmu3[36], double dedmu8[36] ) [virtual]`

Given the momentum `mom`, and the chemical potentials associated with the third and eighth gluons (`mu3` and `mu8`), this computes the eigenvalues of the inverse propagator and the associated derivatives.

Note that this is not the same as [cfl\\_njl\\_eos::eigenvalues\(\)](#) which returns `dedmu` rather `dedqqu`.

14.5.2.3 `virtual int cfl6_eos::make_matrices ( double lmom, double mu3, double mu8, double egv[36], double dedmuu[36], double dedmud[36], double dedmus[36], double dedmu[36], double dedmd[36], double dedms[36], double dedu[36], double dedd[36], double deds[36], double dedmu3[36], double dedmu8[36] ) [virtual]`

This is used by `check_derivatives()` to make sure that the derivative entries are right.

The documentation for this class was generated from the following file:

- `cfl6_eos.h`

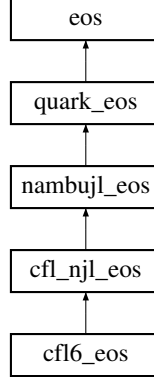
## 14.6 cfl\_njl\_eos Class Reference

Nambu Jona-Lasinio model with a schematic CFL di-quark interaction at finite temperature.

```
#include <cfl_njl_eos.h>
```

Inheritance diagram for `cfl_njl_eos`:





#### 14.6.1 Detailed Description

The variable B0 must be set before use.

The original Lagrangian is

$$\mathcal{L} = \mathcal{L}_{\text{Dirac}} + \mathcal{L}_{4\text{-fermion}} + \mathcal{L}_{6\text{-fermion}} + \mathcal{L}_{\text{CSC1}} + \mathcal{L}_{\text{CSC2}}$$

$$\mathcal{L}_{\text{Dirac}} = \bar{q}_{i\alpha} (i\partial \delta_{ij} \delta_{\alpha\beta} - m_{ij} \delta_{\alpha\beta} - \mu_{ij, \alpha\beta} \gamma^0) q_{j\beta}$$

$$\mathcal{L}_{4\text{-fermion}} = G_S \sum_{a=0}^8 \left[ (\bar{q} \lambda_f^a q)^2 + (\bar{q} i \gamma_5 \lambda_f^a q)^2 \right]$$

$$\mathcal{L}_{6\text{-fermion}} = -G_D \left[ \det_{ij} \bar{q}_{i\alpha} (1 + i\gamma_5) q_{j\beta} + \det_{ij} \bar{q}_{i\alpha} (1 - i\gamma_5) q_{j\beta} \right] \delta_{\alpha\beta}$$

$$\mathcal{L}_{\text{CSC1}} = G_{DIQ} \sum_k \sum_\gamma \left[ \left( \bar{q}_{i\alpha} \epsilon_{ijk} \epsilon_{\alpha\beta\gamma} q_{j\beta}^C \right) \left( \bar{q}_{i'\alpha'}^C \epsilon_{i'j'k} \epsilon_{\alpha'\beta'\gamma} q_{j'\beta'} \right) \right]$$

$$\mathcal{L}_{\text{CSC2}} = G_{DIQ} \sum_k \sum_\gamma \left[ \left( \bar{q}_{i\alpha} i\gamma_5 \epsilon_{ijk} \epsilon_{\alpha\beta\gamma} q_{j\beta}^C \right) \left( \bar{q}_{i'\alpha'}^C i\gamma_5 \epsilon_{i'j'k} \epsilon_{\alpha'\beta'\gamma} q_{j'\beta'} \right) \right],$$

where  $\mu$  is the quark number chemical potential. couplings  $G_S$ ,  $G_D$ , and  $G_{DIQ}$  ultra-violet three-momentum cutoff,  $\Lambda$

The thermodynamic potential is

$$\Omega(\mu_i, \langle \bar{q}q \rangle_i, \langle qq \rangle_i, T) = \Omega_{\text{vac}} + \Omega_{\text{stat}} + \Omega_0$$

where  $i$  runs over all nine (three colors times three flavors) quarks. We assume that the condensates are independent of color and that the quark chemical potentials are of the form  $\mu_Q = \mu_{\text{Flavor}(Q)} + \mu_{\text{Color}(Q)}$  with

$$\mu_{\text{red}} = \mu_3 + \mu_8/\sqrt{3} \quad \mu_{\text{green}} = -\mu_3 + \mu_8/\sqrt{3} \quad \mu_{\text{blue}} = -2\mu_8/\sqrt{3}$$

With these assumptions, the thermodynamic potential as given by the function thd\_potential(), is a function of 12 variables

$$\Omega(\mu_u, \mu_d, \mu_s, \mu_3, \mu_8, \langle \bar{u}u \rangle, \langle \bar{d}d \rangle, \langle \bar{s}s \rangle, \langle ud \rangle, \langle us \rangle, \langle ds \rangle, T)$$

The individual terms are

$$\Omega_{\text{stat}} = -\frac{1}{2} \int \frac{d^3p}{(2\pi)^3} \sum_{i=1}^{72} \left[ \frac{\lambda_i}{2} + T \ln \left( 1 + e^{-\lambda_i/T} \right) \right]$$

$$\Omega_{\text{vac}} = -2G_S \sum_{i=u,d,s} \langle \bar{q}_i q_i \rangle^2 + 4G_D \langle \bar{u} u \rangle \langle \bar{d} d \rangle \langle \bar{s} s \rangle + \sum_k \sum_\gamma \frac{|\Delta^k \gamma|^2}{4G_{D1Q}}$$

where  $\lambda_i$  are the eigenvalues of the (72 by 72) matrix (calculated by the function `eigenvalues()`)

$$D = \begin{bmatrix} -\gamma^0 \vec{\gamma} \cdot \vec{p} - M_i \gamma^0 + \mu_{i\alpha} & \Delta_i \gamma^0 \gamma_5 C \\ i \Delta_i \gamma^0 C \gamma_5 & -\gamma^0 \vec{\gamma}^T \cdot \vec{p} + M_i \gamma^0 - \mu_{i\alpha} \end{bmatrix}$$

and  $C$  is the charge conjugation matrix (in the Dirac representation).

The values of the various condensates are usually determined by the condition

$$\frac{\partial \Omega}{\partial \langle \bar{q} q \rangle_i} = 0 \quad \frac{\partial \Omega}{\partial \langle q q \rangle_i} = 0$$

Note that setting `fixed_mass` to `true` and setting all of the gaps to zero when `gap_limit` is less than zero will reproduce an analog of the bag model with a momentum cutoff.

The variable `nambu_njl_eos::fromqq` is automatically set to `true` in the constructor, as computations with `fromqq=false` are not implemented.

**Idea for Future** This class internally mixes `ovector`, `omatrix`, `gsl_vector` and `gsl_matrix` objects in a confusing and non-optimal way. Fix this.

Allow user to change derivative object? This isn't possible right now because the stepsize parameter of the derivative object is used.

---

## References:

Created for [Steiner02](#).

Definition at line 207 of file `cfl_njl_eos.h`.

## Public Member Functions

- virtual int `set_parameters` (double `lambda`=0.0, double `fourferm`=0.0, double `sixferm`=0.0, double `fourgap`=0.0)  
*Set the parameters and the bag constant 'B0'.*
  - virtual int `calc_eq_temp_p` (`quark` &u, `quark` &d, `quark` &s, double &qq1, double &qq2, double &qq3, double &gap1, double &gap2, double &gap3, double `mu3`, double `mu8`, double &n3, double &n8, `thermo` &qb, double `temper`)  
*Calculate the EOS.*
  - virtual int `test_derivatives` (double `lmom`, double `mu3`, double `mu8`, `test_mgr` &t)  
*Check the derivatives specified by `eigenvalues()`*
  - virtual int `eigenvalues` (double `lmom`, double `mu3`, double `mu8`, double `egv`[36], double `dedmuu`[36], double `dedmud`[36], double `dedmus`[36], double `dedmu`[36], double `dedmd`[36], double `dedms`[36], double `dedu`[36], double `dedd`[36], double `deds`[36], double `dedmu3`[36], double `dedmu8`[36])  
*Calculate the energy eigenvalues as a function of the momentum.*
  - int `set_quartic` (`quartic_real_coeff` &q)  
*Set the routine for solving quartics.*
  - int `test_integration` (`test_mgr` &t)  
*Test the integration routines.*
  - int `test_normal_eigenvalues` (`test_mgr` &t)  
*Test the routine to compute the eigenvalues of non-superfluid fermions.*
  - int `test_gapped_eigenvalues` (`test_mgr` &t)  
*Test the routine to compute the eigenvalues of superfluid fermions.*
  - virtual const char \* `type` ()  
*Return string denoting type ("cfl\_njl\_eos")*
-

## Data Fields

- double [eq\\_limit](#)  
*The equal mass threshold.*
- bool [integ\\_test](#)  
*Set to true to test the integration (default false)*
- **cern\_quartic\_real\_coeff** [def\\_quartic](#)  
*The default quartic routine.*
- double [gap\\_limit](#)  
*Smallest allowable gap (default 0.0)*
- bool [zerot](#)  
*If this is true, then finite temperature corrections are ignored (default false)*
- bool [fixed\\_mass](#)  
*Use a fixed quark mass and ignore the quark condensates.*
- bool [color\\_neut](#)  
*If true, then ensure color neutrality.*
- double [GD](#)  
*Diquark coupling constant (default 3 G/4)*
- double [inte\\_epsabs](#)  
*The absolute precision for the integration (default  $10^{-4}$ )*
- double [inte\\_epsrel](#)  
*The relative precision for the integration (default  $10^{-4}$ )*
- size\_t [inte\\_npoints](#)  
*The number of points used in the last integration (default 0)*

## Protected Member Functions

- virtual int [integrands](#) (double p, double res[])  
*The integrands.*
- int [normal\\_eigenvalues](#) (double m, double lmom, double mu, double lam[2], double dldmu[2], double dldm[2])  
*Compute ungapped eigenvalues and the appropriate derivatives.*
- int [gapped\\_eigenvalues](#) (double m1, double m2, double lmom, double mu1, double mu2, double tdelta, double lam[4], double dldmu1[4], double dldmu2[4], double dldm1[4], double dldm2[4], double dldg[4])  
*Treat the simply gapped quarks in all cases gracefully.*

## For the integration

- double [rescale\\_error](#) (double err, double result\_abs, double result\_asc)  
*The error scaling function for integ\_err.*
- int [integ\\_err](#) (double a, double b, const size\_t nr, **ovector** &res, double &err2)  
*A new version of [gsl\\_inte\\_qng](#) to integrate several functions at the same time.*

## Protected Attributes

- double [temper](#)  
*Temperature.*
- double [smu3](#)  
*3rd gluon chemical potential*
- double [smu8](#)  
*8th gluon chemical potential*

## Numerical methods

- **quartic\_real\_coeff** \* [quartic](#)  
*The routine to solve quartics.*

For computing eigenvalues

- **omatrix\_cx** [iprop](#)  
*Inverse propagator matrix.*
- **omatrix\_cx** [eivec](#)  
*The eigenvectors.*
- **omatrix\_cx** [dipdgapu](#)  
*The derivative of the inverse propagator wrt the ds gap.*
- **omatrix\_cx** [dipdgapd](#)  
*The derivative of the inverse propagator wrt the us gap.*
- **omatrix\_cx** [dipdgaps](#)  
*The derivative of the inverse propagator wrt the ud gap.*
- **ovector** [eval](#)  
*The eigenvalues.*
- **gsl\_eigen\_hermv\_workspace** \* **w**  
*Workspace for eigenvalue computation.*

#### Private Member Functions

- **cfl\_njl\_eos** (const [cfl\\_njl\\_eos](#) &)
- **cfl\_njl\_eos** & **operator=** (const [cfl\\_njl\\_eos](#) &)

#### 14.6.2 Member Function Documentation

**14.6.2.1** `virtual int cfl_njl_eos::set_parameters ( double lambda = 0.0, double fourferm = 0.0, double sixferm = 0.0, double fourgap = 0.0 )`  
[virtual]

This function allows the user to specify the momentum cutoff, *lambda*, the four-fermion coupling *fourferm*, the six-fermion coupling from the 't Hooft interaction *sixferm*, and the color-superconducting coupling, *fourgap*. If 0.0 is given for any of the values, then the default is used ( $\Lambda = 602.3/(\hbar c)$ ,  $G = 1.835/\Lambda^2$ ,  $K = 12.36/\Lambda^5$ ).

If the four-fermion coupling that produces a gap is not specified, it is automatically set to 3/4 G, which is the value obtained from the Fierz transformation.

The value of the shift in the bag constant [nambujl\\_eos::B0](#) is automatically calculated to ensure that the vacuum has zero energy density and zero pressure. The functions [set\\_quarks\(\)](#) and [set\\_thermo\(\)](#) must be used before hand to specify the **quark** and **thermo** objects.

**14.6.2.2** `virtual int cfl_njl_eos::calc_eq_temp_p ( quark & u, quark & d, quark & s, double & qq1, double & qq2, double & qq3, double & gap1, double & gap2, double & gap3, double mu3, double mu8, double & n3, double & n8, thermo & qb, double temper )` [virtual]

Calculate the EOS from the quark condensates in *u*.*qq*, *d*.*qq* and *s*.*qq*. Return the mass gap equations in *qq1*, *qq2*, *qq3*, and the normal gap equations in *gap1*, *gap2*, and *gap3*.

Using *fromqq*=false as in [nambujl\\_eos](#) and [nambujl\\_eos](#) does not work here and will return an error. Also, the quarks must be set through [quark\\_eos::quark\\_set\(\)](#) before use.

If all of the gaps are less than *gap\_limit*, then the [nambujl\\_eos::calc\\_temp\\_p\(\)](#) is used, and *gap1*, *gap2*, and *gap3* are set to equal *u.del*, *d.del*, and *s.del*, respectively.

**Todo** It surprises me that *n3* is not -res[11]. Is there a sign error in the color densities?

Reimplemented in [cfl6\\_eos](#).

**14.6.2.3** `virtual int cfl_njl_eos::eigenvalues ( double lmom, double mu3, double mu8, double egv[36], double dedmuu[36], double dedmud[36], double dedmus[36], double dedmu[36], double dedmd[36], double dedms[36], double dedu[36], double dedd[36], double deds[36], double dedmu3[36], double dedmu8[36] )` [virtual]

Given the momentum *mom*, and the chemical potentials associated with the third and eighth gluons (*mu3* and *mu8*), the energy eigenvalues are computed in *egv*[0] ... *egv*[35].

#### 14.6.2.4 virtual int cfl\_njl\_eos::integrands ( double p, double res[] ) [protected, virtual]

- res[0] is the thermodynamic potential,  $\Omega$
- res[1] is  $d - \Omega/dT$
- res[2] is  $d\Omega/d\mu_u$
- res[3] is  $d\Omega/d\mu_d$
- res[4] is  $d\Omega/d\mu_s$
- res[5] is  $d\Omega/dm_u$
- res[6] is  $d\Omega/dm_d$
- res[7] is  $d\Omega/dm_s$
- res[8] is  $d\Omega/d\Delta_{ds}$
- res[9] is  $d\Omega/d\Delta_{us}$
- res[10] is  $d\Omega/d\Delta_{ud}$
- res[11] is  $d\Omega/d\mu_3$
- res[12] is  $d\Omega/d\mu_8$

Reimplemented in [cfl6\\_eos](#).

#### 14.6.2.5 int cfl\_njl\_eos::gapped\_eigenvalues ( double m1, double m2, double lmom, double mu1, double mu2, double tdelta, double lam[4], double dldmu1[4], double dldmu2[4], double dldm1[4], double dldm2[4], double dldg[4] ) [protected]

This function uses the quarks q1 and q2 to construct the eigenvalues of the inverse propagator, properly handling the either zero or finite quark mass and either zero or finite quark gaps. In the case of finite quark mass and finite quark gaps, the quartic solver is used.

The chemical potentials are separated so we can add the color chemical potentials to the quark chemical potentials if necessary.

This function is used by [eigenvalues\(\)](#). It does not work for the "ur-dg-sb" set of quarks which are paired in a non-trivial way.

**Todo** In the code, the equal mass case seems to be commented out. Why?

### 14.6.3 Field Documentation

#### 14.6.3.1 cern\_quartic\_real\_coeff cfl\_njl\_eos::def\_quartic

Slightly better accuracy (with slower execution times) can be achieved using [gsl\\_poly\\_real\\_coeff](#) which polishes the roots of the quartics. For example

```
cfl_njl_eos cfl;
gsl_poly_real_coeff gp;
cfl.set_quartic(gp);
```

Definition at line 313 of file [cfl\\_njl\\_eos.h](#).

#### 14.6.3.2 double cfl\_njl\_eos::gap\_limit

If any of the gaps are below this value, then it is assumed that they are zero and the equation of state is simplified accordingly. If all of the gaps are less than gap\_limit, then the results from [nambu\\_njl\\_eos](#) are used in [calc\\_eq\\_temp\\_p\(\)](#), [calc\\_temp\\_p\(\)](#) and [thd\\_potential\(\)](#).

Definition at line 334 of file [cfl\\_njl\\_eos.h](#).

**14.6.3.3 bool cfl\_njl\_eos::zerot**

This implements some simplifications in the momentum integration that are not possible at finite temperature.

Definition at line 342 of file cfl\_njl\_eos.h.

**14.6.3.4 double cfl\_njl\_eos::GD**

The default value is the one derived from a Fierz transformation. ([Buballa04](#))

Definition at line 357 of file cfl\_njl\_eos.h.

**14.6.3.5 double cfl\_njl\_eos::inte\_epsabs**

This is analogous to gsl\_inte::epsabs

Definition at line 364 of file cfl\_njl\_eos.h.

**14.6.3.6 double cfl\_njl\_eos::inte\_epsrel**

This is analogous to gsl\_inte::epsrel

Definition at line 371 of file cfl\_njl\_eos.h.

**14.6.3.7 size\_t cfl\_njl\_eos::inte\_npoints**

This returns 21, 43, or 87 depending on the number of function evaluations needed to obtain the desired precision. If it the routine fails to obtain the desired precision, then this variable is set to 88.

Definition at line 381 of file cfl\_njl\_eos.h.

The documentation for this class was generated from the following file:

- cfl\_njl\_eos.h

**14.7 cold\_nstar Class Reference**

Naive static cold neutron star.

```
#include <cold_nstar.h>
```

**14.7.1 Detailed Description**

This uses [hadronic\\_eos::calc\\_e\(\)](#) to compute the equation of state of zero-temperature beta-equilibrated neutron star matter and [tov\\_solve::mvsr\(\)](#) to compute the mass versus radius curve.

The electron and muon are given masses **o2scl\_fm::mass\_electron** and **o2scl\_fm::mass\_muon**, respectively.

There is an example for the usage of this class given in `examples/ex_cold_nstar.cpp`.

**EOS Output**

The function [calc\\_eos\(\)](#) generates an object of type **table\_units**, which contains the following columns

- `ed` in units of  $1/\text{fm}^4$ , the total energy density of neutron star matter
- `pr` in units of  $1/\text{fm}^4$ , the total pressure of neutron star matter
- `nb` in units of  $1/\text{fm}^3$ , the baryon number density
- `mun` in units of  $1/\text{fm}$ , the neutron chemical potential

- `mup` in units of 1/fm, the proton chemical potential
- `mue` in units of 1/fm, the electron chemical potential
- `nn` in units of 1/fm<sup>3</sup>, the neutron number density
- `np` in units of 1/fm<sup>3</sup>, the proton number density
- `ne` in units of 1/fm<sup>3</sup>, the electron number density
- `kfn` in units of 1/fm, the neutron Fermi momentum
- `kfp` in units of 1/fm, the proton Fermi momentum
- `kfe` in units of 1/fm, the electron Fermi momentum.

If `include_muons` is true, the table has additional columns

- `mumu` in units of 1/fm, the muon chemical potential
- `nmu` in units of 1/fm<sup>3</sup>, the muon number density
- `kfm` in units of 1/fm, the muon Fermi momentum

If the energy density is always positive and increasing, and the pressure is always positive and increasing, then the EOS is well-formed and `well_formed` is true. The variable `pressure_flat` records the lowest baryon density where the pressure decreases with increasing density.

After computing the equation of state, `calc_eos()` also adds the following columns

- `cs2` (unitless), the squared speed of sound
- `logp`, the logarithm of the pressure stored in `pr`
- `loge`, the logarithm of the energy density stored in `ed`
- `s` in units of 1/fm, the semi-perimeter of the Urca triangle
- `urca` in units of 1/fm<sup>4</sup>, the squared area of the Urca triangle
- `ad_index`, the adiabatic index

The condition for the direct Urca process is the area of the triangle formed by the neutron, proton, and electron Fermi momenta. Using the definition of the semi-perimeter,

$$s \equiv (k_{F,n} + k_{F,p} + k_{F,e}) / 2$$

Heron's formula gives the triangle area as

$$a = \sqrt{s(s - k_{F,n})(s - k_{F,p})(s - k_{F,e})}.$$

The column in the eos **table** labeled `urca` is  $a^2$ . If this quantity is positive, then direct Urca is allowed. The variable `allow_urca` is the smallest density for which the direct Urca process turns on, and `deny_urca` is the smallest density for which the direct Urca process turns off.

The squared speed of sound (in units of  $c$ ) is calculated by

$$c_s^2 = \frac{dP}{d\varepsilon}$$

and this is placed in the column labeled `cs2`. If the EOS is not well-formed, then this column is set to zero. If `cs2` is larger than 1, the EOS is said to be "acausal". The variables `acausal`, `acausal_ed`, and `acausal_pr` record the baryon density, energy density, and pressure where the EOS becomes acausal. The adiabatic index is calculated by

$$\Gamma = \frac{d \ln P}{d \ln \varepsilon}$$

Note that  $\Gamma$  must be greater than  $4/3$  at the center of the neutron star for stability. (This is a necessary, but not sufficient condition.) If the EOS is not well-formed then this column is set to zero.

---

### TOV Output

The TOV table contains all the columns typically generated for mass versus radius tables in [tov\\_solve](#), as well as columns containing the central values of all the densities and chemical potentials, and all the other columns computed for the EOS above.

---

**Idea for Future** Warn if the EOS becomes pure neutron matter.

Definition at line 159 of file cold\_nstar.h.

### Public Member Functions

#### Basic operation

- int [set\\_eos](#) ([hadronic\\_eos](#) &he)  
*Set the equation of state.*
- int [calc\\_eos](#) (double np\_0=0.0)  
*Calculate the given equation of state.*
- double [calc\\_urca](#) (double np\_0=0.0)  
*Compute the density at which the direct Urca process is allowed.*
- int [calc\\_nstar](#) ()  
*Calculate the M vs. R curve.*

### Data Fields

#### Default objects

- **fermion** [def\\_n](#)  
*The default neutron.*
- **fermion** [def\\_p](#)  
*The default proton.*
- fermion\_zerot [fzt](#)  
*Desc.*
- [tov\\_solve](#) [def\\_tov](#)  
*The default TOV equation solver.*
- **cern\_mroot\_root** < **funct** > [def\\_root](#)  
*The default equation solver for the EOS.*
- [tov\\_interp\\_eos](#) [def\\_tov\\_eos](#)  
*Default EOS object for the TOV solver.*

### Protected Member Functions

- double [solve\\_fun](#) (double x)  
*Solve to ensure zero charge in  $\beta$ -equilibrium.*

### Protected Attributes

- bool [eos\\_set](#)  
*True if equation of state has been set.*
  - **fermion** [e](#)  
*The electron.*
  - **fermion** [mu](#)  
*The muon.*
  - [hadronic\\_eos](#) \* [hep](#)
-



- *A pointer to the equation of state.*
- **fermion** \* [np](#)  
*A pointer to the neutron.*
- **fermion** \* [pp](#)  
*A pointer to the proton.*
- **tov\_solve** \* [tp](#)  
*A pointer to the TOV object.*
- **root**< **funct** > \* [rp](#)  
*A pointer to the solver.*
- **o2\_shared\_ptr**< **table\_units** >::type [eost](#)  
*Storage for the EOS table.*
- double [barn](#)  
*The baryon density.*

#### The thermodynamic information

- **thermo** [hb](#)
- **thermo** [h](#)
- **thermo** [l](#)

#### Output

- bool [solver\\_success](#)  
*If true, the last call of `calc_eos()` succeeded.*
- bool [well\\_formed](#)  
*If true, the energy density of the EOS is monotonically increasing and the pressure is always positive.*
- double [pressure\\_flat](#)  
*The smallest baryon density where the pressure starts to decrease.*
- double [allow\\_urca](#)  
*The smallest density where Urca becomes allowed.*
- double [deny\\_urca](#)  
*The smallest density where Urca becomes disallowed.*
- double [acausal](#)  
*The density at which the EOS becomes acausal.*
- double [acausal\\_pr](#)  
*The pressure at which the EOS becomes acausal.*
- double [acausal\\_ed](#)  
*The energy density at which the EOS becomes acausal.*
- double [solver\\_tol](#)  
*Solver tolerance (default  $10^{-4}$ )*
- int [verbose](#)  
*Verbosity parameter (default 0)*
- **o2\_shared\_ptr**< **table\_units** >::type [get\\_eos\\_results](#) ()  
*Get the eos table (after having called `calc_eos()`)*
- **o2\_shared\_ptr**< **table\_units** >::type [get\\_tov\\_results](#) ()  
*Get the results from the TOV (after having called `calc_nstar()`)*

#### Configuration

- double [nb\\_start](#)  
*The starting baryon density (default 0.05)*
- double [nb\\_end](#)  
*The final baryon density (default 2.0)*
- double [dnb](#)  
*The baryon density stepsize (default 0.01)*
- bool [include\\_muons](#)  
*If true, include muons (default false)*
- bool [err\\_nonconv](#)  
*If true, throw an exception if the calculation fails (default true)*

- int [set\\_n\\_and\\_p](#) (**fermion** &n, **fermion** &p)  
*Set the neutron and proton.*
- int [set\\_root](#) (**root**< **funct** > &rf)  
*Set the equation solver for the EOS.*
- int [set\\_tov](#) (**tov\_solve** &ts)  
*Specify the object for solving the TOV equations.*

## 14.7.2 Member Function Documentation

### 14.7.2.1 int cold\_nstar::set\_eos ( **hadronic\_eos** &he ) [inline]

This should be set before calling [calc\\_eos\(\)](#).

Definition at line 171 of file cold\_nstar.h.

### 14.7.2.2 double cold\_nstar::calc\_urca ( double np\_0=0.0 )

This is faster than using [calc\\_eos\(\)](#) since it does nothing other than computes the critical density. It does not store the equation of state.

### 14.7.2.3 int cold\_nstar::set\_n\_and\_p ( **fermion** &n, **fermion** &p ) [inline]

The default objects are of type **fermion**, with mass **o2scl\_fm::mass\_neutron** and **o2scl\_fm::mass\_proton**. These defaults will give incorrect results for non-relativistic equations of state.

Definition at line 298 of file cold\_nstar.h.

### 14.7.2.4 int cold\_nstar::set\_tov ( **tov\_solve** &ts ) [inline]

The default uses the low-density equation of state with `tov::verbose=0`. In [calc\\_nstar\(\)](#), the units are set by calling [tov\\_solve::set\\_units\(\)](#).

Definition at line 317 of file cold\_nstar.h.

## 14.7.3 Field Documentation

### 14.7.3.1 double cold\_nstar::pressure\_flat

If this is zero after calling [calc\\_eos\(\)](#), then the pressure does not decrease in the specified range of baryon density

Definition at line 211 of file cold\_nstar.h.

### 14.7.3.2 double cold\_nstar::allow\_urca

If this is zero after calling [calc\\_eos\(\)](#), then direct Urca is never allowed.

Definition at line 218 of file cold\_nstar.h.

### 14.7.3.3 double cold\_nstar::deny\_urca

If this is zero after calling [calc\\_eos\(\)](#), then direct Urca is not disallowed at a higher density than it becomes allowed.

Definition at line 226 of file cold\_nstar.h.

### 14.7.3.4 double cold\_nstar::acausal

If this is zero, then the EOS is causal at all baryon densities in the specified range

Definition at line 233 of file cold\_nstar.h.

## 14.7.3.5 double cold\_nstar::acausal\_pr

If this is zero, then the EOS is causal at all baryon densities in the specified range

Definition at line 240 of file cold\_nstar.h.

## 14.7.3.6 double cold\_nstar::acausal\_ed

If this is zero, then the EOS is causal at all baryon densities in the specified range

Definition at line 247 of file cold\_nstar.h.

The documentation for this class was generated from the following file:

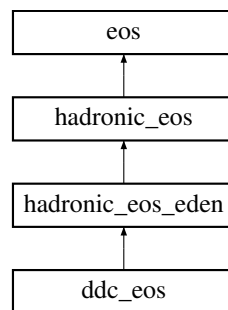
- cold\_nstar.h

## 14.8 ddc\_eos Class Reference

Relativistic mean field EOS with density dependent couplings.

```
#include <ddc_eos.h>
```

Inheritance diagram for ddc\_eos:



## 14.8.1 Detailed Description

Based on [Typel99](#).

**Idea for Future** Implement the finite temperature EOS properly.

Definition at line 46 of file ddc\_eos.h.

## Public Member Functions

- virtual int **calc\_e** (**fermion** &n, **fermion** &p, **thermo** &th)  
*Equation of state as a function of the densities.*
- virtual int **calc\_eq\_e** (**fermion** &neu, **fermion** &p, double sig, double ome, double rho, double &f1, double &f2, double &f3, **thermo** &th)  
*Equation of state and meson field equations as a function of the density.*
- virtual const char \* **type** ()  
*Return string denoting type ("ddc\_eos")*

## Data Fields

- double **rho0**

## Masses

- double **mnuc**  
*nucleon mass*
- double **ms**  
 $\phi$  mass (in fm<sup>-1</sup>)
- double **mw**  
 $A_\omega$  mass (in fm<sup>-1</sup>)
- double **mr**  
 $A_\rho$  mass (in fm<sup>-1</sup>)

## Parameters for couplings

- double **Gs**  
*The coupling  $\Gamma_\sigma(\rho_{\text{sat}})$ .*
- double **Gw**  
*The coupling  $\Gamma_\omega(\rho_{\text{sat}})$ .*
- double **Gr**  
*The coupling  $\Gamma_\rho(\rho_{\text{sat}})$ .*
- double **as**  
 $a_\sigma$
- double **aw**  
 $a_\omega$
- double **ar**  
 $a_\rho$
- double **bs**  
 $b_\sigma$
- double **bw**  
 $b_\omega$
- double **cs**  
 $c_\sigma$
- double **cw**  
 $c_\omega$
- double **ds**  
 $d_\sigma$
- double **dw**  
 $d_\omega$

## Protected Attributes

- fermion\_zerot **fzt**  
*Desc.*

## 14.8.2 Member Function Documentation

14.8.2.1 virtual int ddc\_eos::calc\_eq\_e ( fermion & neu, fermion & p, double sig, double ome, double rho, double & f1, double & f2, double & f3, thermo & th ) [virtual]

This calculates the pressure and energy density as a function of  $\mu_n, \mu_p, \phi, A_\omega, A_\rho$ . When the field equations have been solved, f1, f2, and f3 are all zero.

**Todo** Is the thermodynamic identity is satisfied even when the field equations are not solved? Check this.

The documentation for this class was generated from the following file:

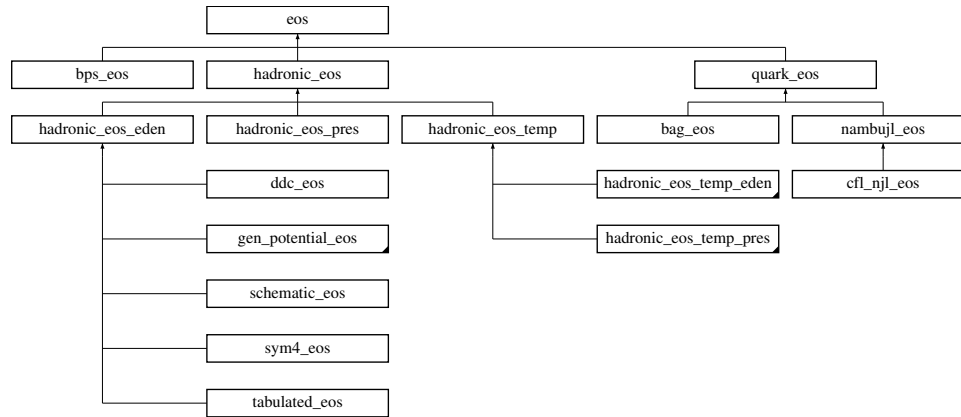
- ddc\_eos.h

## 14.9 eos Class Reference

Equation of state base.

```
#include <eos.h>
```

Inheritance diagram for eos:



### 14.9.1 Detailed Description

A base class for the computation of an equation of state

Definition at line 37 of file eos.h.

#### Public Member Functions

- virtual int [set\\_thermo](#) (**thermo** &th)  
*Set class thermo object.*
- virtual int [get\\_thermo](#) (**thermo** \*&th)  
*Get class thermo object.*
- virtual const char \* [type](#) ()  
*Return string denoting type ("eos")*

#### Data Fields

- **thermo** [def\\_thermo](#)  
*The default thermo object.*

#### Protected Attributes

- **thermo** \* [eos\\_thermo](#)  
*A pointer to the thermo object.*

The documentation for this class was generated from the following file:

- eos.h

## 14.10 ex\_apr\_eos Class Reference

Compute the APR EOS with a Gibbs construction and the mass versus radius curve [Example class].

## 14.10.1 Detailed Description

In succession, calculates nuclear matter, neutron matter, and then neutron star matter with Maxwell and Gibbs constructions.

We could use the more accurate masses in <o2scl/constants.h> here, but APR appears to have been designed to be used with  $m_n = m_p = 939$  MeV.

Definition at line 55 of file ex\_apr\_eos.cpp.

## Public Member Functions

- void **run** ()  
*Main driver, computing the APR EOS and the associated  $M$  vs.  $R$  curve.*

## Protected Member Functions

- int **maxwell\_fig7** (size\_t nv, const **ovector\_base** &x, **ovector\_base** &y)  
*Function for the Maxwell construction in Fig. 7.*
- int **mixedmaxwell** (size\_t nv, const **ovector\_base** &x, **ovector\_base** &y)  
*Maxwell construction of the nuclear matter mixed phase.*
- int **fig7fun** (size\_t nv, const **ovector\_base** &x, **ovector\_base** &y)  
*Function to construct Fig. 7.*
- int **nstar** (size\_t nv, const **ovector\_base** &x, **ovector\_base** &y)  
*Solve for neutron star matter (low-density phase)*
- int **nstar2** (size\_t nv, const **ovector\_base** &x, **ovector\_base** &y)  
*Solve for neutron star matter (high-density phase)*
- int **nstarmixed** (size\_t nv, const **ovector\_base** &x, **ovector\_base** &y)  
*Solve for neutron star matter (mixed phase)*
- void **store\_data** ()  
*Write a line of data to the table.*
- int **nucmixed** (size\_t nv, const **ovector\_base** &x, **ovector\_base** &y)  
*Solve for nuclear matter (mixed phase)*
- int **neutmixed** (size\_t nv, const **ovector\_base** &x, **ovector\_base** &y)  
*Solve for neutron matter (mixed phase)*
- int **nucleimat** (size\_t nv, const **ovector\_base** &ex, **ovector\_base** &ey)  
*Solve for phase transition to nuclei.*
- int **nucleimat\_pdrip** (size\_t nv, const **ovector\_base** &ex, **ovector\_base** &ey)  
*Solve for phase transition to nuclei with a proton drip.*

## Protected Attributes

- **fermion n**  
*Neutron for low-density phase.*
- **fermion p**  
*Proton for low-density phase.*
- **fermion n2**  
*Neutron for high-density phase.*
- **fermion p2**  
*Proton for high-density phase.*
- **fermion\_zerot fzt**  
*Compute zero-temperature thermodynamics.*
- **fermion e**  
*Electron for low-density phase.*
- **fermion mu**  
*Muon for low-density phase.*
- **fermion e2**  
*Electron for high-density phase.*
- **fermion mu2**

- **thermo hb**  
*Muon for high-density phase.*
- **thermo l**  
*Baryon thermodynamics for low-density phase.*
- **thermo hb2**  
*Leptonic thermodynamics for low-density phase.*
- **thermo tot**  
*Baryon thermodynamics for high-density phase.*
- **thermo l2**  
*Total thermodynamics.*
- **gsl\_mroot\_hybrids nd**  
*Leptonic thermodynamics for high-density phase.*
- **gsl\_mroot\_hybrids nd2**  
*Solver.*
- **gsl\_mroot\_hybrids sat\_solver**  
*Solver.*
- **double nb**  
*Solver for saturation properties.*
- **double chi**  
*Baryon density.*
- **double mub**  
*Volume fraction of low-density phase.*
- **double muq**  
*Baryon chemical potential.*
- **double f7x**  
*Charge chemical potential.*
- **int choice**  
*Proton fraction for Fig. 7.*
- **apr\_eos ap**  
*Choice of model from APR.*
- **table\_units at**  
*Base APR EOS.*
- **cern\_deriv<funct> cd**  
*Table for output.*
- **hdf\_file hf**  
*Derivative object.*
- **hdf\_file hf**  
*HDF file for output.*

#### Phase specification

- **int phase**
- **static const int low\_phase = 1**
- **static const int mixed\_phase = 2**
- **static const int high\_phase = 3**

### 14.10.2 Member Function Documentation

#### 14.10.2.1 void ex\_apr\_eos::run ( ) [inline]

Compute matter at densities below the maxwell construction

Definition at line 575 of file ex\_apr\_eos.cpp.

The documentation for this class was generated from the following file:

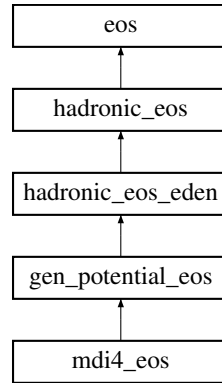
- ex\_apr\_eos.cpp

## 14.11 gen\_potential\_eos Class Reference

Generalized potential model equation of state.

```
#include <gen_potential_eos.h>
```

Inheritance diagram for gen\_potential\_eos:



## 14.11.1 Detailed Description

The single particle energy is defined by the functional derivative of the energy density with respect to the distribution function

$$e_\tau = \frac{\delta \mathcal{E}}{\delta f_\tau}$$

The effective mass is defined by

$$\frac{m^*}{m} = \left( \frac{m}{k} \frac{de_\tau}{dk} \right)^{-1}_{k=k_F}$$

In all of the models, the kinetic energy density is  $\tau_n + \tau_p$  where

$$\tau_i = \frac{2}{(2\pi)^3} \int d^3k \left( \frac{k^2}{2m} \right) f_i(k, T)$$

and the number density is

$$\rho_i = \frac{2}{(2\pi)^3} \int d^3k f_i(k, T)$$

When `form == mdi_form` or `gbd_form`, the potential energy density is given by Das03 :

$$V(\rho, \delta) = \frac{Au}{\rho_0} \rho_n \rho_p + \frac{A_l}{2\rho_0} (\rho_n^2 + \rho_p^2) + \frac{B}{\sigma + 1} \frac{\rho^{\sigma+1}}{\rho_0^\sigma} (1 - x\delta^2) + V_{mom}(\rho, \delta)$$

where  $\delta = 1 - 2\rho_p/(\rho_n + \rho_p)$ . If `form == mdi_form`, then

$$V_{mom}(\rho, \delta) = \frac{1}{\rho_0} \sum_{\tau, \tau'} C_{\tau, \tau'} \int \int d^3k d^3k' \frac{f_\tau(\vec{k}) f_{\tau'}(\vec{k}')}{1 - (\vec{k} - \vec{k}')^2 / \Lambda^2}$$

where  $C_{1/2, 1/2} = C_{-1/2, -1/2} = C_\ell$  and  $C_{1/2, -1/2} = C_{-1/2, 1/2} = C_u$ . Otherwise if `form == gbd_form`, then

$$V_{mom}(\rho, \delta) = \frac{1}{\rho_0} [C_\ell (\rho_n g_n + \rho_p g_p) + C_u (\rho_n g_p + \rho_p g_n)]$$



where

$$g_i = \frac{\Lambda^2}{\pi^2} [k_{F,i} - \Lambda \tan^{-1}(k_{F,i}/\Lambda)]$$

Otherwise, if `form == bgbd_form`, `bpalb_form` or `sl_form`, then the potential energy density is given by Bombaci01 :

$$V(\rho, \delta) = V_A + V_B + V_C$$

$$V_A = \frac{2A}{3\rho_0} \left[ \left(1 + \frac{x_0}{2}\right) \rho^2 - \left(\frac{1}{2} + x_0\right) (\rho_n^2 + \rho_p^2) \right]$$

$$V_B = \frac{4B}{3\rho_0^\sigma} \frac{T}{1 + 4B'T / (3\rho_0^{\sigma-1} \rho^2)}$$

where

$$T = \rho^{\sigma-1} \left[ \left(1 + \frac{x_3}{2}\right) \rho^2 - \left(\frac{1}{2} + x_3\right) (\rho_n^2 + \rho_p^2) \right]$$

The term  $V_C$  is:

$$V_C = \sum_{i=1}^{i_{\max}} \frac{4}{5} (C_i + 2z_i) \rho (g_{n,i} + g_{p,i}) + \frac{2}{5} (C_i - 8z_i) (\rho_n g_{n,i} + \rho_p g_{p,i})$$

where

$$g_{\tau,i} = \frac{2}{(2\pi)^3} \int d^3k f_{\tau}(k, T) g_i(k)$$

For `form == bgbd_form` or `form == bpalb_form`, the form factor is given by

$$g_i(k) = \left(1 + \frac{k^2}{\Lambda_i^2}\right)^{-1}$$

while for `form == sl_form`, the form factor is given by

$$g_i(k) = 1 - \frac{k^2}{\Lambda_i^2}$$

where  $\Lambda_i$  is specified in the parameter `Lambda` when necessary.

See Mathematica notebook at

```
doc/o2scl/extras/gen_potential_eos.nb
doc/o2scl/extras/gen_potential_eos.ps
```

**Bug** The BGBD EOS doesn't work and the effective mass for the GBD EOS doesn't work

**Idea for Future** Calculate the chemical potentials analytically

Definition at line 170 of file `gen_potential_eos.h`.

#### Public Member Functions

- virtual int `calc_e` (**fermion** &ne, **fermion** &pr, **thermo** &lt)  
*Equation of state as a function of density.*
- int `set_mu_deriv` (**deriv** < **funct** > &de)  
*Set the derivative object to calculate the chemical potentials.*
- virtual const char \* `type` ()  
*Return string denoting type ("gen\_potential\_eos")*

## Data Fields

- int [form](#)  
*Form of potential.*
- [gsl\\_deriv](#)< [funct](#) > [def\\_mu\\_deriv](#)  
*The default derivative object for calculating chemical potentials.*

## The parameters for the various interactions

- double [x](#)
- double [Au](#)
- double [Al](#)
- double [rho0](#)
- double [B](#)
- double [sigma](#)
- double [Cl](#)
- double [Cu](#)
- double [Lambda](#)
- double [A](#)
- double [x0](#)
- double [x3](#)
- double [Bp](#)
- double [C1](#)
- double [z1](#)
- double [Lambda2](#)
- double [C2](#)
- double [z2](#)
- double [bpal\\_esym](#)
- int [sym\\_index](#)

## Static Public Attributes

- static const int [mdi\\_form](#) = 1  
*The "momentum-dependent-interaction" form.*
- static const int [bgbd\\_form](#) = 2  
*The modified GBD form.*
- static const int [bpalb\\_form](#) = 3  
*The form from [Prakash88](#) as formulated in [Bombaci01](#).*
- static const int [sl\\_form](#) = 4  
*The "SL" form. See [Bombaci01](#).*
- static const int [gbd\\_form](#) = 5  
*The Gale, Bertsch, Das Gupta from [Gale87](#).*
- static const int [bpal\\_form](#) = 6  
*The form from [Prakash88](#).*

## Protected Member Functions

- double [mom\\_integral](#) (double pft, double pftp)  
*Compute the momentum integral for [mdi\\_form](#).*
- double [energy](#) (double x)  
*Compute the energy.*

## Protected Attributes

- [nonrel\\_fermion](#) [nrf](#)  
*Desc.*
- bool [mu\\_deriv\\_set](#)  
*True of the derivative object has been set.*
- [deriv](#)< [funct](#) > \* [mu\\_deriv\\_ptr](#)  
*The derivative object.*

The mode for the energy() function [protected]

- int **mode**
- static const int **nmode** = 1
- static const int **pmode** = 2
- static const int **normal** = 0

The documentation for this class was generated from the following file:

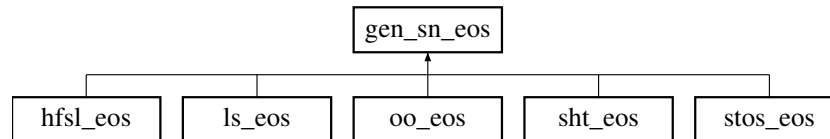
- gen\_potential\_eos.h

## 14.12 gen\_sn\_eos Class Reference

A base class for the supernova EOSs [abstract].

```
#include <gen_sn_eos.h>
```

Inheritance diagram for gen\_sn\_eos:



### 14.12.1 Detailed Description

This class is experimental.

The EOSs are stored in a set of **tensor\_grid3** objects on grids with density in  $\text{fm}^{-3}$ , electron fraction (unitless) and temperature in MeV.

Some tabulated EOSs do not store data for `gen_sn_eos::E`, `gen_sn_eos::F`, `gen_sn_eos::S`, and `gen_sn_eos::P`. In that case, the grid is set for these objects but the data is set to zero. To compute these from the data after loading the EOS table, use `gen_sn_eos::compute_eg()`.

The function `load()` loads the entire EOS into memory. Memory allocation is automatically performed by `load()`, but not deallocated until `free()` or the destructor is called.

After loading, you can interpolate the EOS by using `tensor_grid3::interp` directly. For example, the following returns the mass number at an arbitrary baryon density, electron fraction, and temperature assuming the table is stored in `skm.dat`:

```
ls_eos ls;
ls.load("skm.dat");
double nb=0.01, Ye=0.2, T=10.0;
cout << A.interp(nb,Ye,T) << endl;
```

Interpolation for all EOSs is linear by default.

**Todo** Ensure all chemical potentials are based on the same rest masses?

Allow logarithmic grids for any of nb, Ye, or T.

**Idea for Future** Create a **table** object, possibly using `tensor_grid::vector_slice`.

**Idea for Future** Show how `matrix_slice` and `vector_slice` can be used with this object.

**Idea for Future** Could this be a child of [hadronic\\_eos\\_temp](#) and then directly used in `cold_nstar()`?

**Idea for Future** Add option to load and store a separate lepton/photon EOS

**Idea for Future** Add muons and/or pions

Definition at line 84 of file `gen_sn_eos.h`.

#### Public Member Functions

- `int` [check\\_free\\_energy](#) (`test_mgr` &tm)  
*Test the free energy and store results in  $\tau m$ .*
- virtual void [beta\\_eq\\_T0](#) (size\_t i, double &nb, double &E\_beta, double &P\_beta, double &Ye\_beta, double &Z\_beta, double &A\_beta)=0  
*Compute energy per baryon of matter in beta equilibrium at zero temperature at a fixed grid point [abstract].*
- virtual void [beta\\_eq\\_sfixed](#) (size\_t i, double entr, double &nb, double &E\_beta, double &P\_beta, double &Ye\_beta, double &Z\_beta, double &A\_beta, double &T\_beta)  
*Compute properties of matter in beta equilibrium at  $s=4$ .*

#### Memory allocation

- `int` [alloc](#) ()  
*Allocate memory.*
- `int` [free](#) ()  
*Free allocated memory.*

#### Data Fields

- `int` [verbose](#)  
*Verbosity parameter (default 1)*

#### Data

- `tensor_grid3` [F](#)  
*Total free energy per baryon in MeV.*
- `tensor_grid3` [Fint](#)  
*Free energy per baryon without lepton and photon contributions in MeV.*
- `tensor_grid3` [E](#)  
*Total internal energy per baryon in MeV.*
- `tensor_grid3` [Eint](#)  
*Internal energy per baryon without lepton and photon contributions in MeV.*
- `tensor_grid3` [P](#)  
*Total pressure in MeV/fm<sup>3</sup>.*
- `tensor_grid3` [Pint](#)  
*Pressure without lepton and photon contributions in MeV/fm<sup>3</sup>.*
- `tensor_grid3` [S](#)  
*Total entropy per baryon.*
- `tensor_grid3` [Sint](#)  
*Entropy per baryon without lepton and photon contributions.*
- `tensor_grid3` [mun](#)  
*Neutron chemical potential in MeV.*
- `tensor_grid3` [mup](#)  
*Proton chemical potential in MeV.*
- `tensor_grid3` [Z](#)  
*Proton number.*
- `tensor_grid3` [A](#)  
*Mass number.*
- `tensor_grid3` [Xn](#)  
*Neutron fraction.*

- **tensor\_grid3** [Xp](#)  
*Proton fraction.*
- **tensor\_grid3** [Xalpha](#)  
*Alpha particle fraction.*
- **tensor\_grid3** [Xnuclei](#)  
*Fraction of heavy nuclei.*
- **tensor\_grid3** [other](#) [20]  
*Other data sets.*
- **tensor\_grid3** \* [arr](#) [[n\\_base](#)+20]  
*List of pointers to data.*

#### Load table

- bool [loaded](#)  
*If true, a EOS table was successfully loaded (default false)*
- bool [with\\_leptons\\_loaded](#)  
*True if thermodynamics with leptons has been loaded.*
- bool [baryons\\_only\\_loaded](#)  
*True if baryon-only thermodynamics has been loaded.*
- virtual void [load](#) (std::string fname)=0  
*Load table from filename `fname`.*

#### Grid and data sizes

- size\_t [n\\_nB](#)  
*Size of baryon density grid.*
- size\_t [n\\_Ye](#)  
*Size of electron fraction grid.*
- size\_t [n\\_T](#)  
*Size of temperature grid.*
- size\_t [n\\_oth](#)  
*Number of additional data sets.*
- static const size\_t [n\\_base](#) = 16  
*Number of base data sets.*

#### Interpolation

- **def\_interp\_mgr**< [uvector\\_base](#), [linear\\_interp](#) > [dim1](#)  
*Default interpolation object.*
- **def\_interp\_mgr** < [uvector\\_const\\_subvector](#), [linear\\_interp](#) > [dim2](#)  
*Default interpolation object.*
- int [set\\_interp](#) ([base\\_interp\\_mgr](#)< [uvector\\_base](#) > &bi1, [base\\_interp\\_mgr](#)< [uvector\\_const\\_subvector](#) > &bi2)  
*Set interpolation managers.*

#### Electron and photon contribution

- **boson** [photon](#)  
*Photon.*
- **fermion** [electron](#)  
*Electron;.*
- **rel\_fermion** [relf](#)  
*Desc.*
- **eff\_boson** [effb](#)  
*Desc.*
- int [compute\\_eg](#) ()  
*Compute the electron and photon contribution for the full grid.*

## 14.12.2 Member Function Documentation

14.12.2.1 `int gen_sn_eos::compute_eg ( )`

This function computes the data for `E`, `P`, `S`, and `F` by adding electrons and photons to the baryon contributions stored in `Eint`, `Pint`, `Sint`, and `Fint`.

The electron contribution to the internal energy and free energy computed by this function includes the electron rest mass.

14.12.2.2 `int gen_sn_eos::check_free_energy ( test_mgr & tm )`

This checks that the data in `Fint` is consistent with that in `Eint` and `Sint`.

14.12.2.3 `virtual void gen_sn_eos::beta_eq_T0 ( size_t i, double & nb, double & E_beta, double & P_beta, double & Ye_beta, double & Z_beta, double & A_beta ) [pure virtual]`

Given an index `i` for the baryon grid, between 0 and `n_nB` (inclusive), this computes the properties of matter in beta equilibrium at zero temperature by finding the electron fraction grid point which minimizes `E`. The baryon density is returned in `nb`, the energy per baryon in `E_beta`, the pressure in `P_beta`, the electron fraction in `Ye_beta`, the proton number in `Z_beta` and the mass number of the nucleus in `A_beta`.

Some tables explicitly contain zero-temperature data which is used when it is available. Otherwise, linear interpolation is used to extrapolate down to zero from the lowest temperature grid points.

Implemented in `hfs_l_eos`, `sht_eos`, `stos_eos`, `oo_eos`, and `ls_eos`.

The documentation for this class was generated from the following file:

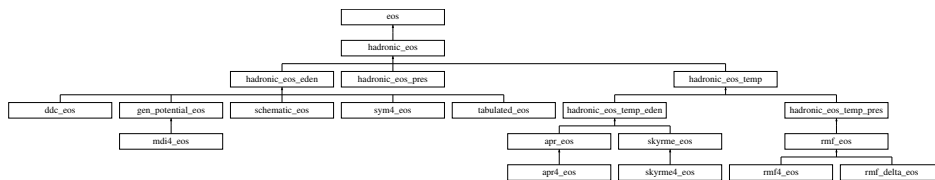
- `gen_sn_eos.h`

## 14.13 hadronic\_eos Class Reference

Hadronic equation of state [abstract base].

```
#include <hadronic_eos.h>
```

Inheritance diagram for `hadronic_eos`:



## 14.13.1 Detailed Description

Denote the number density of neutrons as  $n_n$ , the number density of protons as  $n_p$ , the total baryon density  $n_B = n_n + n_p$ , the asymmetry  $\alpha \equiv (n_n - n_p)/n_B$ , the nuclear saturation density as  $n_0 \approx 0.16 \text{ fm}^{-3}$ , and the quantity  $\eta \equiv (n - n_0)/3n_0$ . Then the energy per baryon of nucleonic matter can be written as an expansion around  $\epsilon = \alpha = 0$

$$E(n_B, \alpha) = -B + \frac{\tilde{K}}{2!} \epsilon^2 + \frac{Q_0}{3!} \epsilon^3 + \alpha^2 \left( S + L\epsilon + \frac{K_{\text{sym}}}{2!} \epsilon^2 + \frac{Q_{\text{sym}}}{3!} \epsilon^3 \right) + E_4(n_B, \alpha) + \mathcal{O}(\alpha^6) \quad (\text{Eq. 1})$$

where  $E_4$  represents the quartic terms

$$E_4(n_B, \alpha) = \alpha^4 \left( S_4 + L_4 \epsilon + \frac{K_4}{2!} \epsilon^2 + \frac{Q_4}{3!} \epsilon^3 \right) \quad (\text{Eq. 2})$$

(Adapted slightly from [Piekarewicz09](#)). From this, one can compute the energy density of nuclear matter  $\varepsilon(n_B, \alpha) = n_B E(n_B, \alpha)$ , the chemical potentials  $\mu_i \equiv (\partial \varepsilon) / (\partial n_i)$  and the pressure  $P = -\varepsilon + \mu_n n_n + \mu_p n_p$ . This expansion motivates the definition of several separate terms. The binding energy  $B$  of symmetric nuclear matter ( $\alpha = 0$ ) is around 16 MeV.

The compression modulus is usually defined by  $\chi = -1/V(dV/dP) = 1/n(dP/dn)^{-1}$ . In nuclear physics it has become common to use the incompressibility (or bulk) modulus with an extra factor of 9,  $K = 9/(n\chi)$  and refer to  $K$  simply as the incompressibility. Here, we define the function

$$K(n_B, \alpha) \equiv 9 \left( \frac{\partial P}{\partial n_B} \right) = 9 n_B \left( \frac{\partial^2 \varepsilon}{\partial n_B^2} \right)$$

This quantity is computed by the function [fcomp\(\)](#) by computing the first derivative of the pressure, which is more numerically stable than the second derivative of the energy density (and most O2scl EOSs compute the pressure exactly). This function is typically evaluated at the point ( $n_B = n_0, \alpha = 0$ ) and is stored in [comp](#). This quantity is not always the same as  $\tilde{K}$ , defined here as

$$\tilde{K}(n_B, \alpha) = 9 n_B^2 \left( \frac{\partial^2 E}{\partial n_B^2} \right) = K(n_B, \alpha) - \frac{1}{n_B} 18 P(n_B, \alpha)$$

We denote  $K \equiv K(n_B = n_0, \alpha = 0)$  and similarly for  $\tilde{K}$ , the quantity in Eq. 1 above. In nuclear matter at saturation, the pressure is zero and  $K = \tilde{K}$ . See [Chabanat97](#) for a discussion of this distinction.

The symmetry energy  $S(n_B, \alpha)$  can be defined as

$$S(n_B, \alpha) \equiv \frac{1}{2 n_B} \frac{\partial^2 \varepsilon}{\partial \alpha^2}$$

and the parameter  $S$  in Eq. 1 is just  $S(n_0, 0)$ . Using

$$\left( \frac{\partial \varepsilon}{\partial \alpha} \right)_{n_B} = \frac{\partial \varepsilon}{\partial n_n} \left( \frac{\partial n_n}{\partial \alpha} \right)_{n_B} + \frac{\partial \varepsilon}{\partial n_p} \left( \frac{\partial n_p}{\partial \alpha} \right)_{n_B} = \frac{n_B}{2} (\mu_n - \mu_p)$$

this can be rewritten

$$S(n_B, \alpha) = \frac{1}{4} \frac{\partial}{\partial \alpha} (\mu_n - \mu_p)$$

where the dependence of the chemical potentials on  $n_B$  and  $\alpha$  is not written explicitly. This quantity is computed by function [fesym\(\)](#). Note that many of the functions in this class are written in terms of the proton fraction  $x_p = (1 - \alpha)/2$  denoted as 'p f' instead of as functions of  $\alpha$ . Frequently,  $S(n_B, \alpha)$  is evaluated at  $\alpha = 0$  to give a univariate function of the baryon density. It is sometimes also evaluated at the point ( $n_B = n_0, \alpha = 0$ ), and this value is denoted by  $S$  above and is typically stored in [esym](#). Alternatively, one can define the symmetry energy by

$$\tilde{S}(n_B) \approx E(n_B, \alpha = 1) - E(n_B, \alpha = 0)$$

which is computed by function [fesym\\_diff\(\)](#). The functions  $S(n_B, \alpha = 0)$  and  $\tilde{S}(n_B)$  are equal when  $\mathcal{O}(\alpha^4)$  terms are zero. In this case,  $\mu_n - \mu_p$  is proportional to  $\alpha$  and so

$$S(n_B) = \tilde{S}(n_B) = \frac{1}{4} \frac{(\mu_n - \mu_p)}{\alpha}.$$

The symmetry energy slope parameter  $L$ , can be defined by

$$L(n_B, \alpha) \equiv 3 n_B \frac{\partial S(n_B, \alpha)}{\partial n_B} = 3 n_B \frac{\partial}{\partial n_B} \left[ \frac{1}{2 n_B} \frac{\partial^2 \varepsilon}{\partial \alpha^2} \right]$$

This can be rewritten as

$$L(n_B, \alpha) = \frac{3 n_B}{4} \frac{\partial}{\partial n_B} \frac{\partial}{\partial \alpha} (\mu_n - \mu_p)$$

(where the derivatives can be evaluated in either order) or alternatively using

$$\left( \frac{\partial \varepsilon}{\partial n_B} \right)_{\alpha} = \frac{\partial \varepsilon}{\partial n_n} \left( \frac{\partial n_n}{\partial n_B} \right)_{\alpha} + \frac{\partial \varepsilon}{\partial n_p} \left( \frac{\partial n_p}{\partial n_B} \right)_{\alpha} = \frac{1}{2} (\mu_n + \mu_p)$$

$L$  can be rewritten

$$\begin{aligned} L(n_B, \alpha) &= 3 n_B \left[ \frac{-1}{2 n_B^2} \frac{\partial^2 \varepsilon}{\partial \alpha^2} + \frac{1}{4 n_B} \frac{\partial^2}{\partial \alpha^2} (\mu_n + \mu_p) \right] \\ &= \frac{3}{4} \frac{\partial^2}{\partial \alpha^2} (\mu_n + \mu_p) - 3 S(n_B, \alpha). \end{aligned}$$

The third derivative with respect to the baryon density is sometimes called the skewness. Here, we define

$$Q_0(n_B, \alpha) = 27n_B^3 \frac{\partial^3}{\partial n_B^3} \left( \frac{\varepsilon}{n_B} \right) = 27n_B^3 \frac{\partial^2}{\partial n_B^2} \left( \frac{P}{n_B^2} \right)$$

and this function is computed in `fkprime()`.

The second derivative of the symmetry energy with respect to the baryon density is

$$K_{\text{sym}}(n_B, \alpha) = 9n_B^2 \frac{\partial^2}{\partial n_B^2} S(n_B, \alpha)$$

The third derivative of the symmetry energy with respect to the baryon density is

$$Q_{\text{sym}}(n_B, \alpha) = 27n_B^3 \frac{\partial^3}{\partial n_B^3} S(n_B, \alpha)$$

Note that solving for the baryon density for which  $P = 0$  gives, to order  $\alpha^2$  (Piekarewicz09)

$$n_B = n_0 \left[ 1 + \frac{6K}{Q} + \alpha^2 \left( \frac{3L}{K} - \frac{6K_{\text{sym}}}{Q} + \frac{6KQ_{\text{sym}}}{Q^2} \right) \right]$$

The quartic symmetry energy  $S_4(n_B, \alpha)$  can be defined as

$$S_4(n_B, \alpha) \equiv \frac{1}{24n_B} \frac{\partial^4 \varepsilon}{\partial \alpha^4}$$

However, fourth derivatives are difficult numerically, and so an alternative quantity is preferable. Instead, one can evaluate the extent to which  $\mathcal{O}(\alpha^4)$  terms are important from

$$\eta(n_B) \equiv \frac{E(n_B, 1) - E(n_B, 1/2)}{3[E(n_B, 1/2) - E(n_B, 0)]}$$

as described in Steiner06. This function can be expressed either in terms of  $\tilde{S}$  or  $S_4$

$$\eta(n_B) = \frac{5\tilde{S}(n_B) - S(n_B, 0)}{\tilde{S}(n_B) + 3S(n_B, 0)} = \frac{5S_4(n_B, 0) + 4S(n_B, 0)}{S_4(n_B, 0) + 4S(n_B, 0)}$$

Evaluating this function at the saturation density gives

$$\eta(n_0) = \frac{4S + 5S_4}{4S + S_4}$$

(Note that  $S_4$  is referred to as  $Q$  in Steiner06). Sometimes it is useful to separate out the kinetic and potential parts of the energy density when computing  $\eta(n_B)$ , and the class `sym4_eos_base` is useful for this purpose.

**Idea for Future** Could write a function to compute the "symmetry free energy" or the "symmetry entropy"

Definition at line 248 of file `hadronic_eos.h`.

#### Public Member Functions

- `int gradient_qij (fermion &n, fermion &p, thermo &th, double &qnn, double &qnp, double &qpp, double &dqnnndnn, double &dqnnndnp, double &dqnpdnn, double &dqnpdnp, double &dqppdnn, double &dqppdnp)`  
Calculate coefficients for gradient part of Hamiltonian.
- `virtual const char * type ()`  
Return string denoting type ("hadronic\_eos")



## Equation of state

- virtual int **calc\_p** (**fermion** &n, **fermion** &p, **thermo** &th)=0  
*Equation of state as a function of the chemical potentials.*
- virtual int **calc\_e** (**fermion** &n, **fermion** &p, **thermo** &th)=0  
*Equation of state as a function of density.*

## EOS properties

- virtual double **fcomp** (double nb, const double &alpha=0.0)  
*Calculate the incompressibility in fm<sup>-1</sup> using **calc\_e()***
- virtual double **feoa** (double nb, const double &alpha=0.0)  
*Calculate the energy per baryon in fm<sup>-1</sup> using **calc\_e()***
- virtual double **fesym** (double nb, const double &alpha=0.0)  
*Calculate symmetry energy of matter in fm<sup>-1</sup> using **calc\_dmu\_alpha()**.*
- virtual double **fesym\_err** (double nb, double &alpha, double &unc)  
*Calculate symmetry energy of matter and its uncertainty.*
- virtual double **fesym\_slope** (double nb, const double &alpha=0.0)  
*The symmetry energy slope parameter.*
- virtual double **fesym\_curve** (double nb, const double &alpha=0.0)  
*The curvature of the symmetry energy.*
- virtual double **fesym\_skew** (double nb, const double &alpha=0.0)  
*The skewness of the symmetry energy.*
- virtual double **fesym\_diff** (double nb)  
*Calculate symmetry energy of matter as energy of neutron matter minus the energy of nuclear matter.*
- virtual double **feta** (double nb)  
*The strength parameter for quartic terms in the symmetry energy.*
- virtual double **fkprime** (double nb, const double &alpha=0.0)  
*Calculate skewness of nuclear matter using **calc\_e()***
- virtual double **fmsom** (double nb, const double &alpha=0.0)  
*Calculate reduced neutron effective mass using **calc\_e()***
- virtual double **fn0** (double alpha, double &leoa)  
*Calculate saturation density using **calc\_e()***
- virtual int **saturation** ()  
*Calculates some of the EOS properties at the saturation density.*

## Functions for calculating physical properties

- double **calc\_dmu\_alpha** (double alpha, const double &nb)  
*Compute the difference between neutron and proton chemical potentials as a function of the isospin asymmetry.*
- double **calc\_musum\_alpha** (double alpha, const double &nb)  
*Compute the sum of the neutron and proton chemical potentials as a function of the isospin asymmetry.*
- double **calc\_pressure\_nb** (double nb, const double &alpha=0.0)  
*Compute the pressure as a function of baryon density at fixed isospin asymmetry.*
- double **calc\_edensity\_nb** (double nb, const double &alpha=0.0)  
*Compute the energy density as a function of baryon density at fixed isospin asymmetry.*
- void **const\_pf\_derivs** (double nb, double pf, double &dednb\_pf, double &dPdnb\_pf)  
*Compute derivatives at constant proton fraction.*
- double **calc\_press\_over\_den2** (double nb, const double &alpha=0.0)  
*Calculate pressure / baryon density squared in nuclear matter as a function of baryon density at fixed isospin asymmetry.*
- double **calc\_edensity\_alpha** (double alpha, const double &nb)  
*Calculate energy density as a function of the isospin asymmetry at fixed baryon density.*

## Other functions

- int **nuc\_matter\_p** (size\_t nv, const **ovector\_base** &x, **ovector\_base** &y, double \*&pa)  
*Nucleonic matter from **calc\_p()***
- int **nuc\_matter\_e** (size\_t nv, const **ovector\_base** &x, **ovector\_base** &y, double \*&pa)  
*Nucleonic matter from **calc\_e()***

## Set auxiliary objects

- virtual int **set\_mroot** (**mroot**< **mm\_funct**<> > &mr)  
*Set class mroot object for use in calculating chemical potentials from densities.*
- virtual int **set\_sat\_root** (**root**< **funct** > &mr)

- *Set class mroot object for use calculating saturation density.*  
virtual int [set\\_sat\\_deriv](#) (**deriv**< **funct** > &de)  
*Set **deriv** object to use to find saturation properties.*
- virtual int [set\\_sat\\_deriv2](#) (**deriv**< **funct** > &de)  
*Set the second **deriv** object to use to find saturation properties.*
- virtual int [set\\_n\\_and\\_p](#) (**fermion** &n, **fermion** &p)  
*Set neutron and proton.*

#### Data Fields

- double [eoa](#)  
*Binding energy.*
- double [comp](#)  
*Compressibility.*
- double [esym](#)  
*Symmetry energy.*
- double [n0](#)  
*Saturation density.*
- double [msom](#)  
*Effective mass (neutron)*
- double [kprime](#)  
*Skewness.*
- **fermion** [def\\_neutron](#)  
*The default neutron.*
- **fermion** [def\\_proton](#)  
*The default proton.*

#### Default solvers and derivative classes

- **gsl\_deriv**< **funct** > [def\\_deriv](#)  
*The default object for derivatives.*
- **gsl\_deriv**< **funct** > [def\\_deriv2](#)  
*The second default object for derivatives.*
- **gsl\_mroot\_hybrids**< **mm\_funct**<> > [def\\_mroot](#)  
*The default solver.*
- **cern\_mroot\_root**< **funct** > [def\\_sat\\_root](#)  
*The default solver for calculating the saturation density.*

#### Protected Member Functions

- double [t1\\_fun](#) (double barn)  
*Compute t1 for [gradient\\_qij\(\)](#).*
- double [t2\\_fun](#) (double barn)  
*Compute t2 for [gradient\\_qij\(\)](#).*

#### Protected Attributes

- **mroot**< **mm\_funct**<> > \* [eos\\_mroot](#)  
*The EOS solver.*
- **root**< **funct** > \* [sat\\_root](#)  
*The solver to compute saturation properties.*
- **deriv**< **funct** > \* [sat\\_deriv](#)  
*The derivative object for saturation properties.*
- **deriv**< **funct** > \* [sat\\_deriv2](#)  
*The second derivative object for saturation properties.*
- **fermion** \* [neutron](#)  
*The neutron object.*
- **fermion** \* [proton](#)  
*The proton object.*

## 14.13.2 Member Function Documentation

14.13.2.1 virtual double hadronic\_eos::fcomp ( double *nb*, const double & *alpha* = 0.0 ) [virtual]

This function computes  $K(n_B, \alpha) = 9n_B \partial^2 \epsilon / (\partial n_B^2) = 9 \partial P / (\partial n_B)$ . The value of  $K(n_0, 0)$ , often referred to as the "compressibility", is stored in [comp](#) by [saturation\(\)](#) and is about 240 MeV at saturation density.

14.13.2.2 virtual double hadronic\_eos::feoa ( double *nb*, const double & *alpha* = 0.0 ) [virtual]

This function computes the energy per baryon of matter without the nucleon rest masses at the specified baryon density, *nb*, and isospin asymmetry *alpha*.

14.13.2.3 virtual double hadronic\_eos::fesym ( double *nb*, const double & *alpha* = 0.0 ) [virtual]

This function computes the symmetry energy,

$$\left( \frac{1}{2n_B} \frac{d^2 \epsilon}{d\alpha^2} \right) = \frac{1}{4} \frac{\partial}{\partial \alpha} (\mu_n - \mu_p)$$

at the value of  $n_B$  given in *nb* and  $\alpha$  given in *alpha*. The symmetry energy at  $\alpha = 0$  at the saturation density and is stored in [esym](#) by [saturation\(\)](#).

14.13.2.4 virtual double hadronic\_eos::fesym\_err ( double *nb*, double & *alpha*, double & *unc* ) [virtual]

This estimates the uncertainty due to the numerical differentiation, assuming that difference between the neutron and proton chemical potentials is computed exactly by [calc\\_dmu\\_alpha\(\)](#).

14.13.2.5 virtual double hadronic\_eos::fesym\_slope ( double *nb*, const double & *alpha* = 0.0 ) [virtual]

This returns the value of the "slope parameter" of the symmetry energy

$$L = 3n_B \left( \frac{\partial E_{sym}}{\partial n_B} \right)$$

in inverse Fermis.

where  $n_B$  is the baryon density. This ranges between about zero and 200 MeV for many EOSs.

14.13.2.6 virtual double hadronic\_eos::fesym\_diff ( double *nb* ) [virtual]

This function returns the energy per baryon of neutron matter minus the energy per baryon of nuclear matter. This will deviate significantly from the results from [fesym\(\)](#) only if the dependence of the symmetry energy on  $\delta$  is not quadratic.

Reimplemented in [apr\\_eos](#).

14.13.2.7 virtual double hadronic\_eos::fkprime ( double *nb*, const double & *alpha* = 0.0 ) [virtual]

The skewness is defined to be  $27n^3 d^3(\epsilon/n)/(dn^3) = 27n^3 d^2(P/n^2)/(dn^2)$  and is denoted 'kprime'. This definition seems to be ambiguous for densities other than the saturation density and is not quite analogous to the compressibility.

14.13.2.8 virtual double hadronic\_eos::fmsom ( double *nb*, const double & *alpha* = 0.0 ) [virtual]

Neutron effective mass (as stored in [part::ms](#)) divided by vacuum mass (as stored in [part::m](#)) in nuclear matter at saturation density. Note that this simply uses the value of *n.ms* from [calc\\_e\(\)](#), so that this effective mass could be either the Landau or Dirac mass depending on the context. Note that this may not be equal to the reduced proton effective mass.

14.13.2.9 virtual double hadronic\_eos::fn0 ( double *alpha*, double & *leoa* ) [virtual]

This function finds the baryon density for which the pressure vanishes.

14.13.2.10 virtual int hadronic\_eos::saturation ( ) [virtual]

**Idea for Future** It would be great to provide numerical uncertainties in the saturation properties.

Reimplemented in [rmf\\_eos](#), and [rmf\\_delta\\_eos](#).

14.13.2.11 int hadronic\_eos::gradient\_qij ( fermion & n, fermion & p, thermo & th, double & qnn, double & qnp, double & qpp, double & dqnnndnn, double & dqnnndnp, double & dqnpdnn, double & dqnpdnp, double & dqppdnn, double & dqppdnp )

Note

This is still somewhat experimental.

We want the gradient part of the Hamiltonian in the form

$$\mathcal{H}_{\text{grad}} = \frac{1}{2} \sum_{i=n,p} \sum_{j=n,p} Q_{ij} \vec{\nabla} n_i \cdot \vec{\nabla} n_j$$

The expression for the gradient terms from [Pethick95](#) is

$$\begin{aligned} \mathcal{H}_{\text{grad}} = & -\frac{1}{4} (2P_1 + P_{1,f} - P_{2,f}) \\ & + \frac{1}{2} (Q_1 + Q_2) (n_n \nabla^2 n_n + n_p \nabla^2 n_p) \\ & + \frac{1}{4} (Q_1 - Q_2) [(\nabla n_n)^2 + (\nabla n_p)^2] \\ & + \frac{1}{2} \frac{dQ_2}{dn} (n_n \nabla n_n + n_p \nabla n_p) \cdot \nabla n \end{aligned}$$

This can be rewritten

$$\begin{aligned} \mathcal{H}_{\text{grad}} = & \frac{1}{2} (\nabla n)^2 \left[ \frac{3}{2} P_1 + n \frac{dP_1}{dn} \right] \\ & - \frac{3}{4} [(\nabla n_n)^2 + (\nabla n_p)^2] \\ & - \frac{1}{2} \nabla n \cdot \nabla n \frac{dQ_1}{dn} \\ & - \frac{1}{4} (\nabla n)^2 P_2 - \frac{1}{4} [(\nabla n_n)^2 + (\nabla n_p)^2] Q_2 \end{aligned}$$

or

$$\begin{aligned} \mathcal{H}_{\text{grad}} = & \frac{1}{4} (\nabla n)^2 \left[ 3P_1 + 2n \frac{dP_1}{dn} - P_2 \right] \\ & - \frac{1}{4} (3Q_1 + Q_2) [(\nabla n_n)^2 + (\nabla n_p)^2] \\ & - \frac{1}{2} \frac{dQ_1}{dn} [n_n \nabla n_n + n_p \nabla n_p] \cdot \nabla n \end{aligned}$$

or

$$\begin{aligned} \mathcal{H}_{\text{grad}} = & \frac{1}{4} (\nabla n)^2 \left[ 3P_1 + 2n \frac{dP_1}{dn} - P_2 \right] \\ & - \frac{1}{4} (3Q_1 + Q_2) [(\nabla n_n)^2 + (\nabla n_p)^2] \\ & - \frac{1}{2} \frac{dQ_1}{dn} [n_n (\nabla n_n)^2 + n_p (\nabla n_p)^2 + n \nabla n_n \cdot \nabla n_p] \end{aligned}$$

Generally, for Skyrme-like interactions

$$P_i = \frac{1}{4}t_i \left( 1 + \frac{1}{2}x_i \right)$$

$$Q_i = \frac{1}{4}t_i \left( \frac{1}{2} + x_i \right).$$

for  $i = 1, 2$ .

This function uses the assumption  $x_1 = x_2 = 0$  to calculate  $t_1$  and  $t_2$  from the neutron and proton effective masses assuming the Skyrme form. The values of  $Q_{ij}$  and their derivatives are then computed.

The functions [set\\_n\\_and\\_p\(\)](#) and [set\\_thermo\(\)](#) will be called by [gradient\\_qij\(\)](#), to facilitate the use of the `n`, `p`, and `th` parameters.

14.13.2.12 `double hadronic_eos::calc_pressure_nb ( double nb, const double & alpha = 0.0 )`

Used by [fcomp\(\)](#).

14.13.2.13 `double hadronic_eos::calc_edensity_nb ( double nb, const double & alpha = 0.0 )`

This function calls [hadronic\\_eos::calc\\_e\(\)](#) with the internally stored neutron and proton objects.

14.13.2.14 `double hadronic_eos::calc_press_over_den2 ( double nb, const double & alpha = 0.0 )`

Used by [fkprime\(\)](#).

14.13.2.15 `double hadronic_eos::calc_edensity_alpha ( double alpha, const double & nb )`

Used by [fesym\(\)](#).

This function calls [hadronic\\_eos::calc\\_e\(\)](#) with the internally stored neutron and proton objects.

14.13.2.16 `virtual int hadronic_eos::set_mroot ( mroot<mm_funct<>> & mr ) [virtual]`

Note

While in principle this allows one to use any **mroot** object, in practice some of the current EOSs require **gsl\_mroot\_hybrids** because it automatically avoids regions where the equations are undefined.

14.13.2.17 `virtual int hadronic_eos::set_sat_root ( root<funct> & mr ) [virtual]`

Note

While in principle this allows one to use any **mroot** object, in practice some of the current EOSs require **gsl\_mroot\_hybrids** because it automatically avoids regions where the equations are undefined.

14.13.2.18 `virtual int hadronic_eos::set_sat_deriv2 ( deriv<funct> & de ) [virtual]`

Computing the slope of the symmetry energy at the saturation density requires two derivative objects, because it has to take an isospin derivative and a density derivative. Thus this second **deriv** object is used in the function [fesym\\_slope\(\)](#).

### 14.13.3 Field Documentation

14.13.3.1 `gsl_deriv<funct> hadronic_eos::def_deriv`

The value of **gsl\_deriv::h** is set to  $10^{-3}$  in the [hadronic\\_eos](#) constructor.

Definition at line 608 of file `hadronic_eos.h`.

## 14.13.3.2 gsl\_deriv&lt;funct&gt; hadronic\_eos::def\_deriv2

The value of `gsl_deriv::h` is set to  $10^{-3}$  in the `hadronic_eos` constructor.

Definition at line 615 of file `hadronic_eos.h`.

## 14.13.3.3 gsl\_mroot\_hybrids&lt;mm\_funct&lt;&gt;&gt; hadronic\_eos::def\_mroot

Used by `calc_e()` to solve `nuc_matter_p()` (2 variables) and by `calc_p()` to solve `nuc_matter_e()` (2 variables).

Definition at line 622 of file `hadronic_eos.h`.

## 14.13.3.4 cern\_mroot\_root&lt;funct&gt; hadronic\_eos::def\_sat\_root

Used by `fn0()` (which is called by `saturation()`) to solve `saturation_matter_e()` (1 variable).

Definition at line 630 of file `hadronic_eos.h`.

The documentation for this class was generated from the following file:

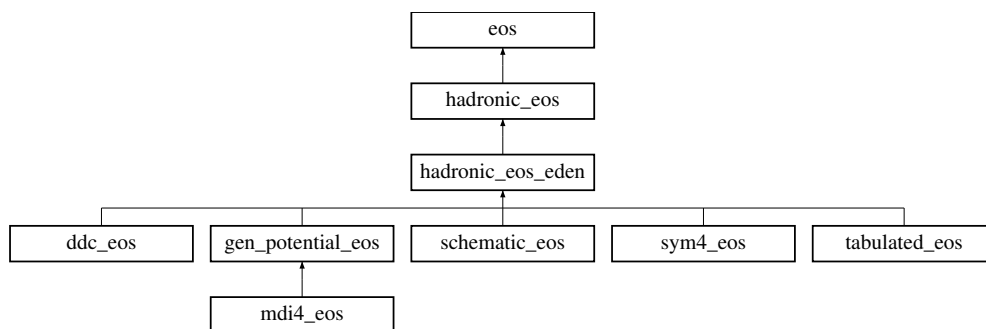
- `hadronic_eos.h`

## 14.14 hadronic\_eos\_edem Class Reference

A hadronic EOS based on a function of the densities [abstract base].

```
#include <hadronic_eos.h>
```

Inheritance diagram for `hadronic_eos_edem`:



## 14.14.1 Detailed Description

Definition at line 669 of file `hadronic_eos.h`.

## Public Member Functions

- virtual int `calc_e` (**fermion** &n, **fermion** &p, **thermo** &th)=0  
*Equation of state as a function of density.*
- virtual int `calc_p` (**fermion** &n, **fermion** &p, **thermo** &th)  
*Equation of state as a function of the chemical potentials.*

The documentation for this class was generated from the following file:

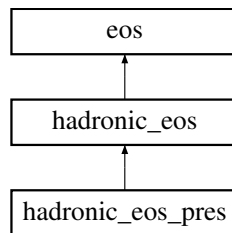
- `hadronic_eos.h`

## 14.15 hadronic\_eos\_pres Class Reference

A hadronic EOS based on a function of the chemical potentials [abstract base].

```
#include <hadronic_eos.h>
```

Inheritance diagram for hadronic\_eos\_pres:



### 14.15.1 Detailed Description

Definition at line 685 of file hadronic\_eos.h.

#### Public Member Functions

- virtual int [calc\\_p](#) (**fermion** &n, **fermion** &p, **thermo** &th)=0  
*Equation of state as a function of the chemical potentials.*
- virtual int [calc\\_e](#) (**fermion** &n, **fermion** &p, **thermo** &th)  
*Equation of state as a function of density.*

The documentation for this class was generated from the following file:

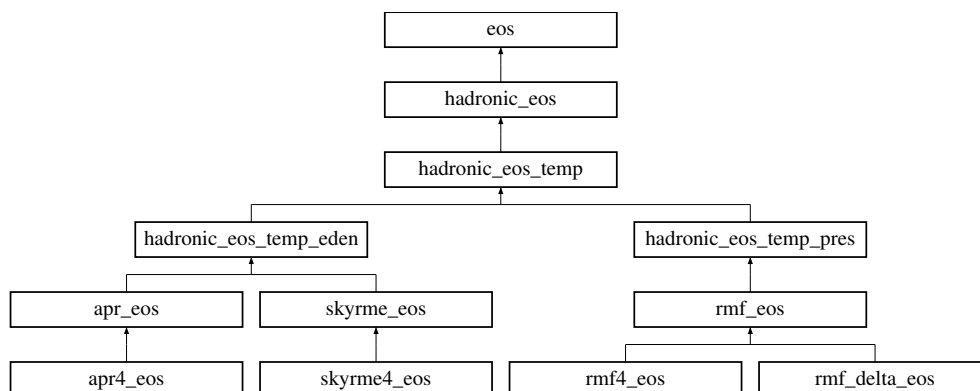
- hadronic\_eos.h

## 14.16 hadronic\_eos\_temp Class Reference

A finite temperature hadronic EOS.

```
#include <hadronic_eos.h>
```

Inheritance diagram for hadronic\_eos\_temp:



## 14.16.1 Detailed Description

Definition at line 699 of file hadronic\_eos.h.

## Public Member Functions

- virtual int [calc\\_e](#) (**fermion** &n, **fermion** &p, **thermo** &th)=0  
*Equation of state as a function of density.*
- virtual int [calc\\_temp\\_e](#) (**fermion** &n, **fermion** &p, double T, **thermo** &th)=0  
*Equation of state as a function of densities at finite temperature.*
- virtual int [calc\\_p](#) (**fermion** &n, **fermion** &p, **thermo** &th)=0  
*Equation of state as a function of the chemical potentials.*
- virtual int [calc\\_temp\\_p](#) (**fermion** &n, **fermion** &p, double T, **thermo** &th)=0  
*Equation of state as a function of the chemical potentials at finite temperature.*

## Data Fields

- **eff\_fermion** [def\\_fet](#)  
*Desc.*

## Protected Member Functions

- int [nuc\\_matter\\_temp\\_e](#) (size\_t nv, const **ovector\_base** &x, **ovector\_base** &y, double \*&pa)  
*Solve for nuclear matter at finite temperature given density.*
- int [nuc\\_matter\\_temp\\_p](#) (size\_t nv, const **ovector\_base** &x, **ovector\_base** &y, double \*&pa)  
*Solve for nuclear matter at finite temperature given mu.*

## Protected Attributes

- **fermion\_eval\_thermo** \* [fet](#)  
*Desc.*
- double [IT](#)  
*The temperature.*

The documentation for this class was generated from the following file:

- hadronic\_eos.h

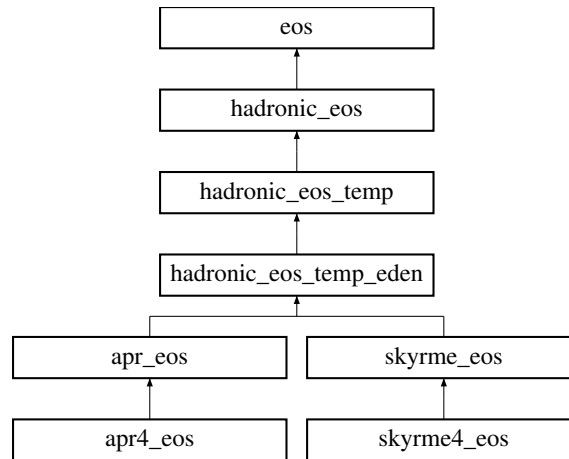
## 14.17 hadronic\_eos\_temp\_edn Class Reference

A hadronic EOS at finite temperature based on a function of the densities [abstract base].

```
#include <hadronic_eos.h>
```

Inheritance diagram for hadronic\_eos\_temp\_edn:





#### 14.17.1 Detailed Description

Definition at line 757 of file `hadronic_eos.h`.

#### Public Member Functions

- virtual int `calc_e` (**fermion** &n, **fermion** &p, **thermo** &th)=0  
*Equation of state as a function of density.*
- virtual int `calc_temp_e` (**fermion** &n, **fermion** &p, double T, **thermo** &th)=0  
*Equation of state as a function of densities at finite temperature.*
- virtual int `calc_p` (**fermion** &n, **fermion** &p, **thermo** &th)  
*Equation of state as a function of the chemical potentials.*
- virtual int `calc_temp_p` (**fermion** &n, **fermion** &p, double T, **thermo** &th)  
*Equation of state as a function of the chemical potentials at finite temperature.*

The documentation for this class was generated from the following file:

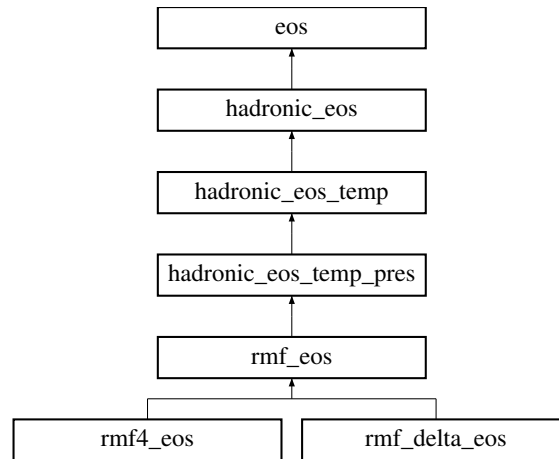
- `hadronic_eos.h`

## 14.18 hadronic\_eos\_temp\_pres Class Reference

A hadronic EOS at finite temperature based on a function of the chemical potentials [abstract base].

```
#include <hadronic_eos.h>
```

Inheritance diagram for `hadronic_eos_temp_pres`:



### 14.18.1 Detailed Description

Definition at line 785 of file hadronic\_eos.h.

#### Public Member Functions

- virtual int [calc\\_p](#) (**fermion** &n, **fermion** &p, **thermo** &th)=0  
*Equation of state as a function of the chemical potentials.*
- virtual int [calc\\_temp\\_p](#) (**fermion** &n, **fermion** &p, double T, **thermo** &th)=0  
*Equation of state as a function of the chemical potentials at finite temperature.*
- virtual int [calc\\_e](#) (**fermion** &n, **fermion** &p, **thermo** &th)  
*Equation of state as a function of density.*
- virtual int [calc\\_temp\\_e](#) (**fermion** &n, **fermion** &p, double T, **thermo** &th)  
*Equation of state as a function of densities at finite temperature.*

The documentation for this class was generated from the following file:

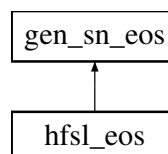
- hadronic\_eos.h

## 14.19 hfsl\_eos Class Reference

The Hempel et al. supernova EOSs.

```
#include <gen_sn_eos.h>
```

Inheritance diagram for hfsl\_eos:



### 14.19.1 Detailed Description

This class is experimental.

## Note

O<sub>2scl</sub> Does not contain the EOS, only provides some code to manipulate it. This class was designed to be used with the files `dd2_frdm_eos_shen98format_v1.02.tab`, `fsg_roca_eos_shen98format_v1.0.tab`, and `nl3_lala_eos_shen98format_v1.0.tab` as obtained from <http://phys-merger.physik.unibas.ch/~hempel/eos.html>.

See also the documentation at [gen\\_sn\\_eos](#).

See [Hempel10](#) and [Hempel11](#).

Definition at line 856 of file `gen_sn_eos.h`.

## Public Member Functions

- virtual void [load](#) (std::string fname)  
*Load table from filename fname.*
- virtual void [beta\\_eq\\_T0](#) (size\_t i, double &nb, double &E\_beta, double &P\_beta, double &Ye\_beta, double &Z\_beta, double &A\_beta)  
*Compute properties of matter in beta equilibrium at zero temperature at a baryon density grid point.*

## Data Fields

- [tensor\\_grid3](#) & [log\\_rho](#)  
*Logarithm of baryon number density in g/cm<sup>3</sup>.*
- [tensor\\_grid3](#) & [nB](#)  
*Baryon number density in 1/fm<sup>3</sup>.*
- [tensor\\_grid3](#) & [log\\_Y](#)  
*Logarithm of proton fraction.*
- [tensor\\_grid3](#) & [Yp](#)  
*Proton fraction.*
- [tensor\\_grid3](#) & [M\\_star](#)  
*Nucleon effective mass in MeV.*
- [tensor\\_grid3](#) & [A\\_light](#)  
*Mass number of light fragments.*
- [tensor\\_grid3](#) & [Z\\_light](#)  
*Proton number of light fragments.*

The documentation for this class was generated from the following file:

- `gen_sn_eos.h`

## 14.20 rmf\_nucleus::initial\_guess Struct Reference

Initial guess structure.

```
#include <rmf_nucleus.h>
```

## 14.20.1 Detailed Description

The initial guess for the meson field profiles is a set of fermi-dirac functions, i.e.

$$\sigma(r) = \sigma_0 / [1 + \exp((r - \text{fermi\_radius})/\text{fermi\_width})]$$

Definition at line 499 of file `rmf_nucleus.h`.

## Data Fields

- double [sigma0](#)  
*Scalar field at  $r=0$ .*
- double [omega0](#)  
*Vector field at  $r=0$ .*
- double [rho0](#)  
*Isuvector field at  $r=0$ .*
- double [A0](#)  
*Coulomb field at  $r=0$ .*
- double [fermi\\_radius](#)  
*The radius for which the fields are half their central value.*
- double [fermi\\_width](#)  
*The "width" of the Fermi-Dirac function.*

The documentation for this struct was generated from the following file:

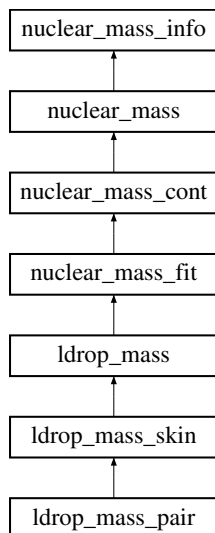
- `rmf_nucleus.h`

## 14.21 ldrop\_mass Class Reference

Simple liquid drop mass formula.

```
#include <ldrop_mass.h>
```

Inheritance diagram for ldrop\_mass:



## 14.21.1 Detailed Description

Includes bulk part plus surface and Coulomb (no pairing) without neutron skin and without any isospin contribution to the surface energy.

The NL4 EOS is loaded by default.

**Warning**

This class sets `part::inc_rest_mass` to true for the particle objects specified in `set_n_and_p()`.

**Central densities**

Given a saturation density,  $n_0$  and a transition density,  $n_t$ , we set  $I = 1 - 2Z/A$ , and then assume  $\delta = I$ . We further assume that the isospin-asymmetric saturation density  $n_L$  is

$$n_L = n_0 + n_1 \delta^2$$

and then we can compute  $n_p = (1 - \delta)/2n_L$  and  $n_n = (1 + \delta)/2n_L$ .

Note that  $\delta = I$  implies no neutron skin. A neutron skin occurs when  $\delta < I$ , and  $\delta = 0$  implies a "maximum skin size" which is occurs when no extra neutrons are in center and all extra neutrons are located in the skin, i.e.  $N_{\text{skin}} = N - Z$ .

**Nuclear radii**

The neutron and proton radii are determined from the central densities with

$$R_n = \left( \frac{3N}{4\pi n_n} \right)^{1/3}$$

$$R_p = \left( \frac{3Z}{4\pi n_p} \right)^{1/3}$$

**Bulk energy contribution**

The bulk binding energy contribution ( $\sim -16$  MeV per nucleon) and the symmetry energy are computed using the hadronic EOS (either `def_had_eos` or the EOS specified in the most recent call to `set_hadronic_eos_temp()`). The bulk energy per baryon is

$$E_{\text{bulk}}/A = \frac{\hbar c}{n_L} [\mathcal{E}(n_n, n_p) - n_n m_n - n_p m_p]$$

**Surface energy contribution**

The surface energy density is (Ravenhall83)

$$\varepsilon = \frac{\chi d\sigma}{R}$$

where  $\sigma$  is the surface tension. The factor  $\chi$  is typically taken care of by the caller, so we ignore it for now. To compute the surface energy per baryon, we divide by the baryon density,  $n_n + n_p$ . We can rewrite this

$$E_{\text{surf}} = \frac{3\sigma}{n_n + n_p} \left[ \frac{3A}{4(n_n + n_p)\pi} \right]^{-1/3}$$

or

$$E_{\text{surf}} = \frac{\sigma}{n_L} \left( \frac{36\pi n_L}{A} \right)^{1/3}$$

where the surface tension  $\sigma$  (in MeV) is given in `surften`.

Taking a typical value,  $\sigma = 1.1$  MeV and  $n_L = 0.16 \text{ fm}^{-3}$ , gives the standard result,  $E_{\text{surf}}/A = 18 \text{ MeV } A^{-1/3}$ .

**Coulomb energy contribution**

The Coulomb energy density (Ravenhall83) is

$$\varepsilon_{\text{Coul}} = \frac{4\pi}{5} n_p^2 e^2 R_p^2$$

The energy per baryon is

$$E_{\text{Coul}}/A = \frac{4\pi}{5n_L} n_p^2 e^2 R_p^2$$

This is the expression used in the code, except for a prefactor `coul_coeff` which is a fit parameter and should be close to unity.

Assuming  $R_p = R$  and  $Z = \frac{4\pi}{3}R^3 n_p$  and  $R = [3A/(4\pi n_L)]^{1/3}$  gives

$$E_{\text{Coul}}/A = \frac{6^{2/3}}{5} \pi^{1/3} e^2 n_L^{1/3} \frac{Z^2}{A^{4/3}}$$

and taking  $n_L = 0.16 \text{ fm}^{-3}$  and  $e^2 = \hbar c/137$  gives the standard result

$$E_{\text{Coul}}/A = 0.76 \text{ MeV } Z^2 A^{-4/3}$$

---

## References

Designed for [Steiner08](#) based on [Lattimer85](#) and [Lattimer91](#).

---

Definition at line 237 of file ldrop\_mass.h.

## Public Member Functions

- virtual double [mass\\_excess\\_d](#) (double Z, double N)  
*Given Z and N, return the mass excess in MeV.*
- virtual double [drip\\_binding\\_energy\\_d](#) (double Z, double N, double npout, double nnout, double chi)  
*Given Z and N, return the binding energy in MeV.*
- virtual const char \* [type](#) ()  
*Return the type, "ldrop\_mass".*

## Fitting functions

- virtual int [fit\\_fun](#) (size\_t nv, const **ovector\_base** &x)  
*Fix parameters from an array for fitting.*
- virtual int [guess\\_fun](#) (size\_t nv, **ovector\_base** &x)  
*Fill array with guess from present values for fitting.*

## Data Fields

- **thermo** [th](#)  
*Energy and pressure.*

## Input parameters

- double [n1](#)  
*Density asymmetry (default 0)*
- double [n0](#)  
*Saturation density ( The default is  $0.16 \text{ fm}^{-3}$ )*
- double [surften](#)  
*Surface tension in MeV (default 1.1 MeV)*
- double [coul\\_coeff](#)  
*Coulomb coefficient (default 1.0)*

## Output quantities

- double [nn](#)  
*Internal average neutron density.*
  - double [np](#)  
*Internal average proton density.*
  - double [Rn](#)  
*Neutron radius.*
  - double [Rp](#)  
*Proton radius.*
  - double [bulk](#)  
*Bulk part of energy.*
  - double [surf](#)  
*Surface part of energy.*
  - double [coul](#)  
*Coulomb part of energy.*
-

## Protected Attributes

- [eff\\_fermion](#) [eff](#)  
*Desc.*
- [fermion](#) \* [n](#)  
*Pointer to neutron.*
- [fermion](#) \* [p](#)  
*Pointer to proton.*
- [hadronic\\_eos\\_temp](#) \* [heos](#)  
*The base EOS for bulk matter.*

## EOS and particle parameters

- [rmf\\_eos](#) [def\\_had\\_eos](#)  
*The default hadronic EOS.*
- [fermion](#) [def\\_neutron](#)  
*Default neutron.*
- [fermion](#) [def\\_proton](#)  
*Default proton.*
- int [set\\_hadronic\\_eos\\_temp](#) ([hadronic\\_eos\\_temp](#) &uhe)  
*Change the base hadronic EOS.*
- void [set\\_n\\_and\\_p](#) ([fermion](#) &un, [fermion](#) &up)  
*Change neutron and proton objects.*

## 14.21.2 Member Function Documentation

14.21.2.1 `virtual double ldrip_mass::mass_excess_d ( double Z, double N )` [virtual]

Implements [nuclear\\_mass\\_cont](#).

14.21.2.2 `virtual double ldrip_mass::drip_binding_energy_d ( double Z, double N, double npout, double nnout, double chi )` [virtual]

This function is currently independent of npout, nnout, and chi.

Reimplemented in [ldrip\\_mass\\_skin](#).

The documentation for this class was generated from the following file:

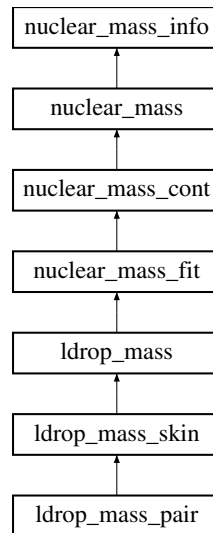
- [ldrip\\_mass.h](#)

## 14.22 ldrip\_mass\_pair Class Reference

Liquid drop model with pairing.

```
#include <ldrip_mass.h>
```

Inheritance diagram for [ldrip\\_mass\\_pair](#):



#### 14.22.1 Detailed Description

This class adds a pairing correction

$$E_{\text{pair}}/A = -\frac{\zeta}{2A^{3/2}} [\cos(Z\pi) + \cos(N\pi)]$$

where  $\zeta$  is stored in [Epair](#). The trigonometric functions give the expected result for integer values of N and Z.

Definition at line 581 of file ldrop\_mass.h.

#### Public Member Functions

- virtual const char \* [type](#) ()  
*Return the type, "ldrop\_mass\_pair".*
- virtual int [fit\\_fun](#) (size\_t nv, const **ovector\_base** &x)  
*Fix parameters from an array for fitting.*
- virtual int [guess\\_fun](#) (size\_t nv, **ovector\_base** &x)  
*Fill array with guess from present values for fitting.*
- virtual double [drip\\_binding\\_energy\\_full\\_d](#) (double Z, double N, double npout, double nnout, double chi, double T)  
*Return the free binding energy of a nucleus in a many-body environment.*

#### Data Fields

- double [Epair](#)  
*Pairing energy coefficient (default 13 MeV)*
- double [pair](#)  
*Most recently computed pairing energy per baryon.*

The documentation for this class was generated from the following file:

- ldrop\_mass.h

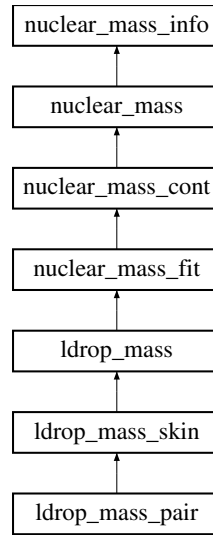
## 14.23 ldrop\_mass\_skin Class Reference

More advanced liquid drop model.

```
#include <ldrop_mass.h>
```



Inheritance diagram for ldrop\_mass\_skin:



#### 14.23.1 Detailed Description

In addition to the physics in [ldrop\\_mass](#), this includes corrections for

- finite temperature
- neutron skin
- an isospin-dependent surface energy
- decrease in the Coulomb energy from external protons

#### Note

The input parameter T should be given in units of inverse Fermis -- this is a bit unusual since the binding energy is returned in MeV, but we keep it for now.

#### Bulk energy

The central densities and radii,  $n_n, n_p, R_n, R_p$  are all determined in the same way as [ldrop\\_mass](#), except that now  $\delta \equiv I\zeta$ , where  $\zeta$  is stored in [doi](#). Note that this means  $N > Z$  iff  $R_n > R_p$ .

If [new\\_skin\\_mode](#) is false, then the bulk energy is also computed as in [ldrop\\_mass](#). Otherwise, the number of nucleons in the core is computed with

$$A_{\text{core}} = Z(n_n + n_p)/n_p \text{ for } N \geq Z$$

$$A_{\text{core}} = N(n_n + n_p)/n_p \text{ for } Z > N$$

and  $A_{\text{skin}} = A - A_{\text{core}}$ . The core contribution to the bulk energy is

$$E_{\text{core}}/A = \left( \frac{A_{\text{core}}}{A} \right) \frac{\hbar c}{n_L} [\mathcal{E}(n_n, n_p) - n_n m_n - n_p m_p]$$

then the skin contribution is

$$E_{\text{skin}}/A = \left( \frac{A_{\text{skin}}}{A} \right) \frac{\hbar c}{n_L} [\mathcal{E}(n_n, 0) - n_n m_n] \text{ for } N > Z$$

and

$$E_{\text{skin}}/A = \left( \frac{A_{\text{skin}}}{A} \right) \frac{\hbar c}{n_L} [\varepsilon(0, n_p) - n_p m_p] \text{ for } Z > N$$

### Surface energy

If `full_surface` is false, then the surface energy is just that from `ldrop_mass`, with an extra factor for the surface symmetry energy

$$E_{\text{surf}} = \frac{\sigma}{n_L} \left( \frac{36\pi n_L}{A} \right)^{1/3} (1 - \sigma_\delta \delta^2)$$

where  $\sigma_\delta$  is unitless and stored in `ss`.

If `full_surface` is true, then the surface energy is modified by a cubic dependence for the medium and contains finite temperature corrections.

### Coulomb energy

The Coulomb energy density ([Ravenhall83](#)) is

$$\varepsilon = 2\pi e^2 R_p^2 n_p^2 f_d(\chi_p)$$

where the function  $f_d(\chi)$  is

$$f_d(\chi_p) = \frac{1}{(d+2)} \left[ \frac{2}{(d-2)} \left( 1 - \frac{d}{2} \chi_p^{(1-2/d)} \right) + \chi_p \right]$$

This class takes  $d = 3$ .

### Todos and Future

**Todo** This is based on LPRL, but it's a little different in Lattimer and Swesty. I should document what the difference is.

**Todo** The testing could be updated.

**Idea for Future** Add translational energy?

**Idea for Future** Remove excluded volume correction and compute nuclear mass relative to the gas rather than relative to the vacuum.

**Idea for Future** In principle,  $T_c$  should be self-consistently determined from the EOS.

**Idea for Future** Does this work if the nucleus is "inside-out"?

---

### References

Designed in [Steiner08](#) and [Souza09](#) based in part on [Lattimer85](#) and [Lattimer91](#).

---

Definition at line 497 of file `ldrop_mass.h`.

---

### Public Member Functions

- virtual const char \* [type](#) ()  
*Return the type, "ldrop\_mass\_skin".*
  - virtual int [fit\\_fun](#) (size\_t nv, const **ovector\_base** &x)  
*Fix parameters from an array for fitting.*
  - virtual int [guess\\_fun](#) (size\_t nv, **ovector\_base** &x)  
*Fill array with guess from present values for fitting.*
  - virtual double [drip\\_binding\\_energy\\_d](#) (double Z, double N, double npout, double nnout, double chi)  
*Return the free binding energy of a nucleus in a many-body environment.*
  - virtual double [drip\\_binding\\_energy\\_full\\_d](#) (double Z, double N, double npout, double nnout, double chi, double T)  
*Return the free binding energy of a nucleus in a many-body environment.*
-

## Data Fields

- bool `full_surface`  
*If true, properly fix the surface for the pure neutron matter limit (default true)*
- bool `new_skin_mode`  
*If true, separately compute the skin for the bulk energy (default false)*
- double `doi`  
*Ratio of  $\delta/I$  (default 0.8).*
- double `ss`  
*Surface symmetry energy (default 0.5)*
- bool `rel_vacuum`  
*If true, define the nuclear mass relative to the vacuum (default true)*
- double `Tchalf`  
*The critical temperature of isospin-symmetric matter in  $\text{fm}^{-1}$  (default  $20.085/(\hbar c)$ .)*

## Input parameters for temperature dependence

- double `pp`  
*Exponent (default 1.25)*
- double `a0`  
*Coefficient (default 0.935)*
- double `a2`  
*Coefficient (default -5.1)*
- double `a4`  
*Coefficient (default -1.1)*

The documentation for this class was generated from the following file:

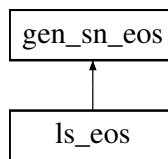
- `ldrop_mass.h`

## 14.24 ls\_eos Class Reference

The Lattimer-Swesty supernova EOS.

```
#include <gen_sn_eos.h>
```

Inheritance diagram for `ls_eos`:



## 14.24.1 Detailed Description

This class is experimental.

## Note

`O2scl` Does not contain the Lattimer-Swesty EOS, only provides some code to manipulate it. This class is designed to be used with the files `ls.dat`, `sk1.dat`, `ska.dat` and `skm.dat` as provided on Jim Lattimer's website, <http://www.astro.sunysb.edu/lattimer/EOS/main.html>.

Note that the tables on this website are different than what is generated from the LS Fortran code. See `oo_eos` to read O'Connor and Ott's tables generated from the LS Fortran code.

The four models are

- LS (K=370, Sv=31)
- SKI' (K=371, Sv=30.4)
- SKa (K=263, Sv=34.5)
- SKM\* (K=217, Sv=31.4)

#### Note

In the original table, the full internal energy per baryon (data section 4 of 26) is apparently based on a rest mass of  $Y_e m_p + (1 - Y_e) m_n$ , while the baryon part of the internal energy per baryon (data section 13 of 26) is based on a rest mass of  $m_n$ . This means that

$$E - E_{\text{int}} = E_{\text{eg}} - Y_e(m_n - m_p)$$

where  $E_{\text{eg}}$  is the energy per baryon of electrons and photons. In order to keep things consistent with the other EOS tables, when the EOS table is loaded, `gen_sn_eos::Eint` is rescaled to a rest mass of  $Y_e m_p + (1 - Y_e) m_n$ .

See also the documentation at [gen\\_sn\\_eos](#).

See [Lattimer91](#) and [Lattimer85](#).

**Todo** There are still a few points for which the electron/photon EOS seems to be off.

Definition at line 371 of file `gen_sn_eos.h`.

#### Public Member Functions

- virtual void [load](#) (std::string fname)  
*Load table from filename fname.*
- int [check\\_eg](#) (test\_mgr &tm)  
*Check electrons and photons.*
- virtual void [beta\\_eq\\_T0](#) (size\_t i, double &nb, double &E\_beta, double &P\_beta, double &Ye\_beta, double &Z\_beta, double &A\_beta)  
*Compute properties of matter in beta equilibrium at zero temperature at a baryon density grid point.*

#### Data Fields

- [tensor\\_grid3](#) & [fill](#)  
*Filling factor for nuclei.*
- [tensor\\_grid3](#) & [nb\\_in](#)  
*Baryon number density inside nuclei in fm<sup>-3</sup>.*
- [tensor\\_grid3](#) & [dPdN](#)  
*Derivative of pressure with respect to baryon density.*
- [tensor\\_grid3](#) & [dPdT](#)  
*Derivative of pressure with respect to temperature.*
- [tensor\\_grid3](#) & [dPdY](#)  
*Derivative of pressure with respect to electron fraction.*
- [tensor\\_grid3](#) & [dsdT](#)  
*Derivative of entropy with respect to temperature.*
- [tensor\\_grid3](#) & [dsdY](#)  
*Derivative of entropy with respect to electron fraction.*
- [tensor\\_grid3](#) & [Nskin](#)  
*Number of neutrons in skin.*
- [tensor\\_grid3](#) & [nb\\_out](#)  
*Baryon density outside nuclei in fm<sup>-3</sup>.*
- [tensor\\_grid3](#) & [x\\_out](#)  
*Proton fraction outside nuclei.*
- [tensor\\_grid3](#) & [mu](#)  
*Out of whackness parameter,  $\mu_n - \mu_p - \mu_e + 1.293$  MeV, in MeV.*

## 14.24.2 Member Function Documentation

## 14.24.2.1 int ls\_eos::check\_eg ( test\_mgr &amp; tm )

This checks that the electron and photon thermodynamics generated by O<sub>2</sub>scl is consistent with the data in E, Eint, F, Fint, P, Pint, S, and Sint.

## 14.24.2.2 virtual void ls\_eos::beta\_eq.T0 ( size\_t i, double &amp; nb, double &amp; E\_beta, double &amp; P\_beta, double &amp; Ye\_beta, double &amp; Z\_beta, double &amp; A\_beta ) [inline, virtual]

This EOS table doesn't have T=0 results, so we extrapolate from the two low-temperature grid points.

Implements [gen\\_sn\\_eos](#).

Definition at line 432 of file gen\_sn\_eos.h.

The documentation for this class was generated from the following file:

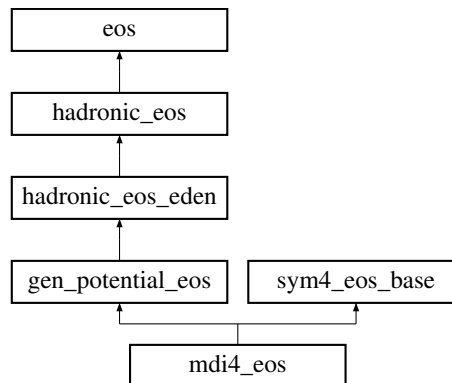
- gen\_sn\_eos.h

## 14.25 mdi4\_eos Class Reference

A version of [gen\\_potential\\_eos](#) to separate potential and kinetic contributions.

```
#include <sym4_eos.h>
```

Inheritance diagram for mdi4\_eos:



## 14.25.1 Detailed Description

**References:**

Created for [Steiner06](#).

Definition at line 163 of file sym4\_eos.h.

**Public Member Functions**

- virtual int [calc\\_e\\_sep](#) (fermion &ne, fermion &pr, double &ed\_kin, double &ed\_pot, double &mu\_n\_kin, double &mu\_p\_kin, double &mu\_n\_pot, double &mu\_p\_pot)  
*Compute the potential and kinetic parts separately.*
- virtual int [test\\_separation](#) (fermion &ne, fermion &pr, test\_mgr &t)  
*Test the separation of the potential and kinetic energy parts.*

## Protected Member Functions

- double [energy\\_kin](#) (double var)  
Compute the kinetic part of the energy density.
- double [energy\\_pot](#) (double var)  
Compute the potential part of the energy density.

The documentation for this class was generated from the following file:

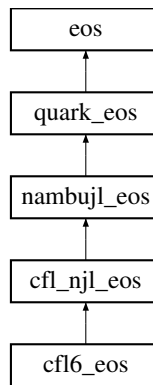
- sym4\_eos.h

## 14.26 nambu\_jl\_eos Class Reference

Nambu Jona-Lasinio EOS at zero temperature.

```
#include <nambu_jl_eos.h>
```

Inheritance diagram for nambu\_jl\_eos:



## 14.26.1 Detailed Description

Calculates everything from the quark condensates ([uds].qq) and the chemical potentials ([uds].mu). If "fromqq" is set to false, then instead it calculates everything from the dynamical masses ([uds].ms) and the chemical potentials. L, G, K, and B0 are fixed constants. [uds].pr returns the pressure due to the Fermi-gas contribution plus the bag pressure contribution. [uds.ed] is the energy density for each quark so that e.g. u.ed+u.pr=u.mu\*u.n. B0 should be fixed using calc\_B0() beforehand to ensure that the energy density and pressure of the vacuum is zero.

The functions [set\\_parameters\(\)](#) should be called first.

The code is based on [Buballa99](#).

The Lagrangian is

$$\mathcal{L} = \bar{q}(i\cancel{D} - \hat{m}_0)q + G \sum_{k=0}^8 [(\bar{q}\lambda_k q)^2 + (\bar{q}i\gamma_5 \lambda_k q)^2] + \mathcal{L}_6$$

$$\mathcal{L}_6 = -K [\det_f(\bar{q}(1 + \gamma_5)q) + \det_f(\bar{q}(1 - \gamma_5)q)].$$

And the corresponding thermodynamic potential is

$$\Omega = \Omega_{FG} + \Omega_{Int}$$

where  $\Omega_{FG}$  is the Fermi gas contribution and

$$\frac{\Omega_{Int}}{V} = -2N_c \sum_{i=u,d,s} \int \frac{d^3p}{(2\pi)^3} \sqrt{m_i^2 + p^2} + \frac{\Omega_V}{V}$$

$$\frac{\Omega_V}{V} = \sum_{i=u,d,s} 2G \langle \bar{q}_i q_i \rangle^2 - 4K \langle \bar{q}_u q_u \rangle \langle \bar{q}_d q_d \rangle \langle \bar{q}_s q_s \rangle + B_0.$$

where  $B_0$  is a constant defined to ensure that the energy density and the pressure of the vacuum is zero.

Unlike [Buballa99](#), the bag constant,  $\Omega_{Int}/V$  is defined without the term

$$\sum_{i=u,d,s} 2N_C \int_0^\Lambda \frac{d^3 p}{(2\pi)^3} \sqrt{m_{0,i}^2 + p^2} dp$$

since this allows an easier comparison to the finite temperature EOS. The constant  $B_0$  in this case is therefore significantly larger, but the energy density and pressure are still zero in the vacuum.

The Feynman-Hellman theorem ([Bernard88](#)), gives

$$\langle \bar{q} q \rangle = \frac{\partial m^*}{\partial m}$$

The functions [calc\\_e\(\)](#) and [calc\\_p\(\)](#) never return a value other than zero, but will give nonsensical results for nonsensical inputs.

### Finite T documentation

Calculates everything from the quark condensates ([uds].qq) and the chemical potentials ([uds].mu). If "frommq" is set to false, then instead it calculates everything from the dynamical masses ([uds].ms) and the chemical potentials. L, G, K, and B0 are fixed constants. [uds].pr returns the pressure due to the Fermi-gas contribution plus the bag pressure contribution. [uds.ed] is the energy density for each quark so that e.g. u.ed+u.pr=u.mu\*u.n. B0 is fixed to ensure that the energy density and pressure of the vacuum is zero.

This implementation includes contributions from antiquarks.

---

### References:

Created for [Steiner00](#). See also [Buballa99](#) and [Hatsuda94](#).

Definition at line 128 of file nambujl\_eos.h.

### Data Structures

- struct [njtp\\_s](#)  
A structure for passing parameters to the integrands.

### Public Types

- typedef struct [nambujl\\_eos::njtp\\_s](#) njtp  
A structure for passing parameters to the integrands.

### Public Member Functions

- virtual int [set\\_parameters](#) (double lambda=0.0, double fourferm=0.0, double sixferm=0.0)  
Set the parameters and the bag constant B0.
  - virtual int [calc\\_p](#) (quark &u, quark &d, quark &s, thermo &lth)  
Equation of state as a function of chemical potentials.
  - virtual int [calc\\_temp\\_p](#) (quark &u, quark &d, quark &s, double T, thermo &th)  
Equation of state as a function of chemical potentials at finite temperature.
  - virtual int [calc\\_eq\\_p](#) (quark &u, quark &d, quark &s, double &gap1, double &gap2, double &gap3, thermo &lth)  
Equation of state and gap equations as a function of chemical potential.
  - virtual int [calc\\_eq\\_e](#) (quark &u, quark &d, quark &s, double &gap1, double &gap2, double &gap3, thermo &lth)  
Equation of state and gap equations as a function of the densities.
  - int [calc\\_eq\\_temp\\_p](#) (quark &tu, quark &td, quark &ts, double &gap1, double &gap2, double &gap3, thermo &qb, double temper)
-

*Equation of state and gap equations as a function of chemical potentials.*

- int [gapfunms](#) (size\_t nv, const **ovector\_base** &x, **ovector\_base** &y)  
*Calculates gap equations in y as a function of the constituent masses in x.*
- int [gapfunqq](#) (size\_t nv, const **ovector\_base** &x, **ovector\_base** &y)  
*Calculates gap equations in y as a function of the quark condensates in x.*
- int [gapfunmsT](#) (size\_t nv, const **ovector\_base** &x, **ovector\_base** &y)  
*Calculates gap equations in y as a function of the constituent masses in x.*
- int [gapfunqqT](#) (size\_t nv, const **ovector\_base** &x, **ovector\_base** &y)  
*Calculates gap equations in y as a function of the quark condensates in x.*
- int [set\\_quarks](#) (**quark** &u, **quark** &d, **quark** &s)  
*Set the quark objects to use.*
- virtual const char \* [type](#) ()  
*Return string denoting type ("nambu\_eos")*
- virtual int [set\\_solver](#) (**mroot**< **mm\_funct**<> > &s)  
*Set solver to use in [set\\_parameters\(\)](#)*
- virtual int [set\\_inte](#) (**inte**< **funct** > &i)  
*Set integration object.*

#### Data Fields

- double [limit](#)  
*Accuracy limit for Fermi integrals for finite temperature.*
- bool [fromqq](#)  
*Calculate from quark condensates if true (default true)*
- double [L](#)  
*The momentum cutoff.*
- double [G](#)  
*The four-fermion coupling.*
- double [K](#)  
*The 't Hooft six-fermion interaction coupling.*
- double [B0](#)  
*The bag constant.*
- **gsl\_mroot\_hybrids**< **mm\_funct**<> > [def\\_solver](#)  
*The default solver.*
- **gsl\_inte\_qag**< **funct** > [def\\_it](#)  
*The default integrator.*

#### The default quark masses

These are the values from [Buballa99](#) which were used to fix the pion and kaon decay constants, and the pion, kaon, and eta prime masses. They are set in the constructor and are in units of fm<sup>-1</sup>.

- double **up\_default\_mass**
- double **down\_default\_mass**
- double **strange\_default\_mass**

#### The default quark objects

The masses are automatically set in the constructor to `up_default_mass`, `down_default_mass`, and `strange_default_mass.c`

- **quark def\_up**
- **quark def\_down**
- **quark def\_strange**



## Protected Member Functions

- `int B0fun (size_t nv, const ovector_base &x, ovector_base &y)`  
*Used by calc\_B0() to compute the bag constant.*
- `void njbag (quark &q)`  
*Calculates the contribution to the bag constant from quark q.*
- `double iqq (double x, const njtp &pa)`  
*The integrand for the quark condensate.*
- `double ide (double x, const njtp &pa)`  
*The integrand for the density.*
- `double ied (double x, const njtp &pa)`  
*The integrand for the energy density.*
- `double ipr (double x, const njtp &pa)`  
*The integrand for the pressure.*

## Protected Attributes

- `inte< funct > * it`  
*The integrator for finite temperature integrals.*
- `mroot< mm_funct<>> * solver`  
*The solver to use for set\_parameters()*
- `quark * up`  
*The up quark.*
- `quark * down`  
*The down quark.*
- `quark * strange`  
*The strange quark.*
- `double cp_temp`  
*The temperature for calc\_temp\_p()*

## 14.26.2 Member Function Documentation

14.26.2.1 `virtual int nambu_l_eos::set_parameters ( double lambda = 0.0, double fourferm = 0.0, double sixferm = 0.0 ) [virtual]`

This function allows the user to specify the momentum cutoff, `lambda`, the four-fermion coupling `fourferm` and the six-fermion coupling from the 't Hooft interaction `sixferm`. If 0.0 is given for any of the values, then the default is used ( $\Lambda = 602.3/(\hbar c)$ ,  $G = 1.835/\Lambda^2$ ,  $K = 12.36/\Lambda^5$ ).

The value of the shift in the bag constant `B0` is automatically calculated to ensure that the energy density and the pressure of the vacuum are zero. The functions `set_quarks()` and `set_thermo()` can be used before hand to specify the **quark** and **thermo** objects.

14.26.2.2 `virtual int nambu_l_eos::calc_p ( quark &u, quark &d, quark &s, thermo &th ) [virtual]`

This function automatically solves the gap equations

Reimplemented from `quark_eos`.

14.26.2.3 `virtual int nambu_l_eos::calc_temp_p ( quark &u, quark &d, quark &s, double T, thermo &th ) [virtual]`

This function automatically solves the gap equations

Reimplemented from `quark_eos`.

14.26.2.4 `int nambu_l_eos::gapfunms ( size_t nv, const ovector_base &x, ovector_base &y )`

The function utilizes the **quark** objects which can be specified in `set_quarks()` and the **thermo** object which can be specified in `eos::set_thermo()`.

14.26.2.5 int nambuyl\_eos::gapfunqq ( size\_t *nv*, const ovector\_base & *x*, ovector\_base & *y* )

The function utilizes the **quark** objects which can be specified in [set\\_quarks\(\)](#) and the **thermo** object which can be specified in [eos::set\\_thermo\(\)](#).

14.26.2.6 int nambuyl\_eos::gapfunmsT ( size\_t *nv*, const ovector\_base & *x*, ovector\_base & *y* )

The function utilizes the **quark** objects which can be specified in [set\\_quarks\(\)](#) and the **thermo** object which can be specified in [eos::set\\_thermo\(\)](#).

14.26.2.7 int nambuyl\_eos::gapfunqqT ( size\_t *nv*, const ovector\_base & *x*, ovector\_base & *y* )

The function utilizes the **quark** objects which can be specified in [set\\_quarks\(\)](#) and the **thermo** object which can be specified in [eos::set\\_thermo\(\)](#).

14.26.2.8 int nambuyl\_eos::set\_quarks ( quark & *u*, quark & *d*, quark & *s* )

The quark objects are used in [gapfunms\(\)](#), [gapfunqq\(\)](#), [gapfunmsT\(\)](#), [gapfunqqT\(\)](#), and [B0fun\(\)](#).

### 14.26.3 Field Documentation

14.26.3.1 double nambuyl\_eos::limit

[limit](#) is used for the finite temperature integrals to ensure that no numbers larger than  $\exp(\text{limit})$  or smaller than  $\exp(-\text{limit})$  are avoided. (Default: 20)

Definition at line 156 of file nambuyl\_eos.h.

14.26.3.2 bool nambuyl\_eos::fromqq

If this is false, then computations are performed using the effective masses as inputs

Definition at line 163 of file nambuyl\_eos.h.

The documentation for this class was generated from the following file:

- nambuyl\_eos.h

## 14.27 nambuyl\_eos::njtp\_s Struct Reference

A structure for passing parameters to the integrands.

```
#include <nambuyl_eos.h>
```

### 14.27.1 Detailed Description

Definition at line 290 of file nambuyl\_eos.h.

#### Data Fields

- double **ms**
- double **m**
- double **mu**
- double **temper**
- double **limit**

The documentation for this struct was generated from the following file:

- `nambuyl_eos.h`

## 14.28 nse\_eos Class Reference

Equation of state for nuclei in statistical equilibrium.

```
#include <nse_eos.h>
```

### 14.28.1 Detailed Description

This class computes the composition of matter in nuclear statistical equilibrium. The chemical potential of a nucleus X with proton number  $Z_X$  and neutron number  $N_X$  is given by

$$\mu_X = N\mu_n + Z\mu_p - E_{\text{bind},X}$$

where  $\mu_n$  and  $\mu_p$  are the neutron and proton chemical potentials and  $E_{\text{bind},X}$  is the binding energy of the nucleus.

The baryon number density and electron fraction are then given by

$$n_B = n_X(N_X + Z_X) \quad Y_e n_B = n_X Z_X$$

where  $n_X$  is the number density which is determined from the chemical potential above.

This implicitly assumes that the nuclei are non-interacting.

**Idea for Future** Right now `calc_density()` needs a very good guess. This could be fixed, probably by solving for the  $\log(\mu/T)$  instead of  $\mu$ .

Definition at line 62 of file `nse_eos.h`.

### Public Member Functions

- `int calc_mu` (double *mun*, double *mup*, double *T*, double &*nb*, double &*Ye*, **thermo** &*th*, **nuclear\_dist** &*nd*)  
*Calculate the equation of state as a function of the chemical potentials.*
- `int calc_density` (double *nb*, double *Ye*, double *T*, double &*mun*, double &*mup*, **thermo** &*th*, **nuclear\_dist** &*nd*)  
*Calculate the equation of state as a function of the densities.*
- `int set_mroot` (**mroot**< **mm\_funct**<> > &*rp*)  
*Set the solver for use in computing the chemical potentials.*

### Data Fields

- `gsl_mroot_hybrids`< **mm\_funct**<> > `def_root`  
*Default solver.*

### 14.28.2 Member Function Documentation

14.28.2.1 `int nse_eos::calc_mu` ( double *mun*, double *mup*, double *T*, double &*nb*, double &*Ye*, **thermo** &*th*, **nuclear\_dist** &*nd* )

Given *mun*, *mup* and *T*, this computes the composition (the individual densities are stored in the distribution *nd*) the baryon number density *nb*, and the electron fraction *Ye*.

This function does not use the solver.

14.28.2.2 `int nse_eos::calc_density ( double nb, double Ye, double T, double & mun, double & mup, thermo & th, nuclear_dist & nd )`

Given the baryon number density `nb`, and the electron fraction `Ye` and the temperature `T`, this computes the composition (the individual densities are stored in the distribution `nd`) and the chemical potentials are given in `mun` and `mup`.

This function uses the solver to self-consistently compute the chemical potentials.

The documentation for this class was generated from the following file:

- `nse_eos.h`

## 14.29 `rmf_nucleus::odparms` Struct Reference

A convenient struct for the solution of the Dirac equations.

```
#include <rmf_nucleus.h>
```

### 14.29.1 Detailed Description

Definition at line 244 of file `rmf_nucleus.h`.

#### Data Fields

- double `eigen`  
*Eigenvalue.*
- double `kappa`  
*Quantum number  $\kappa$ .*
- `umatrix` \* `fields`  
*The meson fields.*
- `uvector` \* `varr`  
*Desc.*

The documentation for this struct was generated from the following file:

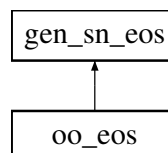
- `rmf_nucleus.h`

## 14.30 `oo_eos` Class Reference

The EOS tables from O'Connor and Ott.

```
#include <gen_sn_eos.h>
```

Inheritance diagram for `oo_eos`:



### 14.30.1 Detailed Description

This class reads the HDF5 EOS tables generated by E. O'Connor and C. Ott in [OConnor10](#). The tables are available from

<http://stellarcollapse.org/equationofstate>

and are available under a creative commons attribution-noncommercial-share alike license. This O<sub>2</sub>scl code to read those tables is licensed (along with all O<sub>2</sub>scl code) under the GPLv3 license (with permission from Evan O'Connor).

The original README file from O'Connor and Ott's EOSdriver code is available in the O<sub>2</sub>scl distribution in `doc/o2scl/eos/extras/scol` \_README and is reproduced below

```
! #####
!
! Copyright C. D. Ott and E. O'Connor, July 2009
!
!
!
! EOS constructed on the basis of the Shen et al. 1998
! relativistic-mean field nuclear EOS table. Electrons (fully
! general, based on TimmesEOS) and Photons added.
!
! Original Shen EOS table extent:
!
! Density: 10^5.1 to 10^15.4 g/cm^3
! Temperature: 0.1 to 100 MeV
! Ye: 0.01 to 0.56
!
! Table extent of current table
! (myshen_test_220r_180t_50y_extT_20090312.h5)
!
! Density: 10^3 to 10^15.36 g/cm^3
! Temperature: 0.01 to 250 MeV
! Ye: 0.015 to 0.56 MeV
!
! This bigger table is realized by extending the
! original Shen table in multiple ways in multiple
! directions:
!
! (a) density:
! Match of pure ideal gas of Ni56 + electrons/positrons +
! photons at densities below 10^7 g/ccm -- at this density
! pressures, energies and entropies match okayish with the
! values in the Shen table. The compositions (A,Z,xh,xa,xp,xn)
! are kept constant in the low-density region and n and p
! chemical potentials are set to 0.
! -- ideally, at low densities, a full NSE EOS with nuclear
! reaction network (at low T) should be stiched onto the
! Shen; working on that, but not yet ready.
!
! (b) temperature (extrapolation):
! At high density: linear extrapolation of everything in T to lower
! temperatures and higher temperatures. At low densities (below 10^7
! g/ccm), ideal gas of Ni56 + electrons/positrons + photons.
!
!
!
! Variable          Units          Description
! pointsrho         dimensionless    number of table points
!                   in log10(rho)
! pointstemp         dimensionless    number of table points
!                   in log10(temp)
! pointsye           dimensionless    number of table point in Y_e
! logrho             log10(rho)       index variable rho
! logtemp            log10(MeV)       index variable temperature
! ye                 number fraction   index variable electron fraction
! Abar               A                average heavy nucleus A
! Zbar               Z                average heavy nucleus Z
! Xa                 number fraction   alpha particle number frac
! Xh                 number fraction   heavy nucleos number frac
! Xn                 number fraction   neutron number frac
! Xn                 number fraction   proton number frac
! cs2                cm^2/s^2         speed of sound squared
! dedt              erg/g/K          C_v
! dpderho            dyn g/cm^2/erg   dp/deps|rho
! dpdrhoe            dyn cm^3/cm^2/g   dp/drho|eps
```

```

! energy_shift      erg/g      Energy shift for table storage
! entropy           k_B/baryon  specific entropy
! gamma             dimensionless Gamma_1
! logenergy         log10(erg/g) specific internal energy
! logpress          log10(dyn/cm^2) pressure
! mu_e             MeV/baryon  electron chemical potential
!                  INCLUDING electron rest mass
! mu_p             MeV/baryon  proton chemical potential
! mu_n             MeV/baryon  neutron chemical potential
! muhat            MeV/baryon  mu_n - mu_p
! munu             MeV/baryon  mu_e - mun + mu_p
!
!
! * energy shift:
!
! In some regions the negative nuclear binding energy
! is larger in magnitude than the thermal/excitation energy. In this
! case the specific internal energy (eps) becomes negative. To allow for
! storage and interpolation of eps in logarithmic fashion, the energy
! is shifted up by an energy shift specified in the variable "energy_shift".
! This energy shift is handled internally in the EOS routines.

```

Definition at line 500 of file gen\_sn\_eos.h.

#### Public Member Functions

- virtual void [load](#) (std::string fname, size\_t mode)  
*Load table from filename fname.*
- virtual void [load](#) (std::string fname)  
*Load table from filename fname.*
- virtual void [beta\\_eq\\_T0](#) (size\_t i, double &nb, double &E\_beta, double &P\_beta, double &Ye\_beta, double &Z\_beta, double &A\_beta)  
*Compute properties of matter in beta equilibrium at zero temperature at a baryon density grid point.*

#### Data Fields

- **tensor\_grid** & [cs2](#)  
*Speed of sound in  $\text{cm}^2/\text{s}^2$ .*
- **tensor\_grid** & [dedt](#)  
 *$C_V$  in erg/g/K.*
- **tensor\_grid** & [dpderho](#)  
*dpderho in  $\text{dyn} \cdot \text{g}/\text{cm}^2/\text{erg}$*
- **tensor\_grid** & [dpdrhoe](#)  
*dpdrhoe in  $\text{dyn cm}^3/\text{cm}^2/\text{g}$*
- **tensor\_grid** & [gamma](#)  
*Gamma.*
- **tensor\_grid** & [mu\\_e](#)  
*Electron chemical potential per baryon including rest mass.*
- **tensor\_grid** & [muhat](#)  
*mun - mup*
- **tensor\_grid** & [munu](#)  
*mue - mun + mup*
- **ovector** [rho](#)  
*The original mass density grid from the table in  $\text{g}/\text{cm}^3$ .*
- double [energy\\_shift](#)  
*Energy shift for table storage in erg/g.*

## Static Public Attributes

- static const size\_t `ls_mode` = 0  
*Use the J. Lattimer et al. method for handling the chemical potentials.*
- static const size\_t `stos_mode` = 1  
*Use the H. Shen et al. method for handling the chemical potentials.*

## 14.30.2 Member Function Documentation

14.30.2.1 virtual void `oo_eos::beta_eq_T0` ( size\_t *i*, double & *nb*, double & *E\_beta*, double & *P\_beta*, double & *Ye\_beta*, double & *Z\_beta*, double & *A\_beta* ) [inline, virtual]

This EOS table doesn't have T=0 results, so we extrapolate from the two low-temperature grid points.

Implements `gen_sn_eos`.

Definition at line 556 of file `gen_sn_eos.h`.

The documentation for this class was generated from the following file:

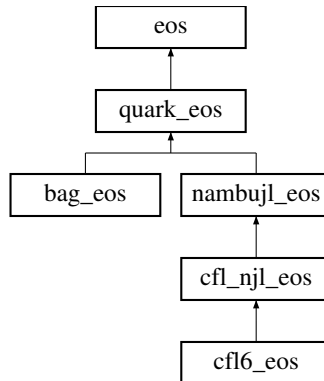
- `gen_sn_eos.h`

## 14.31 quark\_eos Class Reference

Quark matter equation of state base.

```
#include <quark_eos.h>
```

Inheritance diagram for `quark_eos`:



## 14.31.1 Detailed Description

Definition at line 39 of file `quark_eos.h`.

## Public Member Functions

- virtual int `calc_p` (quark &u, quark &d, quark &s, thermo &th)  
*Calculate equation of state as a function of chemical potentials.*
- virtual int `calc_e` (quark &u, quark &d, quark &s, thermo &th)  
*Calculate equation of state as a function of density.*
- virtual int `calc_temp_p` (quark &u, quark &d, quark &s, double temper, thermo &th)  
*Calculate equation of state as a function of chemical potentials.*

- virtual int [calc\\_temp\\_e](#) (**quark** &u, **quark** &d, **quark** &s, double temper, **thermo** &th)  
*Calculate equation of state as a function of density.*
- virtual const char \* [type](#) ()  
*Return string denoting type ("quark\_eos")*

#### Data Fields

- fermion\_eval\_thermo \* [fet](#)  
*Desc.*
- **eff\_fermion** [def\\_fet](#)  
*Desc.*

The documentation for this class was generated from the following file:

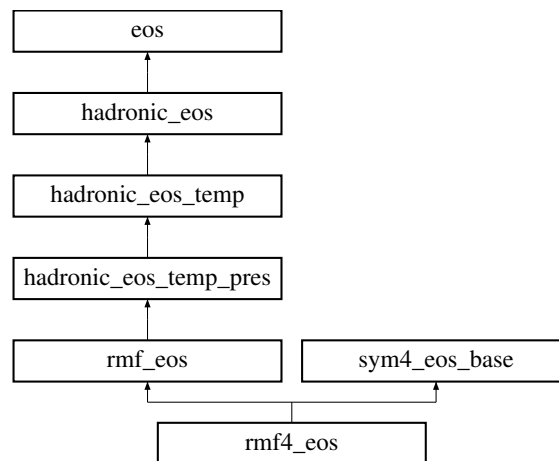
- quark\_eos.h

## 14.32 **rmf4\_eos** Class Reference

A version of [rmf\\_eos](#) to separate potential and kinetic contributions.

```
#include <sym4_eos.h>
```

Inheritance diagram for **rmf4\_eos**:



### 14.32.1 Detailed Description

#### References:

Created for [Steiner06](#).

Definition at line 103 of file sym4\_eos.h.

#### Public Member Functions

- virtual int [calc\\_e\\_sep](#) (**fermion** &ne, **fermion** &pr, double &ed\_kin, double &ed\_pot, double &mu\_n\_kin, double &mu\_p\_kin, double &mu\_n\_pot, double &mu\_p\_pot)  
*Compute the potential and kinetic parts separately.*



The documentation for this class was generated from the following file:

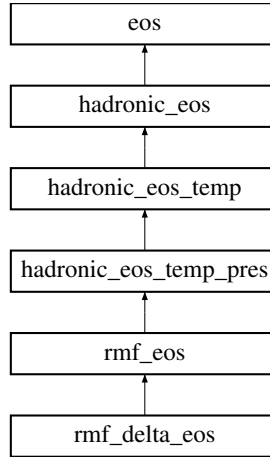
- sym4\_eos.h

### 14.33 rmf\_delta\_eos Class Reference

Field-theoretical EOS with scalar-isovector meson,  $\delta$ .

```
#include <rmf_delta_eos.h>
```

Inheritance diagram for rmf\_delta\_eos:



#### 14.33.1 Detailed Description

This essentially follows the notation in [Kubis97](#), except that our definitions of  $\mathfrak{b}$  and  $\mathfrak{c}$  follow their  $\bar{b}$  and  $\bar{c}$ , respectively.

Also discussed in [Gaitanos04](#), where they take  $m_\delta = 980$  MeV.

The full Lagrangian is:

$$\mathcal{L} = \mathcal{L}_{Dirac} + \mathcal{L}_\sigma + \mathcal{L}_\omega + \mathcal{L}_\rho + \mathcal{L}_\delta$$

$$\begin{aligned}
 \mathcal{L}_{Dirac} &= \bar{\Psi} \left[ i \not{\partial} - g_\omega \not{\omega} - \frac{g_\rho}{2} \vec{\not{\rho}} \cdot \vec{\tau} - M + g_\sigma \sigma - \frac{e}{2} (1 + \tau_3) A_\mu \right] \Psi \\
 \mathcal{L}_\sigma &= \frac{1}{2} (\partial_\mu \sigma)^2 - \frac{1}{2} m_\sigma^2 \sigma^2 - \frac{bM}{3} (g_\sigma \sigma)^3 - \frac{c}{4} (g_\sigma \sigma)^4 \\
 \mathcal{L}_\omega &= -\frac{1}{4} f_{\mu\nu} f^{\mu\nu} + \frac{1}{2} m_\omega^2 \omega^\mu \omega_\mu + \frac{\zeta}{24} g_\omega^4 (\omega^\mu \omega_\mu)^2 \\
 \mathcal{L}_\rho &= -\frac{1}{4} \vec{B}_{\mu\nu} \cdot \vec{B}^{\mu\nu} + \frac{1}{2} m_\rho^2 \vec{\rho}^\mu \cdot \vec{\rho}_\mu + \frac{\xi}{24} g_\rho^4 (\vec{\rho}^\mu \cdot \vec{\rho}_\mu)^2 + g_\rho^2 f(\sigma, \omega) \vec{\rho}^\mu \cdot \vec{\rho}_\mu
 \end{aligned}$$

where the additional terms are

$$\mathcal{L}_\delta = \bar{\Psi} \left( g_\delta \vec{\delta} \cdot \vec{\tau} \right) \Psi + \frac{1}{2} (\partial_\mu \vec{\delta})^2 - \frac{1}{2} m_\delta^2 \vec{\delta}^2$$

The new field equation for the delta meson is

$$m_\delta^2 \vec{\delta} = g_\delta (n_{s,p} - n_{s,n})$$

**Idea for Future** Finish the finite temperature EOS

Definition at line 92 of file `rmf_delta_eos.h`.

#### Public Member Functions

- virtual int `calc_e` (**fermion** &ne, **fermion** &pr, **thermo** &lth)  
*Equation of state as a function of density.*
- virtual int `calc_p` (**fermion** &neu, **fermion** &p, double sig, double ome, double rho, double delta, double &f1, double &f2, double &f3, double &f4, **thermo** &th)  
*Equation of state as a function of chemical potentials.*
- int `calc_temp_p` (**fermion** &ne, **fermion** &pr, double temper, double sig, double ome, double lrho, double delta, double &f1, double &f2, double &f3, double &f4, **thermo** &lth)  
*Finite temperature (unfinished)*
- virtual int `set_fields` (double sig, double ome, double lrho, double delta)  
*Set a guess for the fields for the next call to `calc_e()`, `calc_p()`, or `saturation()`*
- virtual int `saturation` ()  
*Calculate saturation properties for nuclear matter at the saturation density.*

#### Data Fields

- double `md`  
*The mass of the scalar-isovector field.*
- double `cd`  
*The coupling of the scalar-isovector field to the nucleons.*
- double `del`  
*The value of the scalar-isovector field.*

#### Protected Member Functions

- virtual int `calc_e_solve_fun` (size\_t nv, const **ovector\_base** &ex, **ovector\_base** &ey)  
*The function for `calc_e()`*
- virtual int `zero_pressure` (size\_t nv, const **ovector\_base** &ex, **ovector\_base** &ey)  
*Compute matter at zero pressure (for `saturation()`)*

#### Private Member Functions

- virtual int `set_fields` (double sig, double ome, double lrho)  
*Forbid setting the guesses to the fields unless all four fields are specified.*

### 14.33.2 Member Function Documentation

#### 14.33.2.1 virtual int `rmf_delta_eos::saturation` ( ) [virtual]

This requires initial guesses to the chemical potentials, etc.

Reimplemented from `rmf_eos`.

The documentation for this class was generated from the following file:

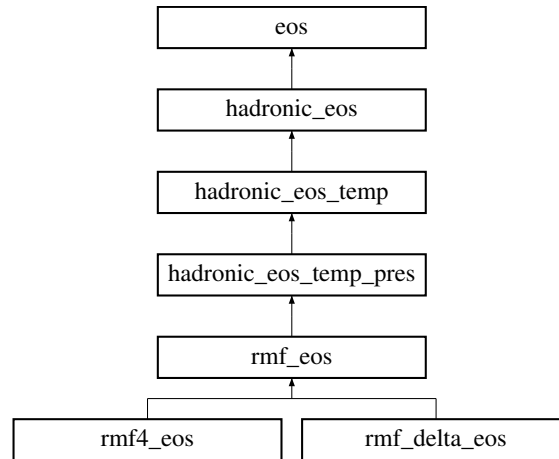
- `rmf_delta_eos.h`

## 14.34 `rmf_eos` Class Reference

Relativistic mean field theory EOS.

```
#include <rmf_eos.h>
```

Inheritance diagram for `rmf_eos`:



### 14.34.1 Detailed Description

This class computes the properties of nucleonic matter using a mean-field approximation to a field-theoretical model.

Before sending neutrons and protons to these member functions, the masses should be set to their vacuum values and the degeneracy factor should be 2. If an internal model is used (using `load()`), then the neutron and proton masses should be set to `mnuc`.

The expressions for the energy densities are often simplified in the literature using the field equations. These expressions are not used in this code since they are only applicable in infinite matter where the field equations hold, and are not suitable for use in applications (such as to finite nuclei in `rmf_nucleus`) where the spatial derivatives of the fields are non-zero. Notice that in the proper expressions for the energy density the similarity between terms in the pressure up to a sign. This procedure allows one to verify the thermodynamic identity even if the field equations are not solved and allows the user to add gradient terms to the energy density and pressure.

#### Note

Since this EOS uses the effective masses and chemical potentials in the fermion class, the values of `part::non_interacting` for neutrons and protons are set to false in many of the functions.

- Todo**
- The functions `fcomp_fields()`, `fkprime_fields()`, and `fesym_fields()` are not quite correct if the neutron and proton masses are different. For this reason, they are currently unused by `saturation()`.
  - The `fix_saturation()` and `calc_cr()` functions use `mnuc`, and should be modified to allow different neutron and proton masses.
  - Check the formulas in the "Background" section
  - There are two `calc_e()` functions that solve. One is specially designed to work without a good initial guess. Possibly the other `calc_e()` function should be similarly designed?
  - Make sure that this class properly handles particles for which `inc_rest_mass` is true/false
  - The error handler is called sometimes when `calc_e()` is used to compute pure neutron matter. This should be fixed.
  - Decide whether to throw an error at [Ref. 1].
  - Put the `err_nonconv` system into `calc_p()`, `calc_temp_e()` and `fix_saturation()`, etc.

**Idea for Future**

- It might be nice to remove explicit reference to the meson masses in functions which only compute nuclear matter since they are unnecessary. This might, however, demand redefining some of the couplings.
- Fix `calc_p()` to be better at guessing
- The number of couplings is getting large, maybe new organization is required.
- Overload `hadronic_eos::fcomp()` with an exact version

**Background**

The full Lagragian is:

$$\mathcal{L} = \mathcal{L}_{Dirac} + \mathcal{L}_{\sigma} + \mathcal{L}_{\omega} + \mathcal{L}_{\rho}$$

$$\begin{aligned}\mathcal{L}_{Dirac} &= \bar{\Psi} \left[ i\partial\!\!\!/ - g_{\omega}\omega\!\!\!/ - \frac{g_{\rho}}{2}\vec{\rho}\vec{\tau} - M + g_{\sigma}\sigma - \frac{e}{2}(1 + \tau_3)A_{\mu} \right] \Psi \\ \mathcal{L}_{\sigma} &= \frac{1}{2}(\partial_{\mu}\sigma)^2 - \frac{1}{2}m_{\sigma}^2\sigma^2 - \frac{bM}{3}(g_{\sigma}\sigma)^3 - \frac{c}{4}(g_{\sigma}\sigma)^4 \\ \mathcal{L}_{\omega} &= -\frac{1}{4}f_{\mu\nu}f^{\mu\nu} + \frac{1}{2}m_{\omega}^2\omega^{\mu}\omega_{\mu} + \frac{\zeta}{24}g_{\omega}^4(\omega^{\mu}\omega_{\mu})^2 \\ \mathcal{L}_{\rho} &= -\frac{1}{4}\vec{B}_{\mu\nu}\cdot\vec{B}^{\mu\nu} + \frac{1}{2}m_{\rho}^2\vec{\rho}^{\mu}\cdot\vec{\rho}_{\mu} + \frac{\xi}{24}g_{\rho}^4(\vec{\rho}^{\mu})\cdot\vec{\rho}_{\mu} + g_{\rho}^2f(\sigma,\omega)\vec{\rho}^{\mu}\cdot\vec{\rho}_{\mu}\end{aligned}$$

The couplings `cs`, `cw`, and `cr` are related to  $g_{\sigma}$ ,  $g_{\omega}$  and  $g_{\rho}$  above by

$$g_s = c_{\sigma}m_{\sigma} \quad g_w = c_{\omega}m_{\omega} \quad \text{and} \quad g_r = c_{\rho}m_{\rho}$$

The coefficients  $b$  and  $c$  are related to the somewhat standard  $\kappa$  and  $\lambda$  by:

$$\kappa = 2Mb \quad \lambda = 6c;$$

The function  $f$  is the coefficient of  $g_{\rho}^2\rho^2$   $f(\sigma,\omega) = b_1\omega^2 + b_2\omega^4 + b_3\omega^6 + a_1\sigma + a_2\sigma^2 + a_3\sigma^3 + a_4\sigma^4 + a_5\sigma^5 + a_6\sigma^6$  where the notation from [Horowitz01](#) is:  $f(\sigma,\omega) = \lambda_4g_s^2\sigma^2 + \lambda_wg_w^2\omega^2$  This implies  $b_1 = \lambda_wg_w^2$  and  $a_2 = \lambda_4g_s^2$

The couplings, `cs`, `cw`, and `cr` all have units of fm, and the couplings `b`, `c`, `zeta` and `xi` are unitless. The additional couplings from [Steiner05b](#),  $a_i$  have units of fm<sup>(i-2)</sup> and the couplings  $b_j$  have units of fm<sup>(2\*j-2)</sup>.

The field equations are:

$$\begin{aligned}0 &= m_{\sigma}^2\sigma - g_{\sigma}(n_{sn} + n_{sp}) + bMg_{\sigma}^3\sigma^2 + cg_{\sigma}^4\sigma^3 - g_{\rho}^2\rho^2\frac{\partial f}{\partial\sigma} \\ 0 &= m_{\omega}^2\omega - g_{\omega}(n_n + n_p) + \frac{\zeta}{6}g_{\omega}^4\omega^3 + g_{\rho}^2\rho^2\frac{\partial f}{\partial\omega} \\ 0 &= m_{\rho}^2\rho + \frac{1}{2}g_{\rho}(n_n - n_p) + 2g_{\rho}^2\rho f + \frac{\xi}{6}g_{\rho}^4\rho^3\end{aligned}$$

When the variable `zm_mode` is true, the effective mass is fixed using the approach of [Zimanyi90](#).

Defining

$$U(\sigma) = \frac{1}{2}m_{\sigma}^2\sigma^2 + \frac{bM}{3}(g_{\sigma}\sigma)^3 + \frac{c}{4}(g_{\sigma}\sigma)^4,$$

the binding energy per particle in symmetric matter at equilibrium is given by

$$\frac{E}{A} = \frac{1}{n_0} \left[ U(\sigma_0) + \frac{1}{2}m_{\omega}\omega_0^2 + \frac{\zeta}{8}(g_{\omega}\omega_0)^4 + \frac{2}{\pi^2} \int_0^{k_F} dk k^2 \sqrt{k^2 + M^{*2}} \right].$$

where the Dirac effective mass is  $M^* = M - g_\sigma \sigma_0$ . The compressibility is given by

$$K = 9 \frac{g_\omega^2}{m_\omega^2} n_0 + 3 \frac{k_F^2}{E_F^*} - 9 n_0 \frac{M^{*2}}{E_F^{*2}} \left[ \left( \frac{1}{g_\sigma^2} \frac{\partial^2}{\partial \sigma_0^2} + \frac{3}{g_\sigma M^*} \frac{\partial}{\partial \sigma_0} \right) U(\sigma_0) - 3 \frac{n_0}{E_F^*} \right]^{-1}.$$

The symmetry energy of bulk matter is given by

$$E_{\text{sym}} = \frac{k_F^2}{6E_F^*} + \frac{n}{8 \left( g_\rho^2/m_\rho^2 + 2f(\sigma_0, \omega_0) \right)}$$

In the above equations, the subscript 0 denotes the mean field values of  $\sigma$  and  $\omega$ . For the case  $f = 0$ , the symmetry energy varies linearly with the density at large densities. The function  $f$  permits variations in the density dependence of the symmetry energy above nuclear matter density.

See also [Muller96](#), [Zimanyi90](#), and [Steiner05b](#).

Definition at line 221 of file `rmf_eos.h`.

### Public Member Functions

- virtual const char \* [type](#) ()  
*Return string denoting type ("rmf\_eos")*
- int [check\\_naturalness](#) (rmf\_eos &re)  
*Set the coefficients of a rmf\_eos object to their limits from naturalness.*
- int [naturalness\\_limits](#) (double value, rmf\_eos &re)  
*Provide the maximum values of the couplings assuming a limit on naturalness.*

### Compute EOS

- virtual int [calc\\_e](#) (fermion &ne, fermion &pr, thermo &lth)  
*Equation of state as a function of density.*
- virtual int [calc\\_e\\_fields](#) (fermion &ne, fermion &pr, thermo &lth, double &sig, double &ome, double &rho)  
*Equation of state as a function of density returning the meson fields.*
- virtual int [calc\\_p](#) (fermion &ne, fermion &pr, thermo &lth)  
*Equation of state as a function of chemical potential.*
- virtual int [calc\\_eq\\_p](#) (fermion &neu, fermion &p, double sig, double ome, double rho, double &f1, double &f2, double &f3, thermo &th)  
*Equation of state and meson field equations as a function of chemical potentials.*
- virtual int [calc\\_eq\\_temp\\_p](#) (fermion &ne, fermion &pr, double temper, double sig, double ome, double rho, double &f1, double &f2, double &f3, thermo &th)  
*Equation of state and meson field equations as a function of chemical potentials at finite temperature.*
- virtual int [calc\\_temp\\_p](#) (fermion &ne, fermion &pr, double T, thermo &lth)  
*Equation of state as a function of chemical potential.*
- int [calc\\_temp\\_e](#) (fermion &ne, fermion &pr, double T, thermo &lth)  
*Equation of state as a function of densities at finite temperature.*

### Saturation properties

- int [fix\\_saturation](#) (double guess\_cs=4.0, double guess\_cw=3.0, double guess\_b=0.001, double guess\_c=-0.001)  
*Calculate cs, cw, cr, b, and c from the saturation properties.*
- virtual int [saturation](#) ()  
*Calculate properties of nuclear matter at the saturation density.*
- double [fesym\\_fields](#) (double sig, double ome, double nb)  
*Calculate symmetry energy assuming the field equations have already been solved.*
- double [fcomp\\_fields](#) (double sig, double ome, double nb)  
*Calculate the compressibility assuming the field equations have already been solved.*
- int [fkprime\\_fields](#) (double sig, double ome, double nb, double &k, double &kprime)  
*Calculate compressibility and kprime assuming the field equations have already been solved.*

## Fields and field equations

- int [field\\_eqs](#) (size\_t nv, const **ovector\_base** &x, **ovector\_base** &y)  
*A function for solving the field equations.*
- int [field\\_eqsT](#) (size\_t nv, const **ovector\_base** &x, **ovector\_base** &y)  
*A function for solving the field equations at finite temperature.*
- virtual int [set\\_fields](#) (double sig, double ome, double lrho)  
*Set a guess for the fields for the next call to [calc\\_e\(\)](#), [calc\\_p\(\)](#), or [saturation\(\)](#)*
- int [get\\_fields](#) (double &sig, double &ome, double &lrho)  
*Return the most recent values of the meson fields.*

## Data Fields

- bool [zm\\_mode](#)  
*Modifies method of calculating effective masses (default false)*
- int [verbose](#)  
*Verbosity parameter.*
- bool [err\\_nonconv](#)  
*If true, throw exceptions when the function [calc\\_e\(\)](#) does not converge (default true)*
- int [last\\_conv](#)  
*The convergence status of the last call to [calc\\_e\(\)](#)*

## Masses

- double [mnuc](#)  
*nucleon mass*
- double [ms](#)  
 *$\sigma$  mass (in fm<sup>-1</sup>)*
- double [mw](#)  
 *$\omega$  mass (in fm<sup>-1</sup>)*
- double [mr](#)  
 *$\rho$  mass (in fm<sup>-1</sup>)*

## Standard couplings (including nonlinear sigma terms)

- double [cs](#)
- double [cw](#)
- double [cr](#)
- double [b](#)
- double [c](#)

## Non-linear terms for omega and rho.

- double [zeta](#)
- double [xi](#)

## Additional isovector couplings

- double [a1](#)
- double [a2](#)
- double [a3](#)
- double [a4](#)
- double [a5](#)
- double [a6](#)
- double [b1](#)
- double [b2](#)
- double [b3](#)

## Protected Member Functions

- int [fix\\_saturation\\_fun](#) (size\_t nv, const **ovector\_base** &x, **ovector\_base** &y)  
*The function for [fix\\_saturation\(\)](#)*
- virtual int [zero\\_pressure](#) (size\_t nv, const **ovector\_base** &ex, **ovector\_base** &ey)

Compute matter at zero pressure (for [saturation\(\)](#))

- virtual int [calc\\_e\\_solve\\_fun](#) (size\_t nv, const **ovector\_base** &ex, **ovector\_base** &ey)  
The function for [calc\\_e\(\)](#)
- virtual int [calc\\_temp\\_e\\_solve\\_fun](#) (size\_t nv, const **ovector\_base** &ex, **ovector\_base** &ey)  
The function for [calc\\_temp\\_e\(\)](#)
- int [calc\\_cr](#) (double sig, double ome, double nb)  
Calculate the *cr* coupling given *sig* and *ome* at the density 'nb'.

#### Protected Attributes

- double [n\\_baryon](#)  
Desc.
- double [n\\_charge](#)  
Temporary charge density.
- double [fe\\_temp](#)  
Temperature for solving field equations at finite temperature.
- bool [ce\\_neut\\_matter](#)  
For [calc\\_e\(\)](#), if true, then solve for neutron matter.
- bool [ce\\_prot\\_matter](#)  
For [calc\\_e\(\)](#), if true, then solve for proton matter.
- bool [guess\\_set](#)  
True if a guess for the fields has been given.
- **mroot**< **mm\_funct**<> > \* [sat\\_mroot](#)  
The solver to compute saturation properties.
- double [ce\\_temp](#)  
Temperature storage for [calc\\_temp\\_e\(\)](#)

#### The meson fields

- double **sigma**
- double **omega**
- double **rho**

#### Solver

- **gsl\_mroot\_hybrids**< **mm\_funct**<> > [def\\_sat\\_mroot](#)  
The default solver for calculating the saturation density.
- virtual int [set\\_sat\\_mroot](#) (**mroot**< **mm\_funct**<> > &mr) *Set class mroot object for use calculating saturation density.*

### 14.34.2 Member Function Documentation

14.34.2.1 virtual int **rmf\_eos::calc\_e** ( fermion & *ne*, fermion & *pr*, thermo & *lth* ) [virtual]

Initial guesses for the chemical potentials are taken from the user-given values. Initial guesses for the fields can be set by [set\\_fields\(\)](#), or default values will be used. After the call to [calc\\_e\(\)](#), the final values of the fields can be accessed through [get\\_fields\(\)](#).

This is a little more robust than the standard version in the parent [hadronic\\_eos](#).

**Idea for Future** Improve the operation of this function when the proton density is zero.

Reimplemented from [hadronic\\_eos\\_temp\\_pres](#).

Reimplemented in [rmf\\_delta\\_eos](#).

14.34.2.2 `virtual int rmf_eos::calc_e_fields ( fermion & ne, fermion & pr, thermo & lth, double & sig, double & ome, double & rho )`  
[virtual]

**Idea for Future** Improve the operation of this function when the proton density is zero.

14.34.2.3 `virtual int rmf_eos::calc_p ( fermion & ne, fermion & pr, thermo & lth )` [virtual]

Solves for the field equations automatically.

#### Note

This may not be too robust. Fix?

Implements [hadronic\\_eos\\_temp\\_pres](#).

14.34.2.4 `virtual int rmf_eos::calc_eq_p ( fermion & neu, fermion & p, double sig, double ome, double rho, double & f1, double & f2, double & f3, thermo & th )` [virtual]

This calculates the pressure and energy density as a function of  $\mu_n, \mu_p, \sigma, \omega, \rho$ . When the field equations have been solved, f1, f2, and f3 are all zero.

The thermodynamic identity is satisfied even when the field equations are not solved.

**Idea for Future** Probably best to have f1, f2, and f3 scaled in some sensible way, i.e. scaled to the fields?

14.34.2.5 `virtual int rmf_eos::calc_eq_temp_p ( fermion & ne, fermion & pr, double temper, double sig, double ome, double rho, double & f1, double & f2, double & f3, thermo & th )` [virtual]

Analogous to [calc\\_eq\\_p\(\)](#) except at finite temperature.

14.34.2.6 `virtual int rmf_eos::calc_temp_p ( fermion & ne, fermion & pr, double T, thermo & lth )` [virtual]

Solves for the field equations automatically.

Implements [hadronic\\_eos\\_temp\\_pres](#).

14.34.2.7 `int rmf_eos::fix_saturation ( double guess_cs = 4.0, double guess_cw = 3.0, double guess_b = 0.001, double guess_c = -0.001 )`

Note that the meson masses and [mnuc](#) must be specified before calling this function.

This function does not give correct results when `bool zm_mode` is true.

`guess_cs`, `guess_cw`, `guess_b`, and `guess_c` are initial guesses for `cs`, `cw`, `b`, and `c` respectively.

- Todo**
- Fix this for `zm_mode=true`
  - Ensure solver is more robust

14.34.2.8 `virtual int rmf_eos::saturation ( )` [virtual]

This function first constructs an initial guess, increasing the chemical potentials if required to ensure the neutron and proton densities are finite, and then uses [rmf\\_eos::sat\\_mroot](#) to solve the field equations and ensure that the neutron and proton densities are equal and the pressure is zero. The quantities [hadronic\\_eos::n0](#), [hadronic\\_eos::eo](#), and [hadronic\\_eos::msom](#) can be computed directly, and the compressibility, the skewness, and the symmetry energy are computed using the functions [fkprime\\_fields\(\)](#) and [fesym\\_fields\(\)](#). This function overrides the generic version in [hadronic\\_eos](#).

If `verbose` is greater than zero, then then this function reports details on the initial iterations to get the initial guess for the solver.

Reimplemented from [hadronic\\_eos](#).

Reimplemented in [rmf\\_delta\\_eos](#).



14.34.2.9 double rmf\_eos::fesym\_fields ( double *sig*, double *ome*, double *nb* )

This may only work at saturation density. Used by [saturation\(\)](#).

14.34.2.10 double rmf\_eos::fcomp\_fields ( double *sig*, double *ome*, double *nb* )

This may only work at saturation density.

14.34.2.11 int rmf\_eos::fkprime\_fields ( double *sig*, double *ome*, double *nb*, double & *k*, double & *kprime* )

This may only work at saturation density. Used by [saturation\(\)](#).

**Todo** Does this work? Fix [fkprime\\_fields\(\)](#) if it does not.

14.34.2.12 int rmf\_eos::field\_eqs ( size\_t *nv*, const ovector\_base & *x*, ovector\_base & *y* )

*x*[0], *x*[1], and *x*[2] should be set to  $\sigma$ ,  $\omega$ , and  $\rho$  on input (in fm<sup>-1</sup>) and on exit, *y*[0], *y*[1] and *y*[2] contain the field equations and are zero when the field equations have been solved. The *pa* parameter is ignored.

14.34.2.13 int rmf\_eos::field\_eqsT ( size\_t *nv*, const ovector\_base & *x*, ovector\_base & *y* )

*x*[0], *x*[1], and *x*[2] should be set to  $\sigma$ ,  $\omega$ , and  $\rho$  on input (in fm<sup>-1</sup>) and on exit, *y*[0], *y*[1] and *y*[2] contain the field equations and are zero when the field equations have been solved. The *pa* parameter is ignored.

14.34.2.14 int rmf\_eos::get\_fields ( double & *sig*, double & *ome*, double & *lrho* ) [inline]

This returns the most recent values of the meson fields set by a call to [saturation\(\)](#), [calc\\_e\(\)](#), or [calc\\_p\(fermion &, fermion &, thermo &\)](#).

Definition at line 488 of file [rmf\\_eos.h](#).

14.34.2.15 int rmf\_eos::check\_naturalness ( rmf\_eos & *re* ) [inline]

As given in [Muller96](#).

The definition of the vector-isovector field and coupling matches what is done here. Compare the Lagrangian above with Eq. 10 from the reference.

The following couplings should all be of the same size:

$$\frac{1}{2c_s^2 M^2}, \frac{1}{2c_v^2 M^2}, \frac{1}{8c_\rho^2 M^2}, \text{ and } \frac{\bar{a}_{ijk} M^{i+2j+2k-4}}{2^{2k}}$$

which are equivalent to

$$\frac{m_s^2}{2g_s^2 M^2}, \frac{m_v^2}{2g_v^2 M^2}, \frac{m_\rho^2}{8g_\rho^2 M^2}, \text{ and } \frac{a_{ijk} M^{i+2j+2k-4}}{g_s^i g_v^{2j} g_\rho^{2k} 2^{2k}}$$

The connection the  $a_{ijk}$  's and the coefficients that are used here is

$$\begin{aligned} \frac{bM}{3} g_\sigma^3 \sigma^3 &= a_{300} \sigma^3 \\ \frac{c}{4} g_\sigma^4 \sigma^4 &= a_{400} \sigma^4 \\ \frac{\zeta}{24} g_\omega^4 \omega^4 &= a_{020} \omega^4 \\ \frac{\xi}{24} g_\rho^4 \rho^4 &= a_{002} \rho^4 \\ b_1 g_\rho^2 \omega^2 \rho^2 &= a_{011} \omega^2 \rho^2 \end{aligned}$$

$$\begin{aligned}
b_2 g_\rho^2 \omega^4 \rho^2 &= a_{021} \omega^4 \rho^2 \\
b_3 g_\rho^2 \omega^6 \rho^2 &= a_{031} \omega^6 \rho^2 \\
a_1 g_\rho^2 \sigma^1 \rho^2 &= a_{101} \sigma^1 \rho^2 \\
a_2 g_\rho^2 \sigma^2 \rho^2 &= a_{201} \sigma^2 \rho^2 \\
a_3 g_\rho^2 \sigma^3 \rho^2 &= a_{301} \sigma^3 \rho^2 \\
a_4 g_\rho^2 \sigma^4 \rho^2 &= a_{401} \sigma^4 \rho^2 \\
a_5 g_\rho^2 \sigma^5 \rho^2 &= a_{501} \sigma^5 \rho^2 \\
a_6 g_\rho^2 \sigma^6 \rho^2 &= a_{601} \sigma^6 \rho^2
\end{aligned}$$

Note that Muller and Serot use the notation

$$\frac{\bar{\kappa} g_s^3}{2} = \frac{\kappa}{2} = b M g_s^3 \quad \text{and} \quad \frac{\bar{\lambda} g_s^4}{6} = \frac{\lambda}{6} = c g_s^4$$

which differs slightly from the "standard" notation above.

We need to compare the values of

$$\begin{aligned}
&\frac{m_s^2}{2g_s^2 M^2}, \frac{m_\pi^2}{2g_\pi^2 M^2}, \frac{m_\rho^2}{8g_\rho^2 M^2}, \frac{b}{3}, \frac{c}{4} \\
&\frac{\zeta}{24}, \frac{\xi}{384}, \frac{b_1}{4g_\omega^2}, \frac{b_2 M^2}{4g_\omega^4}, \frac{b_3 M^4}{4g_\omega^6}, \frac{a_1}{4g_\sigma M}, \\
&\frac{a_2}{4g_\sigma^2}, \frac{a_3 M}{4g_\sigma^3}, \frac{a_4 M^2}{4g_\sigma^4}, \frac{a_5 M^3}{4g_\sigma^5}, \text{ and } \frac{a_6 M^4}{4g_\sigma^6}.
\end{aligned}$$

These values are stored in the variables cs, cw, cr, b, c, zeta, xi, b1, etc. in the specified [rmf\\_eos](#) object. All of the numbers should be around 0.001 or 0.002.

For the scale  $M$ , [mnuc](#) is used.

**Todo** I may have ignored some signs in the above, which are unimportant for this application, but it would be good to fix them for posterity.

Definition at line 594 of file [rmf\\_eos.h](#).

```
14.34.2.16 int rmf_eos::naturalness_limits( double value, rmf_eos & re ) [inline]
```

The limits for the couplings are function of the nucleon and meson masses, except for the limits on b, c, zeta, and xi which are independent of the masses because of the way that these four couplings are defined.

Definition at line 631 of file [rmf\\_eos.h](#).

```
14.34.2.17 int rmf_eos::calc_cr( double sig, double ome, double nb ) [protected]
```

Used by [fix\\_saturation\(\)](#).

## 14.34.3 Field Documentation

### 14.34.3.1 int rmf\_eos::verbose

This is used by [saturation\(\)](#) to report progress towards computing the properties of nuclear matter near saturation

Definition at line 234 of file [rmf\\_eos.h](#).

### 14.34.3.2 double rmf\_eos::mnuc

This need not be exactly equal to the neutron or proton mass, but provides the scale for the coupling b.

Definition at line 252 of file [rmf\\_eos.h](#).

**14.34.3.3** `gsl_mroot_hybrids<mm_funct<>> rmf_eos::def_sat_mroot`

Used by `fn0()` (which is called by `saturation()`) to solve `saturation_matter_e()` (1 variable).

Definition at line 675 of file `rmf_eos.h`.

**14.34.3.4** `double rmf_eos::n_charge` [protected]

**Todo** Should use `hadronic_eos::proton_frac` instead?

Definition at line 689 of file `rmf_eos.h`.

The documentation for this class was generated from the following file:

- `rmf_eos.h`

**14.35** **rmf\_nucleus Class Reference**

Spherical closed-shell nuclei with a relativistic mean-field model in the Hartree approximation.

```
#include <rmf_nucleus.h>
```

**14.35.1 Detailed Description**

This code is very experimental.

This class is based on a code developed by C.J. Horowitz and B.D. Serot, and used in [Horowitz81](#) which was then adapted by P.J. Ellis and used in [Heide94](#) and [Prakash94](#). Ellis and A.W. Steiner adapted it for the parameterization in in `rmf_eos` for [Steiner05b](#), and then converted to C++ by Steiner afterwards.

The standard usage is something like:

```
rmf_nucleus rn;
o2scl_hdf::rmf_load(rn.rmfm, "NL4");
rn.run_nucleus(82, 208, 0, 0);
cout << rn.rnrp << endl;
```

which computes the structure of  $^{208}\text{Pb}$  and outputs the neutron skin thickness using the model 'NL4'.

Potential exceptions are

- Failed to converge
- Failed to solve meson field equations
- Energy not finite (usually a problem in the equation of state)
- Energy not finite in final calculation
- Function `iterate()` called before `init_run()`
- Not a closed-shell nucleus

The initial level pattern is

```
1 S 1/2
// 2 nucleons
1 P 3/2
1 P 1/2
// 8 nucleus
1 D 5/2
```

```

1 D 3/2
2 S 1/2
// 20 nucleons
1 F 7/2
// 28 nucleons
1 F 5/2
2 P 3/2
2 P 1/2
// 40 nucleons
1 G 9/2
// 50 nucleus
1 G 7/2
2 D 5/2
1 H 11/2
2 D 3/2
3 S 1/2
// 82 nucleons
1 H 9/2
2 F 7/2
1 I 13/2
2 F 5/2
3 P 3/2
3 P 1/2
// 126 nucleons
2 G 9/2
1 I 11/2
1 J 15/2
3 D 5/2
4 S 1/2
2 G 7/2
3 D 3/2
// 184 nucleons

```

Below,  $\alpha$  is a generic index for the isospin, the radial quantum number  $n$  and the angular quantum numbers  $\kappa$  and  $m$ . The meson fields are  $\sigma(r)$ ,  $\omega(r)$  and  $\rho(r)$ . The baryon density is  $n(r)$ , the neutron and proton densities are  $n_n(r)$  and  $n_p(r)$ , and the baryon scalar density is  $n_s(r)$ . The nucleon field equations are

$$\begin{aligned}
F'_\alpha(r) - \frac{\kappa}{r} F_\alpha(r) + \left[ \epsilon_\alpha - g_\omega \omega(r) - t_\alpha g_\rho \rho(r) - (t_\alpha + \frac{1}{2}) e A(r) - M + g_\sigma \sigma(r) \right] G_\alpha(r) &= 0 \\
G'_\alpha(r) + \frac{\kappa}{r} G_\alpha(r) - \left[ \epsilon_\alpha - g_\omega \omega(r) - t_\alpha g_\rho \rho(r) - (t_\alpha + \frac{1}{2}) e A(r) + M - g_\sigma \sigma(r) \right] F_\alpha(r) &= 0
\end{aligned}$$

where  $t_\alpha$  is 1/2 for protons and -1/2 for neutrons. The meson field equations are

$$\begin{aligned}
\sigma''(r) + \frac{2}{r} \sigma'(r) - m_\sigma^2 \sigma &= -g_\sigma n_s(r) + b M g_\sigma^3 \sigma^2 + c g_\sigma^4 \sigma^3 - g_\rho^2 \rho^2 \frac{\partial f}{\partial \sigma} \\
\omega''(r) + \frac{2}{r} \omega'(r) - m_\omega^2 \omega &= -g_\omega n(r) + \frac{\zeta}{6} g_\omega^4 \omega^3 + g_\rho^2 \rho^2 \frac{\partial f}{\partial \omega} \\
\rho''(r) + \frac{2}{r} \rho'(r) - m_\rho^2 \rho &= -\frac{g_\rho}{2} [n_n(r) - n_p(r)] + 2 g_\rho^2 \rho f + \frac{\xi}{6} g_\rho^4 \rho^3
\end{aligned}$$

and the Coulomb field equation is

$$A''(r) + \frac{2}{r} A'(r) = -e n_p(r)$$

The densities are

$$\begin{aligned}
n_s &= \sum_\alpha \left\{ \int d^3 r [g(r)^2 - f(r)^2] \right\} \\
n &= \sum_\alpha \left\{ \int d^3 r [g(r)^2 + f(r)^2] \right\} \\
n_i &= \sum_\alpha \left\{ t_\alpha \int d^3 r [g(r)^2 - f(r)^2] \right\}
\end{aligned}$$

$$n_c = \sum_{\alpha} \left\{ \left[ t_{\alpha} + \frac{1}{2} \right] \int d^3r [g(r)^2 - f(r)^2] \right\}$$

Using the Green function

$$D(r, r', m_i) = \frac{-1}{m_i r r'} \sinh(m_i r_{<}) \exp(-m_i r_{>})$$

one can write the meson field

$$\sigma(r) = \int_0^{\infty} r'^2 dr' [-g_{\sigma} \rho_s(r)] D(r, r', m_{\sigma})$$

When  $r > r'$ , and setting  $x = r' - r$ , the Green function is

$$D = \frac{-1}{m_i r(x+r)} \sinh[m_i(x+r)] \exp(-m_i r)$$

When  $x \gg r$ ,

The total energy is

$$E = \sum_{\alpha} \varepsilon_{\alpha} (2j_{\alpha} + 1) - \frac{1}{2} \int d^3r \left[ -g_{\sigma} \sigma(r) \rho_s(r) + g_{\omega} \omega(r) \rho(r) + \frac{1}{2} g_{\rho} \rho(r) + eA(r) n_p(r) \right]$$

The charge density is the proton density folded with the charge density of the proton, i.e.

$$\rho_{\text{ch}}(r) = \int d^3r' \rho_{\text{prot}}(r - r') \rho_p(r)$$

where the proton charge density is assumed to be of the form

$$\rho_{\text{prot}}(r) = \frac{\mu^3}{8\pi} \exp(-\mu|r|)$$

and the parameter  $\mu = (0.71)^{1/2}$  GeV (see Eq. 20b in [Horowitz81](#)). The default value of [a\\_proton](#) is the value of  $\mu$  converted into fm<sup>-1</sup>.

**Todo** Better documentation

Convert [energies\(\)](#) to use EOS and possibly replace [sigma\\_rhs\(\)](#) and related functions by the associated field equation method of [rmf\\_eos](#).

**Todo** Document `hw=3.923+23.265/cbrt(atot);`

**Idea for Future** Sort energy levels at the end by energy

**Idea for Future** Improve the numerical methods

**Idea for Future** Make the neutron and proton orbitals more configurable

**Idea for Future** Generalize to  $m_n \neq m_p$ .

**Idea for Future** Allow more freedom in the integrations

**Idea for Future** Consider converting everything to inverse fermis.

**Idea for Future** Convert to zero-indexed arrays

**Idea for Future** Warn when the level ordering is wrong, and unoccupied levels are lower energy than occupied levels

Definition at line 237 of file `rmf_nucleus.h`.

**Data Structures**

- struct [initial\\_guess](#)  
*Initial guess structure.*
- struct [odparms](#)  
*A convenient struct for the solution of the Dirac equations.*
- struct [shell](#)  
*A shell of nucleons for [rmf\\_nucleus](#).*

**Public Member Functions****Basic operation**

- void [run\\_nucleus](#) (int nucleus\_Z, int nucleus\_N, int unocc\_Z, int unocc\_N)  
*Computes the structure of a nucleus with the specified number of levels.*
- void [set\\_verbose](#) (int v)  
*Set output level.*

**Lower-level interface**

- void [init\\_run](#) (int nucleus\_Z, int nucleus\_N, int unocc\_Z, int unocc\_N)  
*Initialize a run.*
- void [iterate](#) (int nucleus\_Z, int nucleus\_N, int unocc\_Z, int unocc\_N, int &iconverged)  
*Perform an iteration.*
- int [post\\_converge](#) (int nucleus\_Z, int nucleus\_N, int unocc\_Z, int unocc\_N)  
*After convergence, make CM corrections, etc.*

**Data Fields**

- [initial\\_guess](#) [ig](#)  
*Parameters for initial guess.*
- bool [generic\\_ode](#)  
*If true, use the generic ODE solver instead of...*

**Protected Member Functions**

- int [load\\_nl3](#) ([rmf\\_eos](#) &r)  
*Load the default model NL3 into the given [rmf\\_eos](#) object.*
- void [init\\_meson\\_density](#) ()  
*Initialize the meson fields, the densities, etc.*
- void [energies](#) (double xpro, double xnu, double e)  
*Calculate the energy profile.*
- void [center\\_mass\\_corr](#) (double atot)  
*Compute the center of mass correction.*
- double [gunt](#) (double x, double g1, double f1, double &funt, double eigen, double kappa, [uvector](#) &vvarr)  
*Integrate the Dirac equations using a simple inline 4th order Runge-Kutta.*

**Protected Attributes**

- double [a\\_proton](#)  
*The factor for the charge density of the proton (default 4.2707297)*
- [rmf\\_eos](#) \* [rmf](#)  
*The base EOS.*
- [o2\\_shared\\_ptr](#)< [table\\_units](#) >::type [profiles](#)  
*The radial profiles.*
- [o2\\_shared\\_ptr](#)< [table\\_units](#) >::type [chden\\_table](#)  
*The final charge densities.*
- std::vector< [shell](#) > \* [levp](#)

*A pointer to the current vector of levels (either [levels](#) or [unocc\\_levels](#))*

- int [verbose](#)  
*Control output.*
- [shell\\_neutron\\_shells](#) [[n\\_internal\\_levels](#)]  
*The starting neutron levels.*
- [shell\\_proton\\_shells](#) [[n\\_internal\\_levels](#)]  
*The starting proton levels.*
- double [step\\_size](#)  
*The grid step size.*
- double [mnuc](#)  
*The nucleon mass (automatically set in [init\\_fun\(\)](#))*
- **uvector** [energy](#)  
*Energy profile.*
- bool [init\\_called](#)  
*True if [init\(\)](#) has been called.*
- double [ode\\_y](#) [2]  
*ODE functions.*
- double [ode\\_dydx](#) [2]  
*ODE derivatives.*
- double [ode\\_yerr](#) [2]  
*ODE errors.*

#### Density information (protected)

- **umatrix** [xrho](#)  
*The densities.*
- **uvector** [xrhosp](#)  
*The proton scalar density.*
- **uvector** [xrhos](#)  
*The scalar field RHS.*
- **uvector** [xrhov](#)  
*The vector field RHS.*
- **uvector** [xrhov](#)  
*The iso-vector field RHS.*
- **uvector** [chdenl](#)  
*Charge density.*
- **uvector** [chdenr](#)  
*Charge density.*
- **uvector** [arho](#)  
*Baryon density.*

#### Gauss-Legendre integration points and weights

- double [x12](#) [6]
- double [w12](#) [6]
- double [x100](#) [50]
- double [w100](#) [50]

#### Static Protected Attributes

- static const int [n\\_internal\\_levels](#) = 29  
*The total number of shells stored internally.*
- static const int [grid\\_size](#) = 300  
*The grid size.*

#### Results

- int [nlevels](#)  
*The number of levels.*
- std::vector< [shell](#) > [levels](#)

- *The levels (protons first, then neutrons)*
- int **nuolevels**  
*The number of unoccupied levels (equal to `unocc_Z` + `unocc_N`)*
- int **last\_conv**  
*Information on the last convergence error.*
- std::vector< **shell** > **unocc\_levels**  
*The unoccupied levels (protons first, then neutrons)*
- double **stens**  
*Surface tension (in fm<sup>-3</sup>)*
- double **mrp**  
*Skin thickness (in fm)*
- double **nrms**  
*Neutron RMS radius (in fm)*
- double **rprms**  
*Proton RMS radius (in fm)*
- double **etot**  
*Total energy (in MeV)*
- double **r\_charge**  
*Charge radius (in fm)*
- double **r\_charge\_cm**  
*Charge radius corrected by the center of mass (in fm)*
- o2\_shared\_ptr< **table\_units** >::type **get\_profiles** ()  
*Get the radial profiles.*
- o2\_shared\_ptr< **table\_units** >::type **get\_chden** ()  
*The final charge densities.*

#### Equation of state

- **rmf\_eos** **def\_rmf**  
*The default equation of state (default NL3)*
- **thermo** **hb**  
*thermo object for the EOS*
- **fermion** **n**  
*The neutron.*
- **fermion** **p**  
*The proton.*
- int **set\_eos** (**rmf\_eos** &r)  
*Set the base EOS to be used.*

#### Numeric configuration

- typedef double **arr\_t** [2]  
*The array type for the ODE solver.*
- bool **err\_nonconv**  
*If true, call the error handler if the routine does not converge or reach the desired tolerance (default true)*
- int **itmax**  
*Maximum number of total iterations (default 70)*
- int **meson\_itmax**  
*Maximum number of iterations for solving the meson field equations (default 10000)*
- int **dirac\_itmax**  
*Maximum number of iterations for solving the Dirac equations (default 100)*
- double **dirac\_tol**  
*Tolerance for Dirac equations (default  $5 \times 10^{-3}$ ).*
- double **dirac\_tol2**  
*Second tolerance for Dirac equations (default  $5 \times 10^{-4}$ ).*
- double **meson\_tol**  
*Tolerance for meson field equations (default  $10^{-6}$ ).*
- void **set\_step** (**ode\_step**< **ode\_funct**< **arr\_t** >, **arr\_t** > &step)  
*Set the stepper for the Dirac differential equation.*



## The meson fields and field equations (protected)

- **umatrix** [field0](#)  
*Values of the fields from the last iteration.*
- **umatrix** [fields](#)  
*The values of the fields.*
- **umatrix** [gin](#)  
*The Green's functions inside.*
- **umatrix** [gout](#)  
*The Green's functions outside.*
- **int** [surf\\_index](#)  
*The grid index corresponding to the nuclear surface (computed by [init\\_run\(\)](#))*
- **double** [sigma\\_rhs](#) (**double** sig, **double** ome, **double** rho)  
*Scalar density RHS.*
- **double** [omega\\_rhs](#) (**double** sig, **double** ome, **double** rho)  
*Vector density RHS.*
- **double** [rho\\_rhs](#) (**double** sig, **double** ome, **double** rho)  
*Iso-vector density RHS.*
- **void** [mesint](#) ()  
*Calculate meson Green's functions.*
- **void** [meson](#) (**double** ic)  
*Calculate meson fields.*
- **void** [meson\\_solve](#) ()  
*Solve for the meson profiles.*

## Calculating the form factor, etc. (protected)

- **smart\_interp\_vec**< **uvector\_base**, **uvector\_const\_subvector**, **uvector**, **uvector\_alloc** > \* [gi](#)  
*Interpolation object.*
- **void** [pfold](#) (**double** x, **double** &xrhof)  
*Fold in proton form factor.*
- **double** [xpform](#) (**double** x, **double** xp, **double** a)  
*Function representing proton form factor.*
- **void** [gauss](#) (**double** xmin, **double** xmax, **double** x, **double** &xi)  
*Perform integrations for form factor.*
- **double** [xrhof](#) (**double** x1)  
*Desc.*

## Solving the Dirac equations (protected)

- **gsl\_rkck**< **ode\_funct**< [arr\\_t](#) > , [arr\\_t](#), [arr\\_t](#), **array\_alloc** < [arr\\_t](#) > > [def\\_step](#)  
*The default stepper.*
- **ode\_step**< **ode\_funct**< [arr\\_t](#) > , [arr\\_t](#) > \* [ostep](#)  
*The ODE stepper.*
- **void** [dirac](#) (**int** ilevel)  
*Solve the Dirac equations.*
- **void** [dirac\\_step](#) (**double** &x, **double** h, **double** eigen, **double** kappa, **uvector** &varr)  
*Take a step in the Dirac equations.*
- **int** [odefun](#) (**double** x, **size\_t** nv, **const** [arr\\_t](#) &y, [arr\\_t](#) &dydx, **odparms** &op)  
*The form of the Dirac equations for the ODE stepper.*
- **void** [field](#) (**double** x, **double** &s, **double** &v, **uvector\_base** &varr)  
*Compute the fields for the Dirac equations.*

## 14.35.2 Member Function Documentation

14.35.2.1 **void** [rmf\\_nucleus::run\\_nucleus](#) ( **int** *nucleus\_Z*, **int** *nucleus\_N*, **int** *unocc\_Z*, **int** *unocc\_N* )

Note that [rmf](#) must be set before [run\\_nucleus\(\)](#) is called.

This calls [init\\_run\(\)](#), and then [iterate\(\)](#) until [iconverged](#) is 1, and then [post\\_converge\(\)](#).

14.35.2.2 void rmf\_nucleus::init\_run ( int *nucleus\_Z*, int *nucleus\_N*, int *unocc\_Z*, int *unocc\_N* )

Note that `rmf` must be set before `run_nucleus()` is called.

14.35.2.3 o2\_shared\_ptr<table\_units>::type rmf\_nucleus::get\_profiles ( ) [inline]

The profiles are calculated each iteration by `iterate()`.

Definition at line 329 of file `rmf_nucleus.h`.

14.35.2.4 int rmf\_nucleus::set\_eos ( rmf\_eos & *r* ) [inline]

The equation of state must be set before `run_nucleus()` or `init_fun()` are called, including the value of `rmf_eos::mnuc`.

Definition at line 419 of file `rmf_nucleus.h`.

14.35.2.5 void rmf\_nucleus::dirac ( int *ilevel* ) [protected]

Solves the Dirac equation in from 12 fm to the match point and then out from .04 fm and adjusts eigenvalue with

$$\Delta\mathcal{E} = -g(r = \text{match\_point}) \times (f^+ - f^-)$$

### 14.35.3 Field Documentation

14.35.3.1 std::vector<shell> rmf\_nucleus::levels

An array of size `nlevels`

Definition at line 342 of file `rmf_nucleus.h`.

14.35.3.2 std::vector<shell> rmf\_nucleus::unocc\_levels

An array of size `nuolevels`

Definition at line 355 of file `rmf_nucleus.h`.

14.35.3.3 double rmf\_nucleus::stens

Computed in `post_converge()` or automatically in `run_nucleus()`

Definition at line 361 of file `rmf_nucleus.h`.

14.35.3.4 double rmf\_nucleus::rnrp

Computed every iteration in `iterate()` or automatically in `run_nucleus()`

Definition at line 368 of file `rmf_nucleus.h`.

14.35.3.5 double rmf\_nucleus::rnrms

Computed every iteration in `iterate()` or automatically in `run_nucleus()`

Definition at line 375 of file `rmf_nucleus.h`.

14.35.3.6 double rmf\_nucleus::rprms

Computed every iteration in `iterate()` or automatically in `run_nucleus()`

Definition at line 382 of file `rmf_nucleus.h`.

14.35.3.7 double rmf\_nucleus::etot

Computed every iteration in `iterate()` or automatically in `run_nucleus()`

Definition at line 389 of file `rmf_nucleus.h`.

#### 14.35.3.8 `double rmf_nucleus::r_charge`

Computed in `post_converge()` or automatically in `run_nucleus()`

Definition at line 395 of file `rmf_nucleus.h`.

#### 14.35.3.9 `double rmf_nucleus::r_charge_cm`

Computed in `post_converge()` or automatically in `run_nucleus()`

Definition at line 401 of file `rmf_nucleus.h`.

#### 14.35.3.10 `rmf_eos rmf_nucleus::def_rmf`

This is set in the constructor to be the default model, NL3, using the function `load_nl3()`.

Definition at line 412 of file `rmf_nucleus.h`.

#### 14.35.3.11 `thermo rmf_nucleus::hb`

This is just used as temporary storage.

Definition at line 428 of file `rmf_nucleus.h`.

#### 14.35.3.12 `fermion rmf_nucleus::n`

The mass of the neutron is ignored and set by `init_run()` to be `rmf_eos::mnuc` from `rmf`.

Definition at line 435 of file `rmf_nucleus.h`.

#### 14.35.3.13 `fermion rmf_nucleus::p`

The mass of the proton is ignored and set by `init_run()` to be `rmf_eos::mnuc` from `rmf`.

Definition at line 442 of file `rmf_nucleus.h`.

#### 14.35.3.14 `bool rmf_nucleus::err_nonconv`

If this is false, the function proceeds normally and may provide convergence information in `last_conv`.

Definition at line 454 of file `rmf_nucleus.h`.

#### 14.35.3.15 `initial_guess rmf_nucleus::ig`

Default is {310,240,-6,25.9,6.85,0.6}

Definition at line 518 of file `rmf_nucleus.h`.

The documentation for this class was generated from the following file:

- `rmf_nucleus.h`

## 14.36 rms\_radius Class Reference

Compute the RMS radius of a Fermi-Dirac density distribution with fixed diffusiveness.

```
#include <ldrop_mass.h>
```

## 14.36.1 Detailed Description

This class computes the RMS radius given either the central density or the radius specified in the Fermi function. This class assumes the density distribution function is of the form

$$N = 4\pi\rho_0 \int r^2 dr \{1 + \exp[(r - R_{\text{fermi}})]\}^{-1}$$

where  $N$  is the total number of particles, and  $\rho_0$  is the central density.

The radius assuming constant density,

$$R_{\text{cd}} = \left( \frac{3N}{4\pi\rho_0} \right)^{1/3},$$

is also given.

Definition at line 60 of file ldrops.h.

## Public Member Functions

- int [eval\\_rms\\_rho](#) (double rho0, double N, double d, double &Rcd, double &Rfermi, double &Rrms)  
*Compute the RMS radius from the central density.*
- int [eval\\_rms\\_rs](#) (double Rfermi, double N, double d, double &rho0, double &Rcd, double &Rrms)  
*Compute the RMS radius from the Fermi distribution radius.*

## Protected Member Functions

- double [iand](#) (double r)  
*The function  $4\pi r^4 \rho(r)$ .*
- double [iand2](#) (double r)  
*The function  $4\pi r^2 \rho(r)$ .*
- double [solve](#) (double x)  
*The function to fix the total number of particles.*

## Protected Attributes

- double [urho0](#)  
*The central density.*
- double [ud](#)  
*The diffusiveness.*
- double [uRfermi](#)  
*Store the user-specified value of the radius in the Fermi distribution.*
- double [uN](#)  
*The total number of particles.*
- int [pa](#)  
*Blank parameter.*
- [gsl\\_integrator](#) < [funct](#) > [it](#)  
*The integrator.*
- [cern\\_minimizer](#) < [funct](#) > [cr](#)  
*The solver.*

## 14.36.2 Member Function Documentation

## 14.36.2.1 int rms\_radius::eval\_rms\_rho ( double rho0, double N, double d, double &amp; Rcd, double &amp; Rfermi, double &amp; Rrms )

Computes the RMS radius  $R_{\text{rms}}$  from the central density  $\rho_0$ , the number of particles  $N$ , and the diffusiveness  $d$ . This function also computes the radius in the Fermi distribution function,  $R_{\text{fermi}}$  and the radius assuming constant density,  $R_{\text{cd}}$ .

14.36.2.2 `int rms_radius::eval_rms_rsqr( double Rfermi, double N, double d, double & rho0, double & Rcd, double & Rrms )`

Computes the RMS radius `Rrms` from the radius `Rfermi` in the Fermi distribution assuming a total number of particles `N`, a diffusiveness paramter `d`. This function also produces the central density `rho0`, and the radius assuming constant density, `Rcd`.

### 14.36.3 Field Documentation

14.36.3.1 `double rms_radius::uRfermi` [protected]

This is used in the integrands `iand()` and `iand2()`.

Definition at line 73 of file `ldrop_mass.h`.

The documentation for this class was generated from the following file:

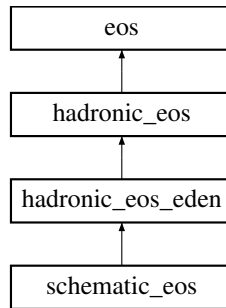
- `ldrop_mass.h`

## 14.37 schematic\_eos Class Reference

Schematic hadronic equation of state.

`#include <schematic_eos.h>`

Inheritance diagram for `schematic_eos`:



### 14.37.1 Detailed Description

A schematic equation of state defined by the energy density:

$$\varepsilon = n_n m_n + n_p m_p + n \left\{ eoa + \frac{comp}{18} (n/n_0 - 1)^2 + \frac{kprime}{162} (n/n_0 - 1)^3 + \frac{kpp}{1944} (n/n_0 - 1)^4 + (1 - 2x)^2 \left[ a \left( \frac{n}{n_0} \right)^{2/3} + b \left( \frac{n}{n_0} \right)^\gamma \right] \right\}$$

Symmetry energy at nuclear matter density is  $a+b$ .

Note that it doesn't really matter what kind of particle object is used, since the `calc_e()` function doesn't use any of the particle thermodynamics functions.

Definition at line 52 of file `schematic_eos.h`.

### Public Member Functions

- virtual int `calc_e(fermion &ln, fermion &lp, thermo &lth)`  
Equation of state as a function of density.
- virtual int `set_kprime_zeroden()`  
Set `kprime` so that the energy per baryon of zero-density matter is zero.

- virtual int `set_kpp_zeroden` ()  
*Set kpp so that the energy per baryon of zero-density matter is zero.*
- virtual int `set_a_from_mstar` (double `u_msom`, double `mnuc`)  
*Fix the kinetic energy symmetry coefficient from the nucleon effective mass and the saturation density.*
- virtual double `eo_a_zeroden` ()  
*Return the energy per baryon of matter at zero density.*
- virtual const char \* `type` ()  
*Return string denoting type ("schematic\_eos")*

#### Data Fields

- double `a`  
*The kinetic energy symmetry coefficient in inverse fm (default 17/hc)*
- double `b`  
*The potential energy symmetry coefficient in inverse fm (default 13/hc)*
- double `kpp`  
*The coefficient of a density to the fourth term (default 0)*
- double `gamma`  
*The exponent of the high-density symmetry energy (default 1.0)*

#### 14.37.2 Member Function Documentation

14.37.2.1 virtual int `schematic_eos::set_a_from_mstar` ( double `u_msom`, double `mnuc` ) [inline, virtual]

This assumes the nucleons are non-relativistic and that the neutrons and protons have equal mass. The relativistic corrections are around 1 part in  $10^6$ .

**Todo** This was computed in `schematic_sym.nb`, which might be added to the documentation?

Definition at line 107 of file `schematic_eos.h`.

14.37.2.2 virtual double `schematic_eos::eo_a_zeroden` ( ) [inline, virtual]

This is inaccessible from `calc_e()` so is available separately here. Using `set_kprime_zeroden()` or `set_kpp_zeroden()` will fix `kprime` or `kpp` (respectively) to ensure that this is zero.

The result provided here does not include the nucleon mass and is given in  $\text{fm}^{-1}$ .

Definition at line 122 of file `schematic_eos.h`.

#### 14.37.3 Field Documentation

14.37.3.1 double `schematic_eos::a`

The default value corresponds to an effective mass of about 0.7.

Definition at line 62 of file `schematic_eos.h`.

The documentation for this class was generated from the following file:

- `schematic_eos.h`

#### 14.38 `rmf_nucleus::shell` Struct Reference

A shell of nucleons for `rmf_nucleus`.

```
#include <rmf_nucleus.h>
```

## 14.38.1 Detailed Description

Definition at line 270 of file `rmf_nucleus.h`.

## Data Fields

- int `twojpl`  
*Degeneracy  $2j+1$ .*
- int `kappa`  
 $\kappa$
- double `energy`  
*Energy eigenvalue.*
- double `isospin`  
*Isospin ( $+1/2$  or  $-1/2$ ).*
- std::string `state`  
*Angular momentum-spin state  $2s+1\ell_j$ .*
- double `match_point`  
*Matching radius (in fm)*
- double `eigen`  
*Desc.*
- double `eigenc`  
*Desc.*
- int `nodes`  
*Number of nodes in the wave function.*

The documentation for this struct was generated from the following file:

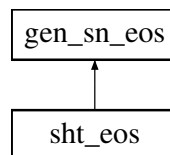
- `rmf_nucleus.h`

## 14.39 sht\_eos Class Reference

A class to manipulate the G. Shen et al. EOS.

```
#include <gen_sn_eos.h>
```

Inheritance diagram for `sht_eos`:



## 14.39.1 Detailed Description

This class is experimental.

## Note

O<sub>2</sub>scl Does not contain the EOS, only provides some code to manipulate it. This class was designed to be used with the FS-U models given at [http://cecelia.physics.indiana.edu/gang\\_shen\\_eos/FSU/fsu.html](http://cecelia.physics.indiana.edu/gang_shen_eos/FSU/fsu.html) and stored in the files named `FSU1.7eos1.01.dat`, `FSU2.1eos1.01.dat`, `FSU1.7eosb1.01.dat`, and `FSU2.1eosb1.01.dat` corresponding to [mode\\_17](#), [mode\\_21](#), [mode\\_17b](#) and [mode\\_21b](#) respectively.

See also the documentation at [gen\\_sn\\_eos](#).

**Todo** More documentation

Definition at line 756 of file gen\_sn\_eos.h.

### Public Member Functions

- virtual void [load](#) (std::string fname, size\_t mode)  
*Load table from filename `fname`.*
- virtual void [load](#) (std::string fname)  
*Load table from filename `fname`.*
- virtual void [beta\\_eq\\_T0](#) (size\_t i, double &nb, double &E\_beta, double &P\_beta, double &Ye\_beta, double &Z\_beta, double &A\_beta)  
*Compute properties of matter in beta equilibrium at zero temperature at a baryon density grid point.*

### Data Fields

- [tensor\\_grid3](#) & [T](#)  
*Temperature in MeV.*
- [tensor\\_grid3](#) & [Yp](#)  
*Proton fraction.*
- [tensor\\_grid3](#) & [nB](#)  
*Baryon number density in  $1/\text{fm}^3$ .*
- [tensor\\_grid3](#) & [mue](#)  
*Electron chemical potential in MeV.*
- [tensor\\_grid3](#) & [M\\_star](#)  
*Nucleon effective mass in MeV.*

### Static Public Attributes

- static const size\_t [mode\\_17](#) = 0  
*1.7 solar masses with leptons and photons*
- static const size\_t [mode\\_21](#) = 1  
*2.1 solar masses with leptons and photons*
- static const size\_t [mode\\_17b](#) = 2  
*1.7 solar masses without leptons and photons*
- static const size\_t [mode\\_21b](#) = 3  
*2.1 solar masses without leptons and photons*
- static const size\_t [mode\\_NL3](#) = 4  
*NL3 model with leptons and photons.*
- static const size\_t [mode\\_NL3b](#) = 5  
*NL3 model with leptons and photons.*

The documentation for this class was generated from the following file:

- gen\_sn\_eos.h

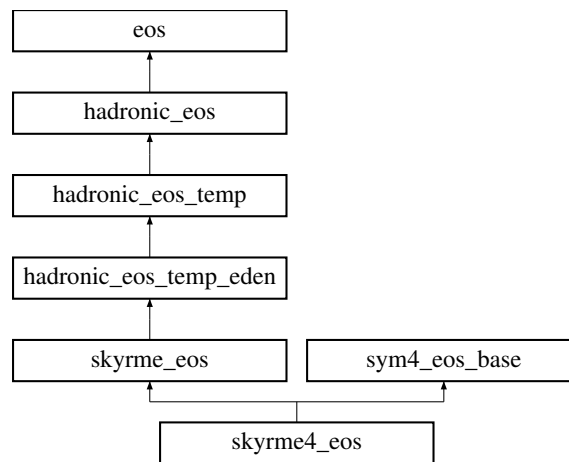
## 14.40 skyrme4\_eos Class Reference

A version of [skyrm\\_eos](#) to separate potential and kinetic contributions.

```
#include <sym4_eos.h>
```

Inheritance diagram for skyrme4\_eos:





#### 14.40.1 Detailed Description

#### References:

Created for [Steiner06](#).

Definition at line 144 of file `sym4_eos.h`.

#### Public Member Functions

- virtual int [calc\\_e\\_sep](#) (**fermion** &ne, **fermion** &pr, double &ed\_kin, double &ed\_pot, double &mu\_n\_kin, double &mu\_p\_kin, double &mu\_n\_pot, double &mu\_p\_pot)  
*Compute the potential and kinetic parts separately.*

The documentation for this class was generated from the following file:

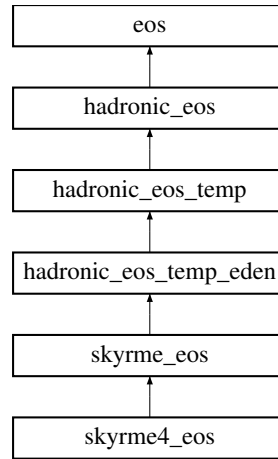
- `sym4_eos.h`

## 14.41 skyrme\_eos Class Reference

Skyrme hadronic equation of state.

```
#include <skyrme_eos.h>
```

Inheritance diagram for `skyrme_eos`:



#### 14.41.1 Detailed Description

Equation of state of nucleonic matter based on the Skyrme interaction from [Skyrme59](#) .

#### Background:

The Hamiltonian is defined (using the notation of [Steiner05b](#) )

$$\mathcal{H} = \mathcal{H}_{k1} + \mathcal{H}_{k2} + \mathcal{H}_{k3} + \mathcal{H}_{p1} + \mathcal{H}_{p2} + \mathcal{H}_{p3} + \mathcal{H}_{g1} + \mathcal{H}_{g2}$$

The kinetic terms are:

$$\begin{aligned}\mathcal{H}_{k1} &= \frac{\tau_n}{2m_n} + \frac{\tau_p}{2m_p} \\ \mathcal{H}_{k2} &= n(\tau_n + \tau_p) \left[ \frac{t_1}{4} \left( 1 + \frac{x_1}{2} \right) + \frac{t_2}{4} \left( 1 + \frac{x_2}{2} \right) \right] \\ \mathcal{H}_{k3} &= (\tau_n n_n + \tau_p n_p) \left[ \frac{t_2}{4} \left( \frac{1}{2} + x_2 \right) - \frac{t_1}{4} \left( \frac{1}{2} + x_1 \right) \right]\end{aligned}$$

where  $\tau_i$  are defined

The potential terms are:

$$\begin{aligned}\mathcal{H}_{p1} &= \frac{t_0}{2} \left[ \left( 1 + \frac{x_0}{2} \right) n^2 - \left( \frac{1}{2} + x_0 \right) (n_n^2 + n_p^2) \right] \\ \mathcal{H}_{p2} &= \frac{at_3}{6} \left[ \left( 1 + \frac{x_3}{2} \right) n^\alpha n_n n_p + 2^{\alpha-2} (1 - x_3) (n_n^{\alpha+2} + n_p^{\alpha+2}) \right] \\ \mathcal{H}_{p3} &= \frac{bt_3}{12} \left[ \left( 1 + \frac{x_3}{2} \right) n^{\alpha+2} - \left( \frac{1}{2} + x_3 \right) n^\alpha (n_n^2 + n_p^2) \right]\end{aligned}$$

The gradient terms are displayed here for completeness even though they are not computed in the code:

$$\begin{aligned}\mathcal{H}_{g1} &= \frac{3}{32} [t_1 (1 - x_1) - t_2 (1 + x_2)] [(\nabla n_n)^2 + (\nabla n_p)^2] \\ \mathcal{H}_{g2} &= \frac{1}{8} \left[ 3t_1 \left( 1 + \frac{x_1}{2} \right) - t_2 \left( 1 + \frac{x_2}{2} \right) \right] \nabla n_n \nabla n_p\end{aligned}$$

The values  $a = 0, b = 1$  give the standard definition of the Skyrme Hamiltonian [Skyrme59](#), while  $a = 1, b = 0$  contains the modifications suggested by [Onsi94](#).

Also, couple useful definitions

$$t'_3 = (a + b)t_3,$$

$$C = \frac{3}{10m} \left( \frac{3\pi^2}{2} \right)^{2/3},$$

and

$$\beta = \frac{M}{2} \left[ \frac{1}{4} (3t_1 + 5t_2) + t_2 x_2 \right]$$

---

### Units:

Quantities which have units containing powers of energy are divided by  $\hbar c$  to ensure all quantities are in units of  $fm$ . The  $x_i$  and  $\alpha$  are unitless, while the original units of the  $t_i$  are:

- $t_0$  - MeV fm<sup>3</sup>
- $t_1$  - MeV fm<sup>5</sup>
- $t_2$  - MeV fm<sup>5</sup>
- $t_3$  - MeV fm<sup>3(1+ $\alpha$ )</sup>

These are stored internally with units of:

- $t_0$  - fm<sup>2</sup>
- $t_1$  - fm<sup>4</sup>
- $t_2$  - fm<sup>4</sup>
- $t_3$  - fm<sup>2+3 $\alpha$</sup>

The class `skyrme_eos_io` uses `o2scl_const::hc_mev_fm` for I/O so that all files contain the parameters in the original units.

---

### Misc:

The functions for the usual saturation properties are based partly on [Brack85](#).

Models are taken from the references: [Bartel79](#), [Beiner75](#), [Chabanat95](#), [Chabanat97](#), [Danielewicz08](#), [Dobaczewski94](#), [Dutta86](#), [Friedrich86](#), [Onsi94](#), [Reinhard95](#), and [Tondeur84](#), and [VanGiai81](#).

See Mathematica notebook at

```
doc/o2scl/extras/skyrme_eos.nb
doc/o2scl/extras/skyrme_eos.ps
```

The variables  $v_n$  and  $v_p$  contain the expressions  $(-\mu_n + V_n)/T$  and  $(-\mu_p + V_p)/T$  respectively, where  $V$  is the potential part of the single particle energy for particle  $i$  (i.e. the derivative of the Hamiltonian w.r.t. density while energy density held constant). Equivalently,  $v_n$  is just  $-k_{F_n}^2/2m^*$ .

### Note

The finite temperature code does not include attempt to include antiparticles and uses `part::calc_density()`.

Since this EOS uses the effective masses and chemical potentials in the fermion class, the values of `part::non_interacting` for neutrons and protons are set to false in many of the functions.

---

**Todo** • Make sure that this class properly handles particles for which `inc_rest_mass` is true/false

---

- What about the spin-orbit units?
- Need to write a function that calculates saturation density?
- Remove use of mnuc in calparfun()?
- The compressibility could probably use some simplification
- Make sure the finite-temperature part is properly tested
- The testing code doesn't work if err\_mode is 2, probably because of problems in load().
- Document load() file format.
- Update reference list.

**Idea for Future** • There is some code duplication between `calc_e()` and `calc_temp_e()` which could be simplified.

---

Definition at line 218 of file skyrme\_eos.h.

#### Public Member Functions

- int `calpar` (double gt0=-10.0, double gt3=70.0, double galpha=0.2, double gt1=2.0, double gt2=-1.0)  
*Calculate  $t_0, t_1, t_2, t_3$  and  $\alpha$  from the saturation properties.*
- int `check_landau` (double nb, double m)  
*Check the Landau parameters for instabilities.*
- int `landau_nuclear` (double n0, double m, double &f0, double &g0, double &f0p, double &g0p, double &f1, double &g1, double &f1p, double &g1p)  
*Calculate the Landau parameters for nuclear matter.*
- int `landau_neutron` (double n0, double m, double &f0, double &g0, double &f1, double &g1)  
*Calculate the Landau parameters for neutron matter.*
- virtual const char \* `type` ()  
*Return string denoting type ("skyrme\_eos")*

#### Basic usage

- `skyrme_eos` ()  
*Create a blank Skyrme EOS.*
- virtual `~skyrme_eos` ()  
*Destructor.*
- virtual int `calc_temp_e` (**fermion** &ne, **fermion** &pr, double temper, **thermo** &th)  
*Equation of state as a function of densities.*
- virtual int `calc_e` (**fermion** &ne, **fermion** &pr, **thermo** &lt)  
*Equation of state as a function of density.*

#### Saturation properties

*These calculate the various saturation properties exactly from the parameters at any density. These routines often assume that the neutron and proton masses are equal.*

- virtual double `feoa` (double nb)  
*Calculate binding energy.*
  - virtual double `fmsom` (double nb)  
*Calculate effective mass.*
  - virtual double `fcomp` (double nb)  
*Calculate compressibility.*
  - virtual double `fesym` (double nb, double alpha=0.0)  
*Calculate symmetry energy.*
  - virtual double `fkprime` (double nb)  
*skewness*
-

## Data Fields

- double **W0**  
*Spin-orbit splitting.*
- bool **parent\_method**  
*Use [hadronic\\_eos](#) methods for saturation properties.*
- bool **mu\_at\_zero\_density**  
*Desc.*

## Basic Skyrme model parameters

- double **t0**
- double **t1**
- double **t2**
- double **t3**
- double **x0**
- double **x1**
- double **x2**
- double **x3**
- double **alpha**
- double **a**
- double **b**

## 14.41.2 Member Function Documentation

14.41.2.1 `virtual int skyrme_eos::calc_temp_e( fermion & ne, fermion & pr, double temper, thermo & th )` [virtual]

## Note

Runs the zero temperature code if `temper` is less than or equal to zero.

Implements [hadronic\\_eos\\_temp\\_edn](#).

14.41.2.2 `virtual double skyrme_eos::feoa( double nb )` [virtual]

$$\frac{E}{A} = Cn_B^{2/3} (1 + \beta n_B) + \frac{3t_0}{8} n_B + \frac{t'_3}{16} n_B^{\alpha+1}$$

14.41.2.3 `virtual double skyrme_eos::fmsom( double nb )` [virtual]

$$M^*/M = (1 + \beta n_B)^{-1}$$

14.41.2.4 `virtual double skyrme_eos::fcomp( double nb )` [virtual]

$$K = 10Cn_B^{2/3} + \frac{27}{4}t_0n_B + 40C\beta n_B^{5/3} + \frac{9t'_3}{16}\alpha(\alpha+1)n_B^{1+\alpha} + \frac{9t'_3}{8}(\alpha+1)n_B^{1+\alpha}$$

14.41.2.5 `virtual double skyrme_eos::fesym( double nb, double alpha=0.0 )` [virtual]

If `pf=0.5`, then the exact expression below is used. Otherwise, the method from class [hadronic\\_eos](#) is used.

$$E_{sym} = \frac{5}{9}Cn^{2/3} + \frac{10Cm}{3} \left[ \frac{t_2}{6} \left( 1 + \frac{5}{4}x_2 \right) - \frac{1}{8}t_1x_1 \right] n^{5/3} - \frac{t'_3}{24} \left( \frac{1}{2} + x_3 \right) n^{1+\alpha} - \frac{t_0}{4} \left( \frac{1}{2} + x_0 \right) n$$

14.41.2.6 virtual double skyrme\_eos::fkprime ( double *nb* ) [virtual]

$$2Cn_B^{2/3} (9 - 5/M^*/M) + \frac{27t'_3}{16} n^{1+\alpha} \alpha (\alpha^2 - 1)$$

14.41.2.7 int skyrme\_eos::calpar ( double *gt0* = -10.0, double *gt3* = 70.0, double *galpha* = 0.2, double *gt1* = 2.0, double *gt2* = -1.0 )

In nuclear matter:

$$E_b = E_b(n_0, M^*, t_0, t_3, \alpha)$$

$$P = P(n_0, M^*, t_0, t_3, \alpha)$$

$$K = K(n_0, M^*, t_3, \alpha) \text{ (the } t_0 \text{ dependence vanishes)}$$

$$M^* = M^*(n_0, t_1, t_2, x_2) \text{ (the } x_1 \text{ dependence cancels),}$$

$$E_{sym} = E_{sym}(x_0, x_1, x_2, x_3, t_0, t_1, t_2, t_3, \alpha)$$

To fix the couplings from the saturation properties, we take  $n_0, M^*, E_b, K$  as inputs, and we can fix  $t_0, t_3, \alpha$  from the first three relations, then use  $M^*, E_b$  to fix  $t_2$  and  $t_1$ . The separation into two solution steps should make for better convergence. All of the  $x$ 's are free parameters and should be set before the function call.

The arguments *gt0*, *gt3*, *galpha*, *gt1*, and *gt2* are used as initial guesses for `skyrme_eos::t0`, `skyrme_eos::t3`, `skyrme_eos::alpha`, `skyrme_eos::t1`, and `skyrme_eos::t2` respectively.

**Todo** Does this work for both 'a' and 'b' non-zero?

**Todo** Compare to similar formulae from [Margueron02](#)

14.41.2.8 int skyrme\_eos::check\_landau ( double *nb*, double *m* )

This returns zero if there are no instabilities.

14.41.2.9 int skyrme\_eos::landau\_nuclear ( double *n0*, double *m*, double & *f0*, double & *g0*, double & *f0p*, double & *g0p*, double & *f1*, double & *g1*, double & *f1p*, double & *g1p* )

Given *n0* and *m*, this calculates the Landau parameters in nuclear matter as given in [Margueron02](#)

**Todo** This needs to be checked.

(Checked once on 11/05/03)

14.41.2.10 int skyrme\_eos::landau\_neutron ( double *n0*, double *m*, double & *f0*, double & *g0*, double & *f1*, double & *g1* )

Given 'n0' and 'm', this calculates the Landau parameters in neutron matter as given in [Margueron02](#)

**Todo** This needs to be checked

(Checked once on 11/05/03)

### 14.41.3 Field Documentation

#### 14.41.3.1 double skyrme\_eos::W0

This is unused, but included for possible future use and present in the internally stored models.

Definition at line 231 of file `skyrme_eos.h`.

## 14.41.3.2 bool skyrme\_eos::parent\_method

This can be set to true to check the difference between the exact expressions and the numerical values from class [hadronic\\_eos](#).

Definition at line 362 of file skyrme\_eos.h.

The documentation for this class was generated from the following file:

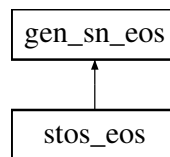
- skyrme\_eos.h

## 14.42 stos\_eos Class Reference

The Shen et al. supernova EOS.

```
#include <gen_sn_eos.h>
```

Inheritance diagram for stos\_eos:



## 14.42.1 Detailed Description

This class is experimental.

**Note**

O<sub>2</sub>scl Does not contain the EOS, only provides some code to manipulate it. This class is designed to be used with the file which was originally called `eos.tab` and now referred to as `eos1.tab` and stored e.g. at <http://user.numazu-ct.ac.jp/~sumi/eos/>.

In order to force the EOS to a uniform grid, linear interpolation is used to recast the variation in baryon density, choosing the grid in baryon density to be the same as the section in the table with  $T=0.1$  MeV and  $Y_p = 0.1$  for all temperature and proton fraction points.

Also, the original EOS is tabulated for constant proton fraction, and this O<sub>2</sub>scl interface assumes that the electron fraction is equal to the proton fraction. Currently, this is a problem only at higher densities where muons might appear.

The data for [gen\\_sn\\_eos::E](#), [gen\\_sn\\_eos::F](#), [gen\\_sn\\_eos::S](#), and [gen\\_sn\\_eos::P](#) is not stored in the table but can be computed with [gen\\_sn\\_eos::compute\\_eg\(\)](#).

See also the documentation at [gen\\_sn\\_eos](#).

See [Shen98](#) and [Shen98b](#).

**Note**

Thanks to Matthias Hempel for providing the correct temperature grid.

**Idea for Future** Add the  $T=0$  and  $Y_e=0$  data to this class

Definition at line 642 of file gen\_sn\_eos.h.

**Public Member Functions**

- virtual void [load](#) (std::string fname)

*Load table from filename `fname`.*

- virtual void [load](#) (std::string `fname`, size\_t `mode`)

*Load table from filename `fname`.*

- virtual void [beta\\_eq\\_T0](#) (size\_t `i`, double &`nb`, double &`E_beta`, double &`P_beta`, double &`Ye_beta`, double &`Z_beta`, double &`A_beta`)

*Compute properties of matter in beta equilibrium at zero temperature at a baryon density grid point.*

#### Data Fields

- [tensor\\_grid3](#) & [log\\_rho](#)

*Logarithm of baryon number density in g/cm<sup>3</sup>.*

- [tensor\\_grid3](#) & [nB](#)

*Baryon number density in fm<sup>-3</sup>.*

- [tensor\\_grid3](#) & [log\\_Y](#)

*Logarithm of proton fraction.*

- [tensor\\_grid3](#) & [Yp](#)

*Proton fraction.*

- [tensor\\_grid3](#) & [M\\_star](#)

*Nucleon effective mass in MeV.*

- [tensor\\_grid3](#) & [quark\\_frac](#)

*Fraction of quark matter.*

#### Static Public Attributes

- static const size\_t [orig\\_mode](#) = 0
- static const size\_t [quark\\_mode](#) = 1

#### 14.42.2 Member Function Documentation

14.42.2.1 virtual void [stos\\_eos::beta\\_eq\\_T0](#) ( size\_t `i`, double & `nb`, double & `E_beta`, double & `P_beta`, double & `Ye_beta`, double & `Z_beta`, double & `A_beta` ) [inline, virtual]

This EOS table doesn't have T=0 results, so we extrapolate from the two low-temperature grid points.

Implements [gen\\_sn\\_eos](#).

Definition at line 687 of file [gen\\_sn\\_eos.h](#).

The documentation for this class was generated from the following file:

- [gen\\_sn\\_eos.h](#)

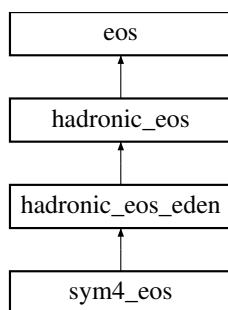
#### 14.43 sym4\_eos Class Reference

Construct an EOS with an arbitrary choice for the terms in the symmetry energy that are quartic in the isospin asymmetry.

```
#include <sym4_eos.h>
```

Inheritance diagram for [sym4\\_eos](#):





### 14.43.1 Detailed Description

#### References:

Created for [Steiner06](#).

Definition at line 198 of file sym4\_eos.h.

#### Public Member Functions

- int [set\\_base\\_eos](#) (sym4\_eos\_base &seb)  
*Set the base equation of state.*
- virtual int [test\\_eos](#) (fermion &ne, fermion &pr, thermo &lth)  
*Test the equation of state.*
- virtual int [calc\\_e](#) (fermion &ne, fermion &pr, thermo &lth)  
*Equation of state as a function of density.*

#### Data Fields

- double [alpha](#)  
*The strength of the quartic terms.*

#### Protected Attributes

- [sym4\\_eos\\_base \\* sp](#)  
*The base equation of state to use.*

### 14.43.2 Member Function Documentation

#### 14.43.2.1 virtual int sym4\_eos::test\_eos ( fermion & ne, fermion & pr, thermo & lth ) [virtual]

This compares the chemical potentials from calc\_e\_sep() to their finite-difference approximations in order to ensure that the separation into potential and kinetic parts is done properly.

The documentation for this class was generated from the following file:

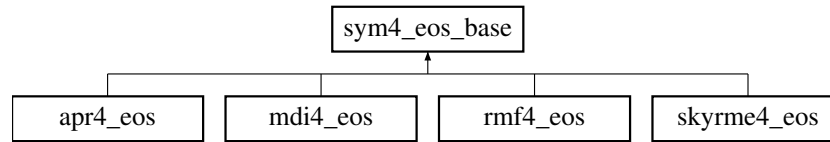
- sym4\_eos.h

## 14.44 sym4\_eos\_base Class Reference

A class to compute quartic contributions to the symmetry energy [abstract base].

```
#include <sym4_eos.h>
```

Inheritance diagram for sym4\_eos\_base:



#### 14.44.1 Detailed Description

The standard usage is that a child class implements the virtual function `calc_e_sep()` which is then used by `calc_e_alpha()` and `calc_muhat()`. These functions are employed by `sym4_eos` to compute the EOS for an arbitrary dependence of the symmetry energy on the isospin.

---

#### References:

Created for [Steiner06](#).

**Bug** Testing was disabled in HDF conversion. Fix this.

Definition at line 55 of file `sym4_eos.h`.

#### Public Member Functions

- virtual int `calc_e_alpha` (**fermion** &ne, **fermion** &pr, **thermo** &lth, double &alphak, double &alphap, double &alphan, double &diff\_kin, double &diff\_pot, double &ed\_kin\_nuc, double &ed\_pot\_nuc)  
*Compute alpha at the specified density.*
- virtual double `calc_muhat` (**fermion** &ne, **fermion** &pr)  
*Compute  $\mu$ , the out-of-whack parameter.*
- virtual int `calc_e_sep` (**fermion** &ne, **fermion** &pr, double &ed\_kin, double &ed\_pot, double &mu\_n\_kin, double &mu\_p\_kin, double &mu\_n\_pot, double &mu\_p\_pot)=0  
*Compute the potential and kinetic parts separately (to be overwritten in children)*

#### Protected Attributes

- fermion** e  
*An electron for the computation of the  $\mu$ .*
- fermion\_zerot `fzt2`  
*Desc.*

The documentation for this class was generated from the following file:

- `sym4_eos.h`

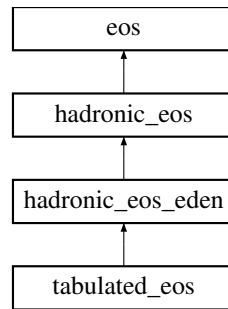
## 14.45 tabulated\_eos Class Reference

EOS from a table.

```
#include <tabulated_eos.h>
```

Inheritance diagram for tabulated\_eos:

---



### 14.45.1 Detailed Description

This assumes a symmetry energy which depends quadratically on the isospin asymmetry in order to construct an EOS from a table of baryon density and energy per baryon for both nuclear and pure neutron matter.

Note: If using a tabulated EOS to compute derivatives (like the compressibility which effectively requires a second derivative), it is important to tabulated the EOS precisely enough to ensure that the derivatives are accurate. In the case of ensuring that the compressibility at saturation density is well reproduced, I have needed the EOS to be specified with at least 6 digits of precision on a grid at least as small as  $0.002 \text{ fm}^{-3}$ .

Definition at line 52 of file tabulated\_eos.h.

#### Public Member Functions

- virtual int **calc\_e** (**fermion** &ne, **fermion** &pr, **thermo** &th)  
*Equation of state as a function of density.*
- template<class vec\_t >  
int **set\_eos** (size\_t n, vec\_t &rho, vec\_t &Enuc, vec\_t &Eneut)  
*Set the EOS through vectors specifying the densities and energies.*
- template<class vec\_t >  
int **set\_eos** (size\_t n\_nuc, vec\_t &rho\_nuc, vec\_t &E\_nuc, size\_t n\_neut, vec\_t &rho\_neut, vec\_t &E\_neut)  
*Set the EOS through vectors specifying the densities and energies.*
- **table** & **get\_nuc\_table** ()  
*Return the internal table.*
- **table** & **get\_neut\_table** ()  
*Return the internal table.*

#### Protected Member Functions

- int **free\_table** ()  
*Free the table memory.*

#### Protected Attributes

- bool **table\_alloc**  
*True if the table has been allocated.*
- bool **one\_table**  
*If true, then tnuc and tneut point to the same table.*

#### The EOS tables

- **table** \* **tnuc**
- **table** \* **tneut**

Strings for the column names

- std::string **srho\_nuc**
- std::string **srho\_neut**
- std::string **snuc**
- std::string **sneut**

The documentation for this class was generated from the following file:

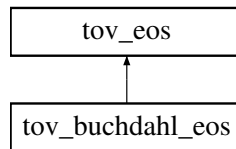
- tabulated\_eos.h

## 14.46 tov\_buchdahl\_eos Class Reference

The Buchdahl EOS for the TOV solver.

```
#include <tov_eos.h>
```

Inheritance diagram for tov\_buchdahl\_eos:



### 14.46.1 Detailed Description

Given the eos

$$\rho = 12\sqrt{p_*P} - 5P$$

the TOV equation has an analytical solution

$$R = (1 - \beta) \sqrt{\frac{\pi}{288p_*G(1 - 2\beta)}}$$

where  $\beta = GM/R$ .

The central pressure and energy density are

$$P_c = 36p_*\beta^2$$

$$\rho_c = 72p_*\beta(1 - 5\beta/2)$$

Physical solutions are obtained only for  $P < 25p_*/144$  and  $\beta < 1/6$ .

Based on [Lattimer01](#).

**Idea for Future** Figure out what to do with the buchfun() function

Definition at line 363 of file tov\_eos.h.

### Public Member Functions

- virtual int [get\\_edens](#) (double P, double &e, double &nb)  
*Given the pressure, produce the energy and number densities.*
- virtual int [get\\_aux](#) (double P, size\_t &np, **ovector\_base** &auxp)  
*Given the pressure, produce all the remaining quantities.*
- virtual int [get\\_names](#) (size\_t &np, std::vector< std::string > &pnames)  
*Fill a list with strings for the names of the remaining quantities.*

## Data Fields

- double `Pstar`  
*The parameter with units of pressure in units of solar masses per km cubed (default value  $3.2 \times 10^{-5}$ )*

## 14.46.2 Member Function Documentation

14.46.2.1 `virtual int tov_buchdahl_eos::get_eden ( double P, double &e, double &nb ) [inline, virtual]`

If the baryon density is not specified, it should be set to zero or `baryon_column` should be set to false

Implements `tov_eos`.

Definition at line 384 of file `tov_eos.h`.

The documentation for this class was generated from the following file:

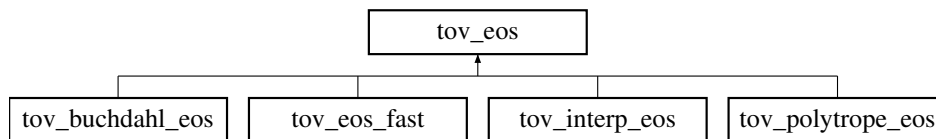
- `tov_eos.h`

14.47 `tov_eos` Class Reference

A EOS base class for the TOV solver.

```
#include <tov_eos.h>
```

Inheritance diagram for `tov_eos`:



## 14.47.1 Detailed Description

**Todo** Fix `read_table_file` and maybe `set_low_density_eos()`.

Definition at line 53 of file `tov_eos.h`.

## Public Member Functions

- virtual int `get_eden` (double *P*, double &*e*, double &*nb*)=0  
*Given the pressure, produce the energy and number densities.*
- virtual int `get_aux` (double *P*, size\_t &*np*, **ovector\_base** &*auxp*)  
*Given the pressure, produce all the remaining quantities.*
- virtual int `get_names_units` (size\_t &*np*, std::vector< std::string > &*pnames*, std::vector< std::string > &*punits*)  
*Fill a list with strings for the names of the remaining quantities.*

## Data Fields

- int `verbose`  
*Control for output (default 1)*
- bool `baryon_column`  
*Set to true if the baryon density is provided in the EOS (default false)*

## Protected Attributes

- double `mev_kg`  
*To convert MeV to kilograms.*
- double `mev_per_fm3_msun_km3`  
*To convert MeV/fm<sup>3</sup> to solar masses per cubic kilometer.*

## 14.47.2 Member Function Documentation

14.47.2.1 `virtual int tov_eos::get_eden ( double P, double &e, double &nb )` [pure virtual]

The arguments *P* and *e* should always be in  $M_{\odot}/\text{km}^3$ .

If the baryon density is not specified, it should be set to zero or `baryon_column` should be set to false.

Implemented in `tov_polytrope_eos`, `tov_buchdahl_eos`, `tov_interp_eos`, and `tov_eos_fast`.

14.47.2.2 `virtual int tov_eos::get_aux ( double P, size_t &np, ovector_base &auxp )` [inline, virtual]

The argument *P* should always be in  $M_{\odot}/\text{km}^3$ .

Reimplemented in `tov_polytrope_eos`, `tov_buchdahl_eos`, and `tov_interp_eos`.

Definition at line 93 of file `tov_eos.h`.

The documentation for this class was generated from the following file:

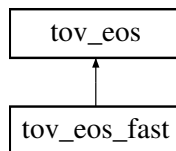
- `tov_eos.h`

14.48 `tov_eos_fast` Class Reference

An EOS.

```
#include <tov_eos_fast.h>
```

Inheritance diagram for `tov_eos_fast`:



## 14.48.1 Detailed Description

Definition at line 38 of file `tov_eos_fast.h`.

## Public Member Functions

- virtual int `get_eden` (double *pres*, double &*ed*, double &*nb*)  
*Given the pressure, produce the energy and number densities.*
- int `set_crust_core` (size\_t *n\_crust*, ovector\_base &*crust\_e*, ovector\_base &*crust\_p*, ovector\_base &*crust\_nb*, size\_t *n\_core*, ovector\_base &*core\_e*, ovector\_base &*core\_p*, ovector\_base &*core\_nb*)  
*Desc.*

## Data Fields

- bool `check_eos`  
*If true, check the EOS for stability and ordering (default true)*

## Protected Member Functions

- void `interp` (const `ovector_base` &x, const `ovector_base` &y, double xx, double &yy, int n1, int n2)  
*Linear EOS interpolation.*

## Protected Attributes

- size\_t `nlines`  
*Desc.*
- `ovector` `edv`  
*Desc.*
- `ovector` `prv`  
*Desc.*
- `ovector` `nbv`  
*Desc.*
- bool `eos_set`  
*Desc.*

The documentation for this class was generated from the following file:

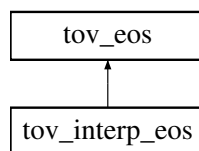
- `tov_eos_fast.h`

14.49 `tov_interp_eos` Class Reference

An EOS for the TOV solver using simple linear interpolation and a default low-density EOS.

```
#include <tov_eos.h>
```

Inheritance diagram for `tov_interp_eos`:



## 14.49.1 Detailed Description

Internally, energy and pressure are stored in units of solar masses per cubic kilometer and baryon density is stored in units of  $\text{fm}^{-3}$ . The user-specified EOS table is left as is, and unit conversion is performed as needed in `get_eden()` and other functions so that results are returned in the units specified by `set_units()`.

## Note

This stores a pointer to the user-specified table, so if that pointer becomes invalid, the interpolation will fail.

The function `set_units()` needs to be called before either of the functions `get_eden()` or `get_eden_ld()` are called. The function `set_units()` may be called after calling either the `read_table()` functions or the `set_low_density_eos()` function.

**Todo** Warn that the pressure in the low-density eos is not strictly increasing! (see at  $P=4.3e-10$ )

It might be useful to exit more gracefully when non-finite values are obtained in interpolation, analogous to the `err_nonconv` mechanism elsewhere.

Definition at line 141 of file `tov_eos.h`.

### Public Member Functions

- virtual int `get_ed` (double `pres`, double &`ed`, double &`nb`)  
*Given the pressure, produce the energy and number densities.*
- virtual int `get_ed_user` (double `pres`, double &`ed`, double &`nb`)  
*Given the pressure, produce the energy and number densities from the user-specified EOS.*
- virtual int `get_ed_ld` (double `pres`, double &`ed`, double &`nb`)  
*Given the pressure, produce the energy and number densities from the low-density EOS.*
- virtual int `get_aux` (double `P`, size\_t &`nv`, **ovector\_base** &`auxp`)  
*Given the pressure, produce all the remaining quantities.*
- virtual int `get_names_units` (size\_t &`np`, std::vector< std::string > &`pnames`, std::vector< std::string > &`punits`)  
*Fill a list with strings for the names of the remaining quantities.*
- int `read_table` (**table\_units** &`eosat`, std::string `s_cole`, std::string `s_colp`, std::string `s_colnb=""`)  
*Specify the EOS through a table.*
- int `set_low_density_eos` (std::string `s_nvpath`, int `s_nvcole=0`, int `s_nvcolp=1`, int `s_nvcolnb=2`)  
*Set the low-density EOS.*
- int `default_low_dens_eos` ()  
*Default low-density EOS.*
- int `no_low_dens_eos` ()  
*Compute with no low-density EOS.*
- int `set_units` (double `s_efactor`, double `s_pfactor`, double `s_nfactor`)  
*Set the units of the user-specified EOS.*
- int `set_units` (std::string `leunits=""`, std::string `lpunits=""`, std::string `lnunits=""`)  
*Set the units of the user-specified EOS.*
- int `get_transition` (double &`plow`, double &`ptrans`, double &`phi`)  
*Return limiting and transition pressures.*
- int `set_transition` (double `ptrans`, double `pw`)  
*Set the transition pressure and "width".*

### Protected Member Functions

- int `check_eos` ()  
*Check that the EOS is valid.*
- void `interp` (const **ovector\_base** &`x`, const **ovector\_base** &`y`, double `xx`, double &`yy`, int `n1`, int `n2`)  
*Linear EOS interpolation.*

### Protected Attributes

#### Low-density EOS

- bool `ldeos`  
*true if we are using the low-density eos (false)*
- bool `ldeos_read`  
*Low-density EOS switch.*
- std::string `ldpath`  
*the path to the low-density EOS*
- int `ldcole`  
*column in low-density eos for energy density*
- int `ldcolp`  
*column in low-density eos for pressure*
- int `ldcolnb`  
*column in low-density eos for baryon density*



- **table\_units** \* [ld\\_eos](#)  
*file containing low-density EOS*
- double [presld](#)  
*Highest pressure in low-density EOS (in  $M_{\odot}/\text{km}^3$ )*
- double [eld](#)  
*Highest energy density in low-density EOS (in  $M_{\odot}/\text{km}^3$ )*
- double [nbld](#)  
*Highest baryon density in low-density EOS (in  $M_{\odot}/\text{km}^3$ )*
- double [prest](#)  
*Transition pressure (in  $M_{\odot}/\text{km}^3$ )*
- double [pwidth](#)  
*Transition width (unitless)*

#### User EOS

- **table\_units** \* [eost](#)  
*file containing eos*
- int [nfile](#)  
*number of lines in eos file*
- int [cole](#)  
*column for energy density in eos file*
- int [colp](#)  
*column for pressure in eos file*
- int [coln](#)  
*column for baryon density in eos file*
- bool [eos\\_read](#)  
*True if an EOS has been specified.*

#### Units

- std::string [eunits](#)  
*Units for energy density.*
- std::string [punits](#)  
*Units for pressure.*
- std::string [nunits](#)  
*Units for baryon density.*
- double [efactor](#)  
*unit conversion factor for energy density (default 1.0)*
- double [pfactor](#)  
*unit conversion factor for pressure (default 1.0)*
- double [nfactor](#)  
*unit conversion factor for baryon density (default 1.0)*

### 14.49.2 Member Function Documentation

14.49.2.1 `virtual int tov_interp_eos::get_edn ( double pres, double & ed, double & nb )` [virtual]

The arguments `P` and `e` should always be in  $M_{\odot}/\text{km}^3$ .

If the baryon density is not specified, it should be set to zero or `baryon_column` should be set to false.

Implements [tov\\_eos](#).

14.49.2.2 `virtual int tov_interp_eos::get_edn_user ( double pres, double & ed, double & nb )` [virtual]

The arguments `pres`, `ed`, and `nb` should be in the user-specified unit system.

14.49.2.3 `virtual int tov_interp_eos::get_edn_ld ( double pres, double & ed, double & nb )` [virtual]

The arguments `pres`, `ed`, and `nb` should be in the user-specified unit system.

14.49.2.4 `virtual int tov_interp_eos::get_aux ( double P, size_t & nv, ovector_base & auxp )` [virtual]

The argument *P* should always be in  $M_{\odot}/\text{km}^3$ .

Reimplemented from [tov\\_eos](#).

14.49.2.5 `int tov_interp_eos::get_transition ( double & plow, double & ptrans, double & phi )`

Returns, in order:

- the highest pressure in the low-density EOS
- the transition pressure
- the lowest pressure in the high-density EOS

14.49.2.6 `int tov_interp_eos::set_transition ( double ptrans, double pw )`

Sets the transition pressure and the width (specified as a number greater than unity in *pw*) of the transition between the two EOSs. The transition is done smoothly using linear interpolation between  $P = \text{ptrans}/\text{pmathrmpw}$  and  $P = \text{ptrans} \times \text{pmathrmpw}$ .

### 14.49.3 Field Documentation

14.49.3.1 `bool tov_interp_eos::ldeos_read` [protected]

This is `true` if the `ldeos` has been read by `set_ldeos`. This is useful, since then we know whether or not we need to free the memory for the LD EOS in the destructor

Definition at line 255 of file `tov_eos.h`.

The documentation for this class was generated from the following file:

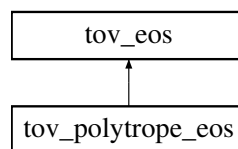
- `tov_eos.h`

## 14.50 `tov_polytrope_eos` Class Reference

Standard polytropic EOS  $p = K\rho^{1+1/n}$ .

```
#include <tov_eos.h>
```

Inheritance diagram for `tov_polytrope_eos`:



### 14.50.1 Detailed Description

Any units are permissible, but if this is to be used with [tov\\_solve](#), then the units of *K* must be consistent with the units set in [tov\\_solve::set\\_units\(\)](#).

Definition at line 426 of file `tov_eos.h`.

## Public Member Functions

- virtual int `get_eden` (double P, double &e, double &nb)  
*Given the pressure, produce the energy and number densities.*
- virtual int `get_aux` (double P, size\_t &np, **ovector\_base** &auxp)  
*Given the pressure, produce all the remaining quantities.*
- virtual int `get_names` (size\_t &np, std::vector< std::string > &pnames)  
*Fill a list with strings for the names of the remaining quantities.*

## Data Fields

- double `K`  
*Coefficient (default 1.0)*
- double `n`  
*Index (default 3.0)*

## 14.50.2 Member Function Documentation

14.50.2.1 virtual int `to_solve::get_eden` ( double P, double & e, double & nb ) [inline, virtual]

If the baryon density is not specified, it should be set to zero or `baryon_column` should be set to false

Implements `to_solve`.

Definition at line 449 of file `to_solve.h`.

The documentation for this class was generated from the following file:

- `to_solve.h`

14.51 `to_solve` Class Reference

Class to solve the Tolman-Oppenheimer-Volkov equations.

```
#include <to_solve.h>
```

## 14.51.1 Detailed Description

**Mathematical background:**

In units where  $c = 1$ , the TOV equations (i.e. Einstein's equations for a static spherically symmetric object) are

$$\frac{dm}{dr} = 4\pi r^2 \epsilon$$

$$\frac{dP}{dr} = -\frac{G\epsilon m}{r^2} \left(1 + \frac{P}{\epsilon}\right) \left(1 + \frac{4\pi P r^3}{m}\right) \left(1 - \frac{2Gm}{r}\right)^{-1}$$

where  $r$  is the radial coordinate,  $m(r)$  is the gravitational mass enclosed within a radius  $r$ , and  $\epsilon(r)$  and  $P(r)$  are the energy density and pressure at  $r$ , and  $G$  is the gravitational constant. The boundary conditions are  $m(r=0) = 0$  the condition  $P(r=R) = 0$  for some fixed radius  $R$ . These boundary conditions give a series of solutions to the TOV equations as a function of the radius, although they do not necessarily have a solution for all radii.

The gravitational mass is given by

$$M_G = \int_0^R 4\pi r^2 \epsilon dr$$

The gravitational potential,  $\Phi(r)$  can be determined from

$$\frac{d\Phi}{dr} = -\frac{1}{\epsilon} \frac{dP}{dr} \left(1 + \frac{P}{\epsilon}\right)^{-1}$$

The proper boundary condition for the gravitational potential is

$$\Phi(r=R) = \frac{1}{2} \ln \left(1 - \frac{2GM}{R}\right)$$

which ensures that  $\phi(r) \rightarrow 0$  as  $r \rightarrow \infty$ .

The surface gravity is

$$g = \frac{Gm}{r^2} \left(1 - \frac{2Gm}{r}\right)^{-1/2}$$

which is given in inverse kilometers and the redshift is

$$z = \left(1 - \frac{2Gm}{r}\right)^{-1/2} - 1$$

which is unitless.

The baryonic mass is typically defined by

$$M_B = \int_0^R 4\pi r^2 n_B m_B \left(1 - \frac{2Gm}{r}\right)^{-1/2} dr$$

where  $n_B(r)$  is the baryon number density at radius  $r$  and  $m_B$  is the mass of one baryon. Then the "binding energy" of the neutron star is defined to be  $M_B - M_G$ . If you can specify the product  $n_B m_B$  in the EOS (analogous to the rest mass energy density), it will compute the associated baryonic mass for you.

In the case of slow rigid rotation with angular velocity  $\Omega$ , the moment of inertia is

$$I = \frac{8\pi}{3} \int_0^R dr r^4 (\epsilon + P) e^{\Phi} \left(\frac{\bar{\omega}}{\Omega}\right) \left(1 - \frac{2Gm}{r}\right)^{-1/2}$$

where the function  $\bar{\omega}(r)$  is the solution of

$$\frac{d}{dr} \left( r^4 j \frac{d\bar{\omega}}{dr} \right) + 4r^3 \frac{dj}{dr} \bar{\omega} = 0$$

and the function  $j(r)$  is defined by

$$j = \left(1 - \frac{2Gm}{r}\right) e^{-\Phi}$$

The boundary conditions for  $\bar{\omega}$  are  $d\bar{\omega}/dr = 0$  at  $r = 0$  and

$$\bar{\omega}(R) = \Omega - \left(\frac{R}{3}\right) \left(\frac{d\bar{\omega}}{dr}\right)_{r=R}.$$

---

### General usage notes:

The equation of state may be changed at any time, by specifying the appropriate `tov_eos` object

Screen output:

- `verbose=0` - Nothing
  - `verbose=1` - Basic information
-

- `verbose=2` - For each profile computation, report solution information at every kilometer.
- `verbose=3` - Report profile information at every 20 grid points. A keypress is required after each profile.

---

### Output tables:

The functions `fixed()` and `max()` produce output tables which represent the profile of the neutron star of the requested mass. The columns are

- `gm`, the enclosed gravitational mass in solar masses
- `r`, the radial coordinate in km
- `gp`, the gravitational potential (unitless) when `calcgpot` is true
- `bm`, the baryonic mass in solar masses (when `tov_eos::baryon_column` is true).
- `pr`, the pressure
- `ed`, the energy density
- `nb`, the baryon density (if `tov_eos::baryon_column` is true)
- `sg`, the local surface gravity (in  $\text{g}/\text{cm}^2$  )
- `rs`, the local redshift (unitless),
- `dmdr`, the derivative of the enclosed gravitational mass in  $M_{\odot}/\text{km}$
- `dlogpdr`, the derivative of the natural logarithm of the pressure
- `dgpdr`, the derivative of the gravitational potential in  $1/\text{km}$  (if `calcgpot` is true)
- `dbmdr`, the derivative of the enclosed baryonic mass (if `tov_eos::baryon_column` is true).

The remaining columns are given by the user-defined columns from the equation of state as determined by `tov_eos::get_names_units()` and `tov_eos::get_aux()`.

The function `mvsvr()` produces a different kind of output table corresponding to the mass versus radius curve. Some points on the curve may correspond to unstable branches.

- `gm`, the total gravitational mass in solar masses
  - `r`, the radius in km
  - `gp`, the gravitational potential at the surface (unitless) when `calcgpot` is true
  - `bm`, total the baryonic mass in solar masses (when `tov_eos::baryon_column` is true).
  - `pr`, the central pressure
  - `ed`, the central energy density
  - `nb`, the central baryon density (if `tov_eos::baryon_column` is true)
  - `sg`, the surface gravity (in  $\text{g}/\text{cm}^2$  )
  - `rs`, the redshift at the surface,
  - `dmdr`, the derivative of the gravitational mass
  - `dlogpdr`, the derivative of the natural logarithm of the pressure
  - `dgpdr`, the derivative of the gravitational potential in  $1/\text{km}$  (if `calcgpot` is true)
-

- `dbmdr`, the derivative of the enclosed baryonic mass (if `to_solve::baryon_column` is true).

The remaining columns are given by the user-defined columns from the equation of state as determined by `to_solve::get_names_units()` and `to_solve::get_aux()`.

---

### Accuracy:

The present code, as demonstrated in the tests, gives the correct central pressure and energy density of the analytical solution by Buchdahl to within less than 1 part in  $10^8$ .

---

### Other details and todos:

#### Note

The function `star_fun()` returns `gsl_efailed` without calling the error handler in the case that the solver can recover gracefully from, for example, a negative pressure.

#### Todo

- baryon mass doesn't work for `fixed()` (This may be fixed. We should make sure it's tested.)
- Combine `maxoutsize` and `kmax`?
- Document column naming issues
- Document surface gravity and redshift
- Standardize `xmev_kg`, etc.
- Use `convert_units`?
- Double check that `fixed()` doesn't give a solution on the unstable branch
- Ensure that this class copies over the units of the user-specified columns to the table

---

Definition at line 243 of file `to_solve.h`.

### Public Member Functions

#### Basic operation

- `int set_eos (to_solve &ter)`  
*Set the EOS to use.*
- `int set_units (double s_efactor=1.0, double s_pfactor=1.0, double s_nbfactor=1.0)`  
*Set units.*
- `int set_units (std::string eunits="", std::string punits="", std::string nunits="")`  
*Set units.*
- `virtual int mvsr ()`  
*Calculate the mass vs. radius curve.*
- `virtual int fixed (double d_tmass)`  
*Calculate the profile of a star with fixed mass.*
- `virtual int max ()`  
*Calculate the profile of the maximum mass star.*
- `o2_shared_ptr<table_units>::type get_results ()`  
*Return the results data table.*
- `int solution_check ()`  
*Check the solution (unfinished)*

#### Control numerical methods

- `int set_kmax (int s_maxoutsize=400, int s_kmax=80000)`  
*Set maximum storage for integration.*
  - `int set_mroot (mroot<mm_funct<>> &s_mrp)`  
*Set solver.*
  - `int set_minimize (minimize<funct> &s_mp)`  
*Set minimizer.*
  - `int set_stepper (adapt_step<ode_funct<>> &sap)`  
*Set the adaptive stepper.*
-

## Data Fields

- `gsl_min_brent`< `funct` > `def_min`  
*The default minimizer.*
- `gsl_mroot_hybrids`< `mm_funct`<> > `def_solver`  
*The default solver.*
- `gsl_astepp`< `ode_funct`<> > `def_stepper`  
*The default adaptive stepper.*
- bool `compute_ang_vel`  
*If true, compute the angular velocity (default false)*
- double `cap_omega`  
*The angular velocity.*

## Basic properties

- double `mass`  
*mass*
- double `rad`  
*radius*
- double `bmass`  
*baryonic mass*
- double `gpot`  
*gravitational potential*

## Solution parameters

*These parameters can be changed at any time.*

- bool `generel`  
*Use general relativistic version (default true)*
- bool `calcgpot`  
*calculate the gravitational potential and the enclosed baryon mass (default false)*
- double `hmin`  
*smallest allowed radial stepsize (default 1.0e-4)*
- double `hmax`  
*largest allowed radial stepsize (default 0.05)*
- double `hstart`  
*initial radial stepsize (default 4.0e-3)*
- int `verbose`  
*control for output (default 1)*
- double `maxradius`  
*maximum radius for integration in km (default 60)*

## Mass versus radius parameters

- double `prbegin`  
*Beginning pressure (default 7.0e-7)*
- double `prend`  
*Ending pressure (default 8.0e-3)*
- double `princ`  
*Increment for pressure (default 1.1)*
- bool `logmode`  
*Use 'princ' as a multiplier, not an additive increment (default true)*
- double `prguess`  
*Guess for central pressure in solar masses per km<sup>3</sup> (default  $5.2 \times 10^{-5}$ )*
- double `max_begin`  
*Beginning pressure for maximum mass guess (default 7.0e-5)*
- double `max_end`  
*Ending pressure for maximum mass guess (default 5.0e-3)*
- double `max_inc`  
*Increment for pressure for maximum mass guess (default 1.3)*

## Protected Member Functions

- int `make_unique_name` (std::string &col, std::vector< std::string > &cnames)  
*Ensure col does not match strings in cnames.*
- int `ang_vel` ()  
*Compute the angular velocity.*
- virtual int `derivs` (double x, size\_t nv, const **ovector\_base** &y, **ovector\_base** &dydx)  
*The ODE step function.*
- virtual int `derivs_ang_vel` (double x, size\_t nv, const **ovector\_base** &y, **ovector\_base** &dydx)  
*The ODE step function for the angular velocity.*
- virtual int `profile_out` (double xx)  
*Construct a stellar profile.*
- virtual double `maxfun` (double maxx)  
*The minimizer function to compute the maximum mass.*
- virtual int `starfun` (size\_t ndvar, const **ovector\_base** &ndx, **ovector\_base** &ndy)  
*The solver function to compute the stellar profile.*

## Protected Attributes

- double `mev_kg`  
*To convert MeV to kilograms.*
- double `mev_per_fm3_msun_km3`  
*To convert MeV / fm<sup>3</sup> to solar masses per cubic kilometer.*
- double `tmass`  
*Target mass.*
- `to_solve::eos * te`  
*The EOS.*
- bool `eos_set`  
*True if the EOS has been set.*
- int `kmax`  
*maximum storage size (default 40000)*
- int `maxoutsize`  
*maximum size of output file (default 400)*
- double `presmin`  
*Smallest allowed pressure for integration (default: -100)*
- `sm_interp smi`  
*Interpolation object for `derivs_ang_vel()`*
- `o2_shared_ptr< table_units >::type out_table`  
*The output table.*
- `mroot< mm_funct<> > * mroot_ptr`  
*The solver.*
- `minimize< funct > * min_ptr`  
*The minimizer.*
- `adapt_step< ode_funct<> > * as_ptr`  
*The adaptive stepper.*
- double `schwarz_km`  
*The schwarzschild radius in km.*

## User EOS units

- std::string `eunits`  
*Units for energy density.*
- std::string `punits`  
*Units for pressure.*
- std::string `nunits`  
*Units for baryon density.*
- double `efactor`  
*unit conversion factor for energy density (default 1.0)*
- double `pfactor`  
*unit conversion factor for pressure (default 1.0)*
- double `nfactor`  
*unit conversion factor for baryon density (default 1.0)*



## Integration storage

- **ovector** `rky` [6]  
*ODE functions.*
- **ovector** `rkx`  
*Radial coordinate (in kilometers)*
- **ovector** `rkdydx` [6]  
*The derivatives of the ODE functions.*

## 14.51.2 Member Function Documentation

14.51.2.1 `int to_v_solve::make_unique_name ( std::string & col, std::vector< std::string > & cnames )` [protected]

Underscores are added to `col` until it matches none of the strings in `cnames`.

14.51.2.2 `virtual int to_v_solve::profile_out ( double xx )` [protected, virtual]

This function constructs a stellar profile in `out_table` from the information in `rkx` and `rky`.

14.51.2.3 `int to_v_solve::set_units ( std::string eunits = " ", std::string punits = " ", std::string nunits = " " )`

Valid entries for the units of energy density and pressure are:

- "g/cm<sup>3</sup>"
- "erg/cm<sup>3</sup>"
- "MeV/fm<sup>3</sup>"
- "1/fm<sup>4</sup>"
- "Msun/km<sup>3</sup>" (i.e. solar masses per cubic kilometer)

Valid entries for the units of baryon density are:

- "1/m<sup>3</sup>"
- "1/cm<sup>3</sup>"
- "1/fm<sup>3</sup>"

14.51.2.4 `int to_v_solve::set_kmax ( int s_maxoutsize = 400, int s_kmax = 80000 )`

The variable `s_kmax` is the maximum number of radial integration stepsk while `s_maxoutsize` is the maximum number of points that will be output for any profile.

If `s_kmax` is less than zero, there is no limit on the number of radial steps.

## 14.51.3 Field Documentation

14.51.3.1 `double to_v_solve::tmass` [protected]

Negative values to indicate a mass measured relative to the maximum mass. For example, if the EOS has a maximum mass of 2.0, then -0.15 will give the profile of a 1.85 solar mass star.

Definition at line 261 of file `to_v_solve.h`.

14.51.3.2 `double to_v_solve::presmin` [protected]

This quantity can't be much smaller than -100 since we need to compute numbers near  $e^{-presmin}$

Definition at line 303 of file `to_v_solve.h`.

---

### 14.51.3.3 `ovector tov_solve::rky[6]` [protected]

The vector `rky[0]` is the gravitational mass in solar masses, and the vector `rky[1]` is the natural logarithm of the pressure in solar masses per cubic kilometer. When `calcpot` is true, the next vector in this list is the gravitational potential (which is unitless), and when `tov_eos::baryon_column` is true, the next vector in this list is the baryonic mass in solar masses.

Definition at line 317 of file `tov_solve.h`.

### 14.51.3.4 `double tov_solve::prguess`

This guess is used in the function `fixed()`.

Definition at line 427 of file `tov_solve.h`.

The documentation for this class was generated from the following file:

- `tov_solve.h`

## 15 File Documentation

### 15.1 `hdf_eos_io.h` File Reference

File for HDF input of the O<sub>2</sub>scl `skyrme_eos` and `rmf_eos` data files.

```
#include <hdf5.h> #include <o2scl/constants.h> #include <o2scl/hdf_file.h> #include <o2scl/lib-
_settings.h> #include <o2scl/skyrme_eos.h> #include <o2scl/rmf_eos.h>
```

#### 15.1.1 Detailed Description

Definition in file `hdf_eos_io.h`.

#### Functions

- `int rmf_load (rmf_eos &rmf, std::string model, bool external=false)`  
*Input a `rmf_eos` object from an HDF file.*
- `int skyrme_load (skyrme_eos &sk, std::string model, bool external=false)`  
*Input a `skyrme_eos` object from an HDF file.*