

O₂scl_part - Particle Sub-Library for O₂scl

Version 0.905

Copyright © 2006, 2007, 2008, 2009 Andrew W. Steiner

Permission is granted to copy, distribute and/or modify this document under the terms of the GNU Free Documentation License, Version 1.2 or any later version published by the Free Software Foundation; with no Invariant Sections, no Front-Cover Texts, and no Back-Cover Texts. A copy of the license is included in the section entitled “License Information”.

Contents

1	Main Page	1
1.1	Quick Reference to User's Guide	1
1.2	Particles	1
1.3	Atomic nuclei	2
1.4	Example source code	3
1.5	Bibliography	5
2	Ideas for future development	5
3	Todo List	6
4	Data Structure Documentation	7
4.1	ame_entry Struct Reference	7
4.2	ame_entry03_io_type Class Reference	8
4.3	ame_entry95_io_type Class Reference	8
4.4	ame_mass Class Reference	9
4.5	boson Class Reference	11
4.6	classical Class Reference	12
4.7	deriv_part Class Reference	13
4.8	eff_boson Class Reference	16
4.9	eff_fermion Class Reference	18
4.10	eff_quark Class Reference	21
4.11	fermion Class Reference	23
4.12	fermion_T Class Reference	24
4.13	full_dist Class Reference	27
4.14	hfb_mass Class Reference	27
4.15	hfb_mass_entry Struct Reference	29
4.16	mass_fit Class Reference	29
4.17	mnmsk_mass Class Reference	30
4.18	mnmsk_mass_entry Struct Reference	32
4.19	mnmsk_mass_exp Class Reference	34
4.20	nonrel_fermion Class Reference	34
4.21	nonrel_fermion_zerot Class Reference	36
4.22	nuclear_dist Class Reference	36
4.23	nuclear_dist::iterator Class Reference	37
4.24	nuclear_mass Class Reference	38
4.25	nuclear_mass_cont Class Reference	41
4.26	nuclear_mass_disc Class Reference	41
4.27	nuclear_mass_fit Class Reference	42
4.28	nuclear_mass_info Class Reference	43
4.29	nuclear_mass_info::string_less_than Struct Reference	45
4.30	nuclear_reaction Class Reference	45
4.31	nucleus Class Reference	46
4.32	part Class Reference	47
4.33	quark Class Reference	49
4.34	reaction_lib Class Reference	50
4.35	rel_boson Class Reference	51
4.36	rel_fermion Class Reference	52
4.37	semi_empirical_mass Class Reference	55
4.38	simple_dist Class Reference	57
4.39	sn_classical Class Reference	58
4.40	sn_fermion Class Reference	59
4.41	sn_nr_fermion Class Reference	63
4.42	thermo Class Reference	66

5 File Documentation	67
5.1 part.h File Reference	67

1 Main Page

1.1 Quick Reference to User's Guide

- [Particles](#)
- [Atomic nuclei](#)
- [Example source code](#)
- [Bibliography](#)

1.2 Particles

These classes in the library `O2scl_part` calculate the thermodynamic properties of interacting and non-interacting quantum and classical particles.

The class `part` is the basic structure for a particle:

- `part::m` - mass
- `part::g` - degeneracy factor (e.g. $2j + 1$)
- `part::n` - number density
- `part::ed` - energy density
- `part::pr` - pressure
- `part::en` - entropy density
- `part::ms` - effective mass, i.e. m^*
- `part::nu` - effective chemical potential, i.e. ν
- `part::inc_rest_mass` - True if the rest mass is included
- `part::non_interacting` - False if the particle includes interactions

The data members `part::ms` and `part::nu` allow one to specify modifications to the mass and the chemical potential due to interactions. This allows one to calculate the properties of particle due to interactions so long as the basic form of the free-particle dispersion relation is unchanged, i.e.

$$\sqrt{k^2 + m^2} - \mu \rightarrow \sqrt{k^2 + m^{*2}} - \nu$$

If the particle is non-interacting, then `part::nu` and `part::ms` are sometimes used by `O2scl_part` functions for temporary storage.

If `part::inc_rest_mass` is `true` (as is the default in all of the classes except `nucleus`), then all functions include the rest mass energy density in the energy density, the "mu" functions expect that the rest mass is included in `part::mu` or `part::nu` as input and the "density" functions output `part::mu` or `part::nu` including the rest mass.

When `part::inc_rest_mass` is true, antiparticles are implemented by choosing the antiparticle chemical potential to be $-\mu$, and when `inc_rest_mass` is false, antiparticles are implemented by choosing the chemical potential of the antiparticles to be $-\mu - 2m$.

The thermodynamic identity used to compute the pressure for interacting particles is

$$P = -\varepsilon + sT + \nu n$$

where `part::nu` is used. This way, the particle class doesn't need to know about the structure of the interactions to ensure that the thermodynamic identity is satisfied. Note that in the `O2scl_eos` library, where in the equations of state the normal thermodynamic identity is used

$$P = -\varepsilon + sT + \mu n$$

Frequently, the interactions which create an effective chemical potential which is different than `part::mu` thus create extra terms in the pressure and the energy density for the given equation of state.

At zero temperature, fermions and bosons can be treated exactly in the classes `fermion` and `boson`. The `quark` class is a descendant of the `fermion` class which contains extra data members for the `quark` condensate and the contribution to the bag constant. The `classical` is a descendant of both `fermion` and `boson` and calculates everything in the `classical` limit.

At finite temperature, there are different classes corresponding to different approaches to computing the integrals over the distribution functions. The approximation scheme from [Johns96](#) is used in `eff_boson`, `eff_fermion`, and `eff_quark`. An exact method employing direct integration of the distribution functions is used in `rel_boson` and `rel_fermion`, but these are necessarily quite a bit slower.

The class `eff_fermion` usually works to within about 1 `part` in 10^4 , but can be as bad as 1 `part` in 10^2 in some more extreme cases. The default settings for `rel_fermion` give an accuracy of at least 1 `part` in 10^6 (and frequently better than this). For `rel_fermion`, some additional accuracy may be obtained by decreasing the integration tolerances.

The class `nonrel_fermion` assumes a non-relativistic dispersion relation for fermions. It includes zero-temperature methods and an exact method for finite temperatures. The non-relativistic integrands are much simpler and `nonrel_fermion` uses the appropriate GSL functions (which are nearly exact) to compute them.

Units:

Factors of \hbar , c and k_B have been removed everywhere, so that mass, energy, and temperature all have the same units. Number and entropy densities have units of mass cubed (or energy cubed). This particle classes can be used with any system of units which is based on powers of one unit, i.e. $[n] = [T]^3 = [m]^3 = [P]^{3/4} = [\varepsilon]^{3/4}$, etc.

Derivative information:

Sometimes it is useful to know derivatives like ds/dT in addition to the energy and pressure. There are three classes which compute these derivatives for fermions and `classical` particles. The class `sn_classical` handles the nondegenerate limit, `sn_fermion` handles fermions and `sn_nr_fermion` handles nonrelativistic fermions. These classes compute the derivatives

$$\left(\frac{dn}{d\mu}\right)_T, \quad \left(\frac{dn}{dT}\right)_\mu, \quad \text{and} \quad \left(\frac{ds}{dT}\right)_\mu$$

All other first derivatives of the thermodynamic functions can be written in terms of these three. To see how to compute the specific heat, for example, see the discussion in the documentation of `deriv_part`.

1.3 Atomic nuclei

Nuclei

Atomic nuclei, class `nucleus`, are implemented as descendants of `classical`. This class sets the value of `nucleus::inc_rest_mass` to false by default.

Nuclear mass formulas are given as children of `nuclear_mass`. The class `ame_mass` provides the experimental data from [Audi95](#) or [Audi03](#), the class `mnmsk_mass` provides the mass formula from [Moller95](#), and the class `hfb_mass` provides the mass formula from [Goriely02](#), [Samyn04](#), or [Goriely07](#). A simple semi-empirical mass formula is given in `semi_empirical_mass` and this can be fit to experimentally measured masses using `mass_fit`.

The class `nuclear_dist` provides an generic base class for a collection of several nuclei with an STL-like iterator. There are two implementations of this base class, `simple_dist` which provides a simple distribution and `full_dist` which enumerates all the nuclei for a given mass formula.

1.4 Example source code

1.4.1 exlist_subsect

- [Particle example](#)
- [Nuclear mass fit example](#)

1.4.2 Particle example

```
/* Example: ex_part.cpp
-----
*/

#include <cmath>
#include <o2scl/test_mgr.h>
#include <o2scl/constants.h>
#include <o2scl/eff_fermion.h>
#include <o2scl/rel_fermion.h>
#include <o2scl/classical.h>

using namespace std;
using namespace o2scl;
using namespace o2scl_const;

int main(void) {
    test_mgr t;
    t.set_output_level(1);

    // Create two different electrons, one using the exact method from
    // rel_fermion, and the other from the approximate scheme used in
    // eff_fermion. We work in units of inverse Fermis, so that energy
    // density is fm-4. We also use a classical particle, to compare
    // to the nondegenerate approximation.
    eff_fermion e(o2scl_fm::mass_electron,2.0);
    rel_fermion e2(o2scl_fm::mass_electron,2.0);
    classical e3(o2scl_fm::mass_electron,2.0);

    // Compute the pressure at a density of 0.0001 fm-3 and a
    // temperature of 10 MeV. At these temperatures, the electrons are
    // non-degenerate, and Boltzmann statistics nearly applies.
    e.n=0.0001;
    e.calc_density(10.0/hc_mev_fm);
    e2.n=0.0001;
    e2.calc_density(10.0/hc_mev_fm);
    e3.n=0.0001;
    e3.calc_density(10.0/hc_mev_fm);

    cout << e.pr << " " << e2.pr << " " << e3.pr << " "
         << e.n*10.0/hc_mev_fm << endl;

    // Test
    t.test_rel(e.pr,e2.pr,1.0e-2,"EFF vs. exact");
    t.test_rel(e2.pr,e3.pr,4.0e-1,"classical vs. exact");
    t.test_rel(e.n*10.0/hc_mev_fm,e3.pr,1.0e-1,"classical vs. ideal gas law");

    // Compute the pressure at a density of 0.1 fm-3 and a
    // temperature of 1 MeV. At these temperatures, the electrons are
    // strongly degenerate
    e.n=0.0001;
```

```

e.calc_density(10.0/hc_mev_fm);
e2.n=0.0001;
e2.calc_density(10.0/hc_mev_fm);
cout << e.pr << " " << e2.pr << endl;

// Test
t.test_rel(e.pr,e2.pr,1.0e-2,"EFF vs. exact");

// Now add the contribution to the pressure from positrons using the
// implementation of part::pair_density()
e.n=0.0001;
e.pair_density(10.0/hc_mev_fm);
e2.n=0.0001;
e2.pair_density(10.0/hc_mev_fm);
cout << e.pr << " " << e2.pr << endl;

// Test
t.test_rel(e.pr,e2.pr,1.0e-2,"EFF vs. exact");

t.report();
return 0;
}
// End of example

```

1.4.3 Nuclear mass fit example

```

/* Example: ex_mass_fit.cpp
-----
*/

#include <iostream>
#include <o2scl/test_mgr.h>
#include <o2scl/mass_fit.h>

using namespace std;
using namespace o2scl;
using namespace o2scl_const;
using namespace o2scl_fm;

int main(void) {
    test_mgr t;
    t.set_output_level(1);

    cout.setf(ios::scientific);

    // The RMS deviation of the fit
    double res;
    // The mass formula to be fitted
    semi_empirical_mass sem;
    // The fitting class
    mass_fit mf;

    // Perform the fit
    mf.fit(sem,res);

    // Output the results
    cout << sem.B << " " << sem.Sv << " " << sem.Ss << " "
        << sem.Ec << " " << sem.Epair << endl;
    cout << res << endl;
    t.test_gen(res<4.0,"Successful fit.");

    t.report();
    return 0;
}
// End of example

```

1.5 Bibliography

Some of the references which contain links should direct you to the work referred to directly through dx.doi.org.

Audi95: [G. Audi and A. H. Wapstra](#), Nucl. Phys. A **595** (1995) 409-480.

Audi03: [G. Audi, A. H. Wapstra and C. Thibault](#), Nucl. Phys. A **729** (2003) 337.

Callen H. B. Callen, "Thermodynamics and an Introduction to Thermostatistics.", 2nd edition.

Eggleton73: [P. P. Eggleton, J. Faulkner, and B. P. Flannery](#), Astron. and Astrophys. **23** (1973) 325.

Goriely02: [S. Goriely, M. Samyn, P.-H. Heenen, J. M. Pearson, and F. Tondeur](#), Phys. Rev. C **66** (2002) 024326.

Goriely07: [S. Goriely, M. Samyn, and J. M. Pearson](#), Phys. Rev. C **75** (2007) 064312.

Johns96: [Johns, P.J. Ellis, and J.M. Lattimer](#), Astrophys. J. **473** (1996) 1020.

Landau: L. D. Landau, "Statistical Physics: Part 1", 3rd Edition.

Moller95: [P. Moller, J.R. Nix, W.D. Myers, and W.J. Swiatecki](#), At. Data Nucl. Data Tables **59** (1995) 185.

Samyn04: [M. Samyn, S. Goriely, M. Bender and J. M. Pearson](#), Phys. Rev. C **70** (2004) 044309.

2 Ideas for future development

Class [ame_mass](#) Create a caching and more intelligent search system for the **table**. The **table** is sorted by A and then N, so we could probably just copy the search routine from [mnmsk_mass](#), which is sorted by Z and then N.

Class [ame_mass](#) Use the atomic mass unit and other constants defined in the evaluation

Class [classical](#) Write a `calc_density_zerot()` function for completeness?

Class [eff_fermion](#) Use bracketing to speed up one-dimensional **root** finding.

Class [fermion](#) Use `hypot()` and other more accurate functions for the analytic expressions for the zero temperature integrals. [Progress has been made, but there are probably other functions which may break down for small but finite masses and temperatures]

Class [fermion_T](#) Create a Chebyshev approximation for inverting the the Fermi functions for `massless_calc_density()` functions?

Global [fermion_T::massless_pair_density\(const double temper\)](#) This could be improved by including more terms in the expansions.

Class [mass_fit](#) Convert to a real fit with errors and covariance, etc.

Class [nonrel_fermion](#) This could be improved by performing a Chebyshev approximation (for example) to invert the density integral so that we don't need to use a solver.

Class [nuclear_mass](#) Make the treatment of the electron binding energy contribution more consistent.

Class `nuclear_mass` It might be useful to consider a fudge factor to ensure no problems with finite precision arithmetic when converting `double` to `int`.

Global `nuclear_mass_info::parse_elstring(std::string ela, int &Z, int &N, int &A)` Allow A to precede Z.

Global `nuclear_mass_info::parse_elstring(std::string ela, int &Z, int &N, int &A)` Right now, n4 is interpreted incorrectly as Nitrogen-4, rather than the tetraneutron.

Class `rel_fermion` Improve the asymptotics of the non-degenerate integrands.

Class `rel_fermion` It appears this doesn't compute the uncertainty in the chemical potential or density with `calc_density()`. This could be fixed.

Class `simple_dist` Make the vector constructor into a template so it accepts any type. Do the same for `set_dist()`.

Class `sn_fermion` It might be worth coding up direct differentiation, or differentiating the eff results, as these may succeed more generally.

Class `sn_fermion` This class will have difficulty with extremely degenerate or extremely non-degenerate systems. Fix this.

Class `sn_fermion` Create a more intelligent method for dealing with bad initial guesses for the chemical potential in `calc_density()`.

3 Todo List

Class `eff_boson` Better documentation (see `eff_fermion`)

Class `eff_boson` Remove the 'meth2' stuff

Class `eff_boson` Remove static variables `fix_density` and `stat_temper`

Class `eff_fermion` There's still `def_err_hnd.set_mode(0)` in the testing code, probably because the solver has a hard time for extreme values.

Global `eff_fermion::calc_mu(const double temper)` Should see if the function actually works if $(\mu - m)/T = -199$.

Class `eff_quark` Add testing.

Class `nonrel_fermion` Check behaviour of `calc_density()` at zero density, and compare with that from `eff_fermion`, `rel_fermion`, and `classical`.

Class `nonrel_fermion` Make sure to test with non-interacting equal to true or false, and document whether or not it works with both `inc_rest_mass` equal to true or false

Class `rel_boson` Testing not completely finished.

Class `sn_classical` This does not work with `inc_rest_mass=true`

Class `sn_fermion` This needs to be corrected to calculate $\sqrt{k^2 + m^{*2}} - m$ gracefully when $m^* \approx m \ll k$.

Class `sn_fermion` Call error handler if `inc_rest_mass` is true or update to properly treat the case when `inc_rest_mass` is true.

4 Data Structure Documentation

4.1 ame_entry Struct Reference

Atomic mass entry structure.

```
#include <nuclear_mass.h>
```

4.1.1 Detailed Description

Definition at line 554 of file `nuclear_mass.h`.

Data Fields

- int `NMZ`
N-Z.
- int `N`
Neutron number.
- int `Z`
Proton number.
- int `A`
Atomic number.
- std::string `el`
Element name.
- std::string `orig`
Data origin.
- double `mass`
Mass excess.
- double `dmass`
Mass excess uncertainty.
- double `be`
Binding energy (given in the '95 data).
- double `dbe`
Binding energy uncertainty (given in the '95 data).
- double `beoa`
Binding energy / A (given in the '03 data).
- double `dbeoa`
Binding energy / A uncertainty (given in the '03 data).
- std::string `bdmode`

Beta decay mode.

- double [bde](#)

Beta-decay energy.

- double [dbde](#)

Beta-decay energy uncertainty.

- int [A2](#)

?

- double [amass](#)

Atomic mass.

- double [damass](#)

Atomic mass uncertainty.

The documentation for this struct was generated from the following file:

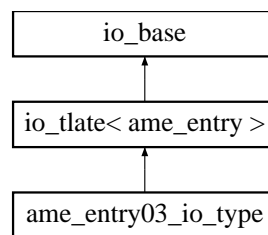
- nuclear_mass.h

4.2 ame_entry03_io_type Class Reference

A support class for I/O of the 2003 AME data.

```
#include <nuclear_mass.h>
```

Inheritance diagram for ame_entry03_io_type::



4.2.1 Detailed Description

Definition at line 623 of file nuclear_mass.h.

Public Member Functions

- int **input** (**cinput** *co, **in_file_format** *ins, [ame_entry](#) *t)
- int **output** (**coutput** *co, **out_file_format** *outs, [ame_entry](#) *t)
- virtual const char * **type** ()

The documentation for this class was generated from the following file:

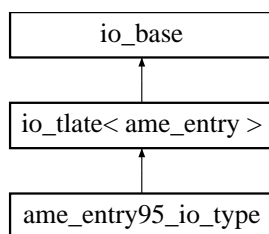
- nuclear_mass.h

4.3 ame_entry95_io_type Class Reference

A support class for I/O of the 1995 AME data.

```
#include <nuclear_mass.h>
```

Inheritance diagram for ame_entry95_io_type::



4.3.1 Detailed Description

Definition at line 614 of file nuclear_mass.h.

Public Member Functions

- `int input (cinput *co, in_file_format *ins, ame_entry *t)`
- `int output (coutput *co, out_file_format *outs, ame_entry *t)`
- `virtual const char * type ()`

The documentation for this class was generated from the following file:

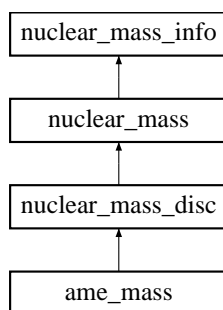
- nuclear_mass.h

4.4 ame_mass Class Reference

Mass formula from the Atomic Mass Evaluation (2005 and 1993).

```
#include <nuclear_mass.h>
```

Inheritance diagram for ame_mass::



4.4.1 Detailed Description

This class provides an interface to the atomic mass **table** using data from [Audi95](#) and [Audi03](#).

There are four data sets, selected by the specification of the `version` string in the constructor.

- "95rmd" - "Recommended" data from [Audi95](#) (ame95rmd.o2)
- "95exp" - "Experimental" data from [Audi95](#) (ame95exp.o2)
- "03round" - "Rounded" data from [Audi03](#) (ame03round.o2)
- "03" - Data from [Audi03](#) (default) (ame03.o2)

If any string other than these four is used, the default data is loaded. If the constructor cannot find the data file (e.g. because of a broken installation), then `ame::is_loaded()` returns false.

The 1995 data provided the binding energy stored in `ame_entry::be` and `ame_entry::dbe`, while the 2003 data provided the binding energy divided by the atomic number stored in `ame_entry::beoa` and `ame_entry::dbeoa`. When the 1995 data is used `ame_entry::beoa` and `ame_entry::dbeoa` are calculated automatically, and when the 2003 data is used `ame_entry::be` and `ame_entry::dbe` are calculated automatically.

Note that blank entries in the original **table** that correspond to columns represented by the type `double` are set to zero arbitrarily.

Note that all uncertainties are 1 sigma uncertainties.

Warning:

There are strict definitions of the atomic mass unit and other constants that are defined by the 1995 and 2003 atomic mass evaluations which are not used at present.

Idea for future

Create a caching and more intelligent search system for the **table**. The **table** is sorted by A and then N, so we could probably just copy the search routine from `mnmsk_mass`, which is sorted by Z and then N.

Idea for future

Use the atomic mass unit and other constants defined in the evaluation

Definition at line 673 of file `nuclear_mass.h`.

Public Member Functions

- `ame_mass` (`std::string version=""`)
*Create a **collection** specified by version.*
- virtual `const char * type ()`
Return the type, "ame_mass".
- virtual `bool is_included (int Z, int N)`
*Return false if the mass formula does not include specified **nucleus**.*
- virtual `double mass_excess (int Z, int N)`
Given Z and N, return the mass excess in MeV.
- `ame_entry get_ZN (int l_Z, int l_N)`
Get element with Z=l_Z and N=l_N (e.g. 82,126).
- `ame_entry get_ZA (int l_Z, int l_A)`
Get element with Z=l_Z and A=l_A (e.g. 82,208).
- `ame_entry get_elA (std::string l_el, int l_A)`
Get element with name l_el and A=l_A (e.g. "Pb",208).
- `ame_entry get (std::string nucleus)`
Get element with string (e.g. "Pb208").
- `bool is_loaded ()`
Returns true if the constructor succesfully loaded the data.

Data Fields

- `int n`
The number of entries (about 3000).
- `std::string * short_names`
The short names of the columns (length 16).
- `std::string * col_names`
The long names of the columns (length 16).
- `std::string reference`

The reference for the original data.

- `ame_entry * mass`

The array containing the mass data of length `ame::n`.

Protected Attributes

- bool `loaded`

True if loading the data was successful.

The documentation for this class was generated from the following file:

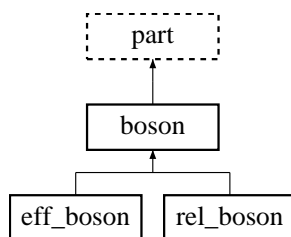
- `nuclear_mass.h`

4.5 boson Class Reference

Boson class [abstract base].

```
#include <boson.h>
```

Inheritance diagram for boson::



4.5.1 Detailed Description

For bosons:

- if either `nu` or `mu` is greater than `ms`, then they are taken to be equal to `ms`
- All contributions from any type of condensate are ignored.

This Mathematica notebook contains the series expansions for the bosonic integrals. functions.

```
doc/o2scl/extras/boson.nb
doc/o2scl/extras/boson.pdf
```

Definition at line 66 of file `boson.h`.

Public Member Functions

- `boson` (double `m`=0.0, double `g`=0.0)
Create a `boson` with mass `m` and degeneracy `g`.
- virtual int `calc_mu` (const double `temper`)=0
- virtual int `calc_density` (const double `temper`)=0
- virtual int `pair_mu` (const double `temper`)=0

- virtual int **pair_density** (const double temper)=0
- virtual int **massless_calc_mu** (const double temper)
Calculate properties of massless bosons.
- virtual const char * **type** ()
Return string denoting type ("boson").

Data Fields

- double **co**
The condensate.

4.5.2 Member Function Documentation

4.5.2.1 virtual int massless_calc_mu (const double temper) [virtual]

The expressions used are exact. The chemical potentials are ignored and the scalar density is set to zero

4.5.3 Field Documentation

4.5.3.1 double co

The condensate variable is mostly ignored by class **boson** and its descendants, and is provided for user storage.

Definition at line 76 of file boson.h.

The documentation for this class was generated from the following file:

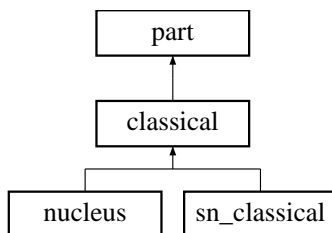
- boson.h

4.6 classical Class Reference

Classical particle class.

```
#include <classical.h>
```

Inheritance diagram for classical::



4.6.1 Detailed Description

Idea for future

Write a calc_density_zerot() function for completeness?

Definition at line 47 of file classical.h.

Public Member Functions

- `classical` (double `m`=0.0, double `g`=0.0)
Create a *classical* particle with mass `m` and degeneracy `g`.
- virtual int `calc_mu` (const double `temper`)
Calculate properties as function of chemical potential.
- virtual int `calc_density` (const double `temper`)
Calculate properties as function of density.
- virtual const char * `type` ()
Return string denoting type ("classical").

The documentation for this class was generated from the following file:

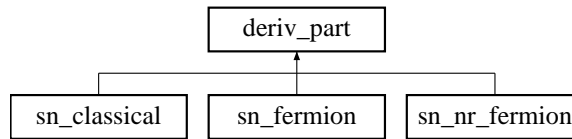
- `classical.h`

4.7 deriv_part Class Reference

Storage for derivatives wrt μ and T .

```
#include <deriv_part.h>
```

Inheritance diagram for `deriv_part`:



4.7.1 Detailed Description

The variables `dndmu`, `dndT`, and `dsdT` correspond to

$$\left(\frac{dn}{d\mu}\right)_{T,V}, \quad \left(\frac{dn}{dT}\right)_{\mu,V}, \quad \text{and} \quad \left(\frac{ds}{dT}\right)_{\mu,V}$$

respectively. All other derivatives can be expressed simply in terms of these three. Note that volume derivatives are trivial, since both the entropy and number scale linearly with the volume.

Derivatives wrt to chemical potential and temperature:

There is a Maxwell relation

$$\left(\frac{ds}{d\mu}\right)_{T,V} = \left(\frac{dn}{dT}\right)_{\mu,V}$$

The pressure derivatives are trivial

$$\left(\frac{dP}{d\mu}\right)_{T,V} = n, \quad \left(\frac{dP}{dT}\right)_{\mu,V} = s$$

The energy density derivatives are related through the thermodynamic identity:

$$\begin{aligned} \left(\frac{d\varepsilon}{d\mu}\right)_{T,V} &= \mu \left(\frac{dn}{d\mu}\right)_{T,V} + T \left(\frac{ds}{d\mu}\right)_{T,V} \\ \left(\frac{d\varepsilon}{dT}\right)_{\mu,V} &= \mu \left(\frac{dn}{dT}\right)_{\mu,V} + T \left(\frac{ds}{dT}\right)_{\mu,V} \end{aligned}$$

Other derivatives:

Note that the derivative of the entropy with respect to the temperature above is not the specific heat per particle, c_V . The specific heat per particle is

$$c_V = \frac{T}{N} \left(\frac{\partial S}{\partial T} \right)_{V,N} = \frac{T}{n} \left(\frac{\partial s}{\partial T} \right)_{V,n}$$

As noted in [Particles](#) in the User's Guide for `O2scl_part`, we work in units so that $\hbar = c = k_B = 1$. In this case, c_V is unitless as defined here. To compute c_V in terms of the derivatives above, note that the descendants of `deriv_part` provide all of the thermodynamic functions in terms of μ, V and T , so we have

$$s = s(\mu, T, V) \quad \text{and} \quad n = n(\mu, T, V).$$

We can then construct a function

$$s = s[\mu(n, T, V), T, V]$$

and then write the required derivative directly

$$\left(\frac{\partial s}{\partial T} \right)_{n,V} = \left(\frac{\partial s}{\partial \mu} \right)_{T,V} \left(\frac{\partial \mu}{\partial T} \right)_{n,V} + \left(\frac{\partial s}{\partial T} \right)_{\mu,V}.$$

Now we use the identity

$$\left(\frac{\partial \mu}{\partial T} \right)_{n,V} = - \left(\frac{\partial n}{\partial T} \right)_{\mu,V} \left(\frac{\partial n}{\partial \mu} \right)_{T,V}^{-1},$$

and the Maxwell relation above to give

$$C_V = \frac{T}{n} \left[\left(\frac{\partial s}{\partial T} \right)_{\mu,V} - \left(\frac{\partial n}{\partial T} \right)_{\mu,V}^2 \left(\frac{\partial n}{\partial \mu} \right)_{T,V}^{-1} \right]$$

which expresses the specific heat in terms of the three derivatives which are given.

For, c_P , defined as

$$c_P = \frac{T}{N} \left(\frac{\partial S}{\partial T} \right)_{N,P}$$

(which is also unitless) we can write functions

$$S = S(N, T, V) \quad \text{and} \quad V = V(N, P, T)$$

which imply

$$\left(\frac{\partial S}{\partial T} \right)_{N,P} = \left(\frac{\partial S}{\partial T} \right)_{N,V} + \left(\frac{\partial S}{\partial V} \right)_{N,T} \left(\frac{\partial V}{\partial T} \right)_{N,P}.$$

Thus we require the derivatives

$$\left(\frac{\partial S}{\partial T} \right)_{N,V}, \left(\frac{\partial S}{\partial V} \right)_{N,T}, \quad \text{and} \quad \left(\frac{\partial V}{\partial T} \right)_{N,P}.$$

To compute the new entropy derivatives, we can write

$$S = S(\mu(N, T, V), T, V)$$

to get

$$\left(\frac{\partial S}{\partial T} \right)_{N,V} = \left(\frac{\partial S}{\partial \mu} \right)_{T,V} \left(\frac{\partial \mu}{\partial T} \right)_{N,V} + \left(\frac{\partial S}{\partial T} \right)_{\mu,V},$$

and

$$\left(\frac{\partial S}{\partial V} \right)_{N,T} = \left(\frac{\partial S}{\partial \mu} \right)_{T,V} \left(\frac{\partial \mu}{\partial V} \right)_{N,T} + \left(\frac{\partial S}{\partial V} \right)_{\mu,T}.$$

These require the chemical potential derivatives which have associated Maxwell relations

$$\left(\frac{\partial \mu}{\partial T}\right)_{N,V} = -\left(\frac{\partial S}{\partial N}\right)_{T,V} \quad \text{and} \quad \left(\frac{\partial \mu}{\partial V}\right)_{N,T} = -\left(\frac{\partial P}{\partial N}\right)_{T,V}.$$

Finally, we can rewrite the derivatives on the right hand sides in terms of derivatives of functions of μ , V and T ,

$$\left(\frac{\partial S}{\partial N}\right)_{T,V} = \left(\frac{\partial S}{\partial \mu}\right)_{T,V} \left(\frac{\partial N}{\partial \mu}\right)_{T,V}^{-1},$$

and

$$\left(\frac{\partial P}{\partial N}\right)_{T,V} = \left(\frac{\partial P}{\partial \mu}\right)_{T,V} \left(\frac{\partial N}{\partial \mu}\right)_{T,V}^{-1}.$$

The volume derivative,

$$\left(\frac{\partial V}{\partial T}\right)_{N,P},$$

is related to the coefficient of thermal expansion, sometimes called α ,

$$\alpha \equiv \frac{1}{V} \left(\frac{\partial V}{\partial T}\right)_{N,P}.$$

We can rewrite the derivative

$$\left(\frac{\partial V}{\partial T}\right)_{N,P} = -\left(\frac{\partial P}{\partial T}\right)_{N,V} \left(\frac{\partial P}{\partial V}\right)_{N,T}^{-1}.$$

The first term can be computed from the Maxwell relation

$$\left(\frac{\partial P}{\partial T}\right)_{N,V} = \left(\frac{\partial S}{\partial V}\right)_{N,T},$$

where the entropy derivative was computed above. The second term (related to the inverse of the isothermal compressibility, $\kappa_T \equiv (-1/V)(\partial V/\partial P)_{T,N}$ can be computed from the function $P = P(\mu(N, V, T), V, T)$

$$\left(\frac{\partial P}{\partial V}\right)_{N,T} = \left(\frac{\partial P}{\partial \mu}\right)_{T,V} \left(\frac{\partial \mu}{\partial V}\right)_{N,T} + \left(\frac{\partial P}{\partial V}\right)_{\mu,T}$$

where the chemical potential derivative was computed above.

The results above can be collected to give

$$\left(\frac{\partial S}{\partial T}\right)_{N,P} = \left(\frac{\partial S}{\partial T}\right)_{\mu,V} + \frac{S^2}{N^2} \left(\frac{\partial N}{\partial \mu}\right)_{T,V} - \frac{2S}{N} \left(\frac{\partial N}{\partial T}\right)_{\mu,V},$$

which implies

$$c_P = \frac{T}{n} \left(\frac{\partial s}{\partial T}\right)_{\mu,V} + \frac{s^2 T}{n^3} \left(\frac{\partial n}{\partial \mu}\right)_{T,V} - \frac{2sT}{n^2} \left(\frac{\partial n}{\partial T}\right)_{\mu,V},$$

This derivation also gives the well-known relationship between the specific heats at constant volume and constant pressure,

$$c_P = c_V + \frac{T\alpha^2}{n\kappa_T}.$$

No derivative with respect to the bare mass is given, since classes cannot know how to relate the effective mass to the bare mass.

Definition at line 260 of file deriv_part.h.

Data Fields

- double [dndmu](#)
Derivative of number density with respect to chemical potential.
- double [dndT](#)
Derivative of number density with respect to temperature.
- double [dsdT](#)
Derivative of entropy density with respect to temperature.
- double [dndm](#)
Derivative of number density with respect to the effective mass.

The documentation for this class was generated from the following file:

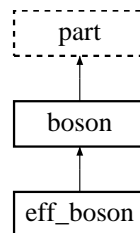
- `deriv_part.h`

4.8 eff_boson Class Reference

Boson class from fitting method.

```
#include <eff_boson.h>
```

Inheritance diagram for eff_boson::



4.8.1 Detailed Description

Todo

Better documentation (see [eff_fermion](#))

Todo

Remove the 'meth2' stuff

Todo

Remove static variables `fix_density` and `stat_temper`

Definition at line 51 of file `eff_boson.h`.

Public Member Functions

- [eff_boson](#) (double [m](#)=0.0, double [g](#)=0.0)
Create a [boson](#) with mass [m](#) and degeneracy [g](#).
- virtual int [calc_mu](#) (const double [temper](#))
Calculate thermodynamic properties as function of chemical potential.
- virtual int [calc_density](#) (const double [temper](#))
Calculate thermodynamic properties as function of density.

- virtual int [pair_mu](#) (const double temper)
Calculate thermodynamic properties with antiparticles as function of chemical potential.
- virtual int [pair_density](#) (const double temper)
Calculate thermodynamic properties with antiparticles as function of density.
- int [set_psi_root](#) ([root](#)< double, [funct](#)< double > > &rp)
Set the solver for use in calculating ψ .
- int [set_density_mroot](#) ([mroot](#)< int, [mm_funct](#)< int > > &rp)
Set the solver for use in calculating the chemical potential from the density.
- int [set_meth2_root](#) ([root](#)< int, [funct](#)< int > > &rp)
Set the solver for use in calculating the chemical potential from the density (meth2=true).
- virtual const char * [type](#) ()
Return string denoting type ("boson").

Static Public Member Functions

- static int [load_coefficients](#) (int ctype)
Load coefficients for finite-temperature approximation.

Data Fields

- [gsl_mroot_hybrids](#)< int, [mm_funct](#)< int > > [def_density_mroot](#)
The default solver for [calc_density\(\)](#) and [pair_density\(\)](#).
- [cern_mroot_root](#)< double, [funct](#)< double > > [def_psi_root](#)
The default solver for ψ .
- [cern_mroot_root](#)< int, [funct](#)< int > > [def_meth2_root](#)
The default solver for [calc_density\(\)](#) and [pair_density\(\)](#).

Static Public Attributes

- static const int [cf_boselat3](#) = 1
A set of coefficients from Jim Lattimer.
- static const int [cf_bosejel21](#) = 2
A set of coefficients from [Johns96](#).
- static const int [cf_bosejel22](#) = 3
A set of coefficients from [Johns96](#).
- static const int [cf_bosejel34](#) = 4
A set of coefficients from [Johns96](#).
- static const int [cf_bosejel34cons](#) = 5
The set of coefficients from [Johns96](#) which retains better thermodynamic consistency.

Protected Member Functions

- int [solve_fun](#) (double x, double &y, double &psi)
The function which solves for h from ψ .
- int [density_fun](#) (size_t nv, const [ovector_base](#) &x, [ovector_base](#) &y, int &pa)
Fix density for [calc_density\(\)](#).
- int [pair_density_fun](#) (size_t nv, const [ovector_base](#) &x, [ovector_base](#) &y, int &pa)
Fix density for [pair_density\(\)](#).

Protected Attributes

- [mroot](#)< int, [mm_funct](#)< int > > * [density_mroot](#)
The solver for [calc_density\(\)](#).
- [root](#)< double, [funct](#)< double > > * [psi_root](#)

The solver to compute h from ψ .

- **root**< int, **funct**< int > > * [meth2_root](#)
The solver for [calc_density\(\)](#).

Static Protected Attributes

- static double ** [Pmnb](#)
The coefficients.
- static int [sizem](#)
The number of coefficient rows.
- static int [sizen](#)
The number of coefficient columns.
- static double [parma](#)
The parameter, a .
- static double [fix_density](#)
Temporary storage.
- static double [stat_temper](#)
Temporary storage.

4.8.2 Member Function Documentation

4.8.2.1 static int load_coefficients (int ctype) [static]

Presently acceptable values of `fn` are: `boselat3` from Lattimer's notes `bosejel21`, `bosejel22`, `bosejel34`, and `bosejel34cons` from [Johns96](#).

The documentation for this class was generated from the following file:

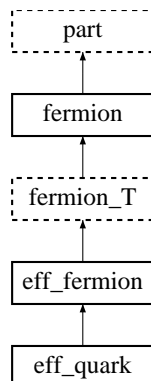
- `eff_boson.h`

4.9 eff_fermion Class Reference

Fermion class from fitting method.

```
#include <eff_fermion.h>
```

Inheritance diagram for `eff_fermion`:



4.9.1 Detailed Description

Based on the fitting method of [Johns96](#) which is an update of the method from [Eggleton73](#). This method is approximate, but very fast. For a more accurate (but slower) method, use [rel_fermion](#).

Given the chemical potential and the temperature the functions `calc_mu()` and `pair_mu()` work by solving the equation (c.f. Eq. 15 in [Johns96](#))

$$\psi = 2\sqrt{1 + f/a} + \log \left(\frac{\sqrt{1 + f/a} - 1}{\sqrt{1 + f/a} + 1} \right)$$

for f given $\psi = (\mu - m)/T$. If $f/a < 10^{-10}$, then the alternative expression

$$\psi = 2(1 + f/(2a)) + \log \left[\frac{f/(2a)}{(1 + f/(2a))} \right]$$

is used. The pressure, energy density, and entropy, are determined as polynomials in f with a set of precomputed coefficients as done in [Johns96](#).

If ψ is too small (less than about -200), the above procedure fails. To handle this, this class uses the classical result if $\psi < \text{min_psi}$, where `min_psi` defaults to -200.

When the density and temperature is given instead (`calc_density()` and `pair_density()`), then there are two ways to proceed.

- Use the density to solve for f .
- Use the density to solve for the chemical potential.

Because the density is a complicated polynomial in f , the former procedure does not work very well even though it might be less time consuming. In this class, the density is solved for the effective chemical potential instead. The initial guess is just taken from the present value of `part::nu`.

Note:

It is important to note that the coefficients are static and apply to all objects of type `eff_fermion`.

The coefficients are stored in static data, and there is a small possibility of a memory leak in a multithreaded program if more than one instance of `eff_fermion` tries to initialize these coefficients in the constructor. This can easily be avoided by ensuring that the static coefficients are initialized by a single thread beforehand.

Todo

There's still `def_err_hnd.set_mode(0)` in the testing code, probably because the solver has a hard time for extreme values.

Idea for future

Use bracketing to speed up one-dimensional **root** finding.

Definition at line 101 of file `eff_fermion.h`.

Coefficients for finite-temperature approximation

- static const int `cf_fermilat3` = 1
A set of coefficients from Jim Lattimer.
- static const int `cf_fermijel2` = 2
The smaller set of coefficients from [Johns96](#).
- static const int `cf_fermijel3` = 3
The larger set of coefficients from [Johns96](#).
- static const int `cf_fermijel3cons` = 4
The set of coefficients from [Johns96](#) which retains better thermodynamic consistency.
- static int `load_coefficients` (int ctype)
Load coefficients.

Public Member Functions

- `eff_fermion` (double mass=0.0, double dof=0.0)
Create a `fermion` with mass `mass` and degeneracy `dof`.
- virtual int `calc_mu` (const double temper)
Calculate thermodynamic properties as function of chemical potential.
- virtual int `calc_density` (const double temper)
Calculate thermodynamic properties as function of density.
- virtual int `pair_mu` (const double temper)
Calculate thermodynamic properties with antiparticles as function of chemical potential.
- virtual int `pair_density` (const double temper)
Calculate thermodynamic properties with antiparticles as function of density.
- int `set_psi_root` (**root**< double, **funct**< double > > &rp)
Set the solver for use in calculating ψ .
- int `set_density_root` (**root**< double, **funct**< double > > &rp)
Set the solver for use in calculating the chemical potential from the density.
- virtual const char * `type` ()
Return string denoting type ("eff_fermion").

Data Fields

- double `tlimit`
If the temperature is less than `tlimit` then the zero-temperature functions are used (default 0).
- **cern_mroot_root**< double, **funct**< double > > `def_psi_root`
The default solver for ψ .
- **cern_mroot_root**< double, **funct**< double > > `def_density_root`
The default solver for `calc_density()` and `pair_density()`.
- double `min_psi`
The minimum value of ψ (default -200).

Protected Member Functions

- int `solve_fun` (double x, double &y, double &psi)
The function which solves for f from ψ .
- int `density_fun` (double x, double &y, double &temper)
Fix density for `calc_density()`.
- int `pair_density_fun` (double x, double &y, double &temper)
Fix density for `pair_density()`.

Protected Attributes

- **root**< double, **funct**< double > > * `psi_root`
The solver for ψ .
- **root**< double, **funct**< double > > * `density_root`
The other solver for `calc_density()`.

Static Protected Attributes

- static double ** `Pmnf`
The matrix of coefficients.
- static double `parma`
The parameter a .
- static int `sizem`
The array row size.
- static int `sizen`
The array column size.

4.9.2 Member Function Documentation

4.9.2.1 virtual int calc_density (const double *temper*) [virtual]

Warning:

This function needs a guess for the chemical potential, and will fail if that guess is not sufficiently accurate.

Implements [fermion_T](#).

Reimplemented in [eff_quark](#).

4.9.2.2 virtual int calc_mu (const double *temper*) [virtual]

If the quantity $(\mu - m)/T$ (or $(\nu - m^*)/T$ in the case of interacting particles) is less than -200, then this quietly sets the density, the scalar density, the energy density, the pressure and the entropy to zero and exits.

Todo

Should see if the function actually works if $(\mu - m)/T = -199$.

Implements [fermion_T](#).

Reimplemented in [eff_quark](#).

4.9.2.3 static int load_coefficients (int *ctype*) [static]

The argument `ctype` Should be one of the constants below.

4.9.2.4 virtual int pair_mu (const double *temper*) [virtual]

Warning:

This function needs a guess for the chemical potential, and will fail if that guess is not sufficiently accurate.

Implements [fermion_T](#).

Reimplemented in [eff_quark](#).

The documentation for this class was generated from the following file:

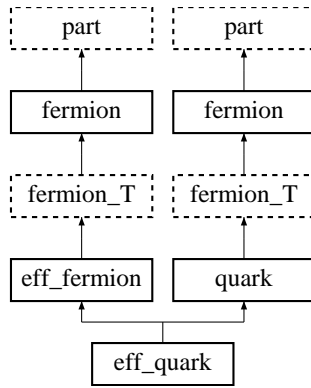
- `eff_fermion.h`

4.10 eff_quark Class Reference

Quark class from fitting method.

```
#include <eff_quark.h>
```

Inheritance diagram for `eff_quark`:



4.10.1 Detailed Description

Todo

Add testing.

Definition at line 45 of file eff_quark.h.

Public Member Functions

- `eff_quark` (double `m`=0.0, double `g`=0.0)
Create a *quark* with mass `m` and degeneracy `g`.
- virtual int `calc_mu` (const double `temper`)
Calculate thermodynamic properties as function of chemical potential.
- virtual int `calc_density` (const double `temper`)
Calculate thermodynamic properties as function of density.
- virtual int `pair_mu` (const double `temper`)
Calculate thermodynamic properties with antiparticles as function of chemical potential.
- virtual int `pair_density` (const double `temper`)
Calculate thermodynamic properties with antiparticles as function of density.
- virtual const char * `type` ()
Return string denoting type ("eff_quark").

4.10.2 Member Function Documentation

4.10.2.1 virtual int calc_density (const double *temper*) [virtual]

Warning:

This function needs a guess for the chemical potential, and will fail if that guess is not sufficiently accurate.

Reimplemented from [eff_fermion](#).

4.10.2.2 virtual int calc_mu (const double *temper*) [virtual]

If the quantity $(\mu - m)/T$ (or $(\nu - m^*)/T$ in the case of interacting particles) is less than -200, then this quietly sets the density, the scalar density, the energy density, the pressure and the entropy to zero and exits.

Todo

Should see if the function actually works if $(\mu - m)/T = -199$.

Reimplemented from [eff_fermion](#).

4.10.2.3 virtual int pair_mu (const double *temper*) [virtual]

Warning:

This function needs a guess for the chemical potential, and will fail if that guess is not sufficiently accurate.

Reimplemented from [eff_fermion](#).

The documentation for this class was generated from the following file:

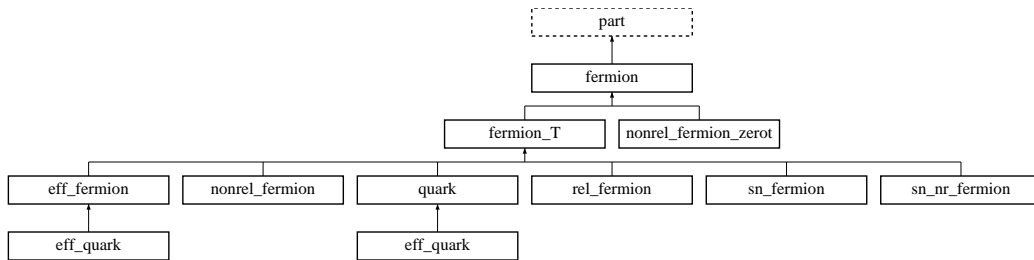
- `eff_quark.h`

4.11 fermion Class Reference

Fermion class.

```
#include <fermion.h>
```

Inheritance diagram for fermion::



4.11.1 Detailed Description

This is a base class for the computation of fermionic statistics at zero temperature. The more general case of finite temperature is taken care of by `fermion_T` objects. The primary functions are `calc_mu_zerot()` and `calc_density_zerot()` which compute all the thermodynamic quantities as a function of the chemical potential, or the density, respectively.

This class also adds two member data variables, `kf` and `del`, for the Fermi momentum and the gap.

Idea for future

Use `hypot()` and other more accurate functions for the analytic expressions for the zero temperature integrals. [Progress has been made, but there are probably other functions which may break down for small but finite masses and temperatures]

Definition at line 63 of file `fermion.h`.

Public Member Functions

- `fermion` (double mass=0, double dof=0)
Create a `fermion` with mass `mass` and degeneracy `dof`.
- `virtual const char * type ()`
Return string denoting type ("`fermion`").

Zero-temperature fermions

- `int kf_from_density ()`
Calculate the Fermi momentum from the density.

- int [energy_density_zerot \(\)](#)
Energy density at T=0 from [kf](#) and [ms](#).
- int [pressure_zerot \(\)](#)
Pressure at T=0 from [kf](#) and [ms](#).
- virtual int [calc_mu_zerot \(\)](#)
Zero temperature fermions from [nu](#) and [ms](#).
- virtual int [calc_density_zerot \(\)](#)
Zero temperature fermions from [n](#) and [ms](#).

Data Fields

- double [kf](#)
Fermi momentum.
- double [del](#)
Gap.

4.11.2 Member Function Documentation

4.11.2.1 virtual int calc_density_zerot () [virtual]

This function always returns `gsl_success`.

Reimplemented in [nonrel_fermion](#), and [nonrel_fermion_zerot](#).

4.11.2.2 virtual int calc_mu_zerot () [virtual]

This function always returns `gsl_success`.

Reimplemented in [nonrel_fermion](#), and [nonrel_fermion_zerot](#).

4.11.2.3 int energy_density_zerot ()

Calculates the integral

$$\varepsilon = \frac{g}{2\pi^2} \int_0^{k_F} k^2 \sqrt{k^2 + m^{*2}} dk$$

4.11.2.4 int kf_from_density ()

Uses the relation $k_F = (6\pi^2 n/g)^{1/3}$

4.11.2.5 int pressure_zerot ()

Calculates the integral

$$P = \frac{g}{6\pi^2} \int_0^{k_F} \frac{k^4}{\sqrt{k^2 + m^{*2}}} dk$$

The documentation for this class was generated from the following file:

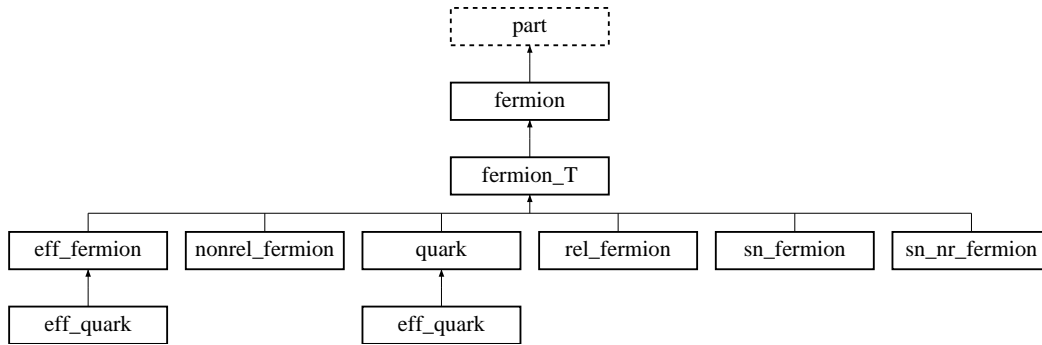
- [fermion.h](#)

4.12 fermion_T Class Reference

Fermion with finite-temperature thermodynamics [abstract base].

```
#include <fermion.h>
```

Inheritance diagram for fermion_T:



4.12.1 Detailed Description

This is an abstract base for the computation of finite-temperature fermionic statistics. Different children (e.g. [eff_fermion](#) and [rel_fermion](#)) use different techniques to computing the momentum integrations.

Because massless fermions at finite temperature are much simpler, there are separate member functions included in this class to handle them. The functions [massless_calc_density\(\)](#) and [massless_calc_mu\(\)](#) compute the thermodynamics of massless fermions at finite temperature given the density or the chemical potentials. The functions [massless_pair_density\(\)](#) and [massless_pair_mu\(\)](#) perform the same task, but automatically include antiparticles.

The function [massless_calc_density\(\)](#) uses a **root** object to solve for the chemical potential as a function of the density. The default is an object of type **cern_mroot_root**. The function [massless_pair_density\(\)](#) does not need to use the **root** object because of the simplification afforded by the inclusion of antiparticles.

Idea for future

Create a Chebyshev approximation for inverting the the Fermi functions for [massless_calc_density\(\)](#) functions?

This Mathematica notebook contains the derivations of related series expansions and some algebra for the [massless_pair\(\)](#) functions.

```
doc/o2scl/extras/fermion.nb
doc/o2scl/extras/fermion.pdf
```

Definition at line 174 of file fermion.h.

Public Member Functions

- [fermion_T](#) (double mass=0, double dof=0)
Create a [fermion](#) with mass `mass` and degeneracy `dof`.
- virtual int [calc_mu](#) (const double temper)=0
Calculate properties as function of chemical potential.
- virtual int [calc_density](#) (const double temper)=0
Calculate properties as function of density.
- virtual int [pair_mu](#) (const double temper)=0
Calculate properties with antiparticles as function of chemical potential.
- virtual int [pair_density](#) (const double temper)=0
Calculate properties with antiparticles as function of density.
- int [set_massless_root](#) (**root**< const double, **funct**< const double > > &rp)
Set the solver for use in [massless_calc_density\(\)](#).
- virtual const char * [type](#) ()

Return string denoting type ("fermion_T").

Massless fermions

- virtual int `massless_calc_mu` (const double temper)
Finite temperature massless fermions.
- virtual int `massless_calc_density` (const double temper)
Finite temperature massless fermions.
- int `massless_pair_mu` (const double temper)
Finite temperature massless fermions and antifermions.
- int `massless_pair_density` (const double temper)
Finite temperature massless fermions and antifermions.

Data Fields

- `cern_mroot_root` < const double, `funct` < const double > > `def_massless_root`
The default solver for `massless_calc_density()`.

4.12.2 Member Function Documentation

4.12.2.1 int massless_pair_density (const double temper)

In the cases $n^3 \gg T$ and $T \gg n^3$, expansions are used instead of the exact formulas to avoid loss of precision.

In particular, using the parameter

$$\alpha = \frac{g^2 \pi^2 T^6}{243 n^2}$$

and defining the expression

$$\text{cbt} = \alpha^{-1/6} (-1 + \sqrt{1 + \alpha})^{1/3}$$

we can write the chemical potential as

$$\mu = \frac{\pi T}{\sqrt{3}} \left(\frac{1}{\text{cbt}} - \text{cbt} \right)$$

These expressions, however, do not work well when α is very large or very small, so series expansions are used whenever $\alpha > 10^4$ or $\alpha < 3 \times 10^{-4}$. For small α ,

$$\left(\frac{1}{\text{cbt}} - \text{cbt} \right) \approx \frac{2^{1/3}}{\alpha^{1/6}} - \frac{\alpha^{1/6}}{2^{1/3}} + \frac{\alpha^{5/6}}{6 \cdot 2^{2/3}} + \frac{\alpha^{7/6}}{12 \cdot 2^{1/3}} - \frac{\alpha^{11/6}}{18 \cdot 2^{2/3}} - \frac{5\alpha^{13/6}}{144 \cdot 2^{1/3}} + \frac{77\alpha^{17/6}}{2592 \cdot 2^{2/3}}$$

and for large α ,

$$\left(\frac{1}{\text{cbt}} - \text{cbt} \right) \approx \frac{2}{3} \sqrt{\frac{1}{\alpha}} - \frac{8}{81} \left(\frac{1}{\alpha} \right)^{3/2} + \frac{32}{729} \left(\frac{1}{\alpha} \right)^{5/2}$$

This approach works to within about 1 part in 10^{12} , and is tested in `fermion_ts.cpp`.

Idea for future

This could be improved by including more terms in the expansions.

4.12.3 Field Documentation

4.12.3.1 cern_mroot_root<const double, funct<const double> > def_massless_root

We default to `cern_mroot_root` here since we don't have a bracket or a derivative.

Definition at line 304 of file `fermion.h`.

The documentation for this class was generated from the following file:

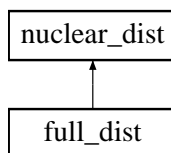
- `fermion.h`

4.13 full_dist Class Reference

Full distribution including all nuclei from a discrete mass formula.

```
#include <nuclear_dist.h>
```

Inheritance diagram for full_dist::



4.13.1 Detailed Description

For example, to create a collection of all nuclei from the most recent (2003) Atomic Mass Evaluation, and then output all the nuclei in the collection

```

ame_mass ame;
full_dist fd(&ame);
for(nuclear_dist::iterator ndi=fd.begin(); ndi!=fd.end(); ndi++) {
    cout << ndi->Z << " " << ndi->A << " " << ndi->m << endl;
}

```

Definition at line 247 of file nuclear_dist.h.

Public Member Functions

- [full_dist](#) ([nuclear_mass](#) &nm, int maxA=400, bool include_neutron=false)
Create a distribution including all nuclei with atomic numbers less than maxA from the mass formula nm.
- int [set_dist](#) ([nuclear_mass](#) &nm, int maxA=400, bool include_neutron=false)
Set the distribution to all nuclei with atomic numbers less than maxA from the mass formula nm.
- virtual [iterator begin](#) ()
The beginning of the distribution.
- virtual [iterator end](#) ()
The end of the distribution.
- virtual size_t [size](#) ()
The number of nuclei in the distribution.

4.13.2 Member Function Documentation

4.13.2.1 int set_dist (nuclear_mass & nm, int maxA = 400, bool include_neutron = false)

The information for the previous distribution is cleared before a new distribution is set.

The documentation for this class was generated from the following file:

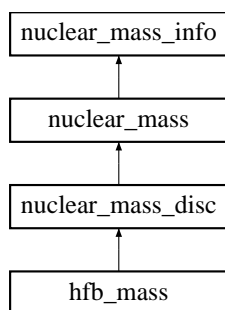
- nuclear_dist.h

4.14 hfb_mass Class Reference

HFB Mass formula.

```
#include <nuclear_mass.h>
```

Inheritance diagram for hfb_mass::



4.14.1 Detailed Description

Definition at line 1019 of file nuclear_mass.h.

Public Member Functions

- [hfb_mass](#) (size_t model=14)
Create a new mass formula object using the specified model number.
- virtual bool [is_included](#) (int Z, int N)
Return false if the mass formula does not include specified [nucleus](#).
- virtual double [mass_excess](#) (int Z, int N)
Given Z and N, return the mass excess in MeV.
- [hfb_mass_entry](#) [get_ZN](#) (int l_Z, int l_N)
Get the entry for the specified proton and neutron number.
- bool [is_loaded](#) ()
*Verify that the constructor properly loaded the **table**.*
- double [blank](#) ()
The value which corresponds to a blank entry.
- virtual const char * [type](#) ()
Return the type, "hfb_mass".

Data Fields

- int [n](#)
*The number of **table** entries.*
- [hfb_mass_entry](#) * [mass](#)
*The array containing the **table**.*

Protected Attributes

- bool [loaded](#)
*True if the **table** was successfully loaded.*
- int [last](#)
*The last **table** index for caching.*

4.14.2 Constructor & Destructor Documentation

4.14.2.1 hfb_mass (size_t model = 14)

Valid values of `model` at present are 2, 8, and 14, corresponding to the HFB2 ([Goriely02](#)), HFB8 ([Samyn04](#)), and HFB14 ([Goriely07](#)). If a number other than these three is given, the error handler is called.

4.14.3 Member Function Documentation

4.14.3.1 hfb_mass_entry get_ZN (int *l_Z*, int *l_N*)

This method searches the **table** using a cached binary search algorithm. It is assumed that the **table** is sorted first by proton number and then by neutron number.

The documentation for this class was generated from the following file:

- nuclear_mass.h

4.15 hfb_mass_entry Struct Reference

Mass formula entry structure for HFB mass formula.

```
#include <nuclear_mass.h>
```

4.15.1 Detailed Description

Definition at line 976 of file nuclear_mass.h.

Data Fields

- int [N](#)
Neutron number.
- int [Z](#)
Proton number.
- int [A](#)
Atomic number.
- double [bet2](#)
Beta 2 deformation.
- double [bet4](#)
Beta 4 deformation.
- double [Rch](#)
RMS charge radius.
- double [def_wig](#)
Deformation and Wigner energies.
- double [Sn](#)
Neutron separation energy.
- double [Sp](#)
Proton separation energy.
- double [Qbet](#)
Beta-decay energy.
- double [Mcal](#)
Calculated mass excess.
- double [Err](#)
Error between experimental and calculated mass excess.

The documentation for this struct was generated from the following file:

- nuclear_mass.h

4.16 mass_fit Class Reference

Fit a nuclear mass formula.

```
#include <mass_fit.h>
```

4.16.1 Detailed Description

There is an example of the usage of this class given in [Nuclear mass fit example](#).

Idea for future

Convert to a real fit with errors and covariance, etc.

Definition at line 46 of file mass_fit.h.

Public Member Functions

- virtual int [fit](#) ([nuclear_mass_fit](#) &n, double &res)
Fit the nuclear mass formula.
- virtual int [eval](#) ([nuclear_mass](#) &n, double &res)
Evaluate quality without fitting.
- int [set_mmin](#) ([multi_min](#)< int, [multi_funct](#)< int > > &umm)
Change the minimizer for use in the fit.
- int [set_dist](#) ([nuclear_dist](#) &uexp)
Set the distribution of nuclei to fit.

Data Fields

- bool [even_even](#)
If true, then only fit doubly-even nuclei (default false).
- int [minZ](#)
Minimum proton number to fit (default 8).
- int [minN](#)
Minimum neutron number to fit (default 8).
- [gsl_mmin_simp](#)< int, [multi_funct](#)< int > > [def_mmin](#)
The default minimizer.
- [full_dist](#) [def_dist](#)
The default distribution of nuclei to fit (defaults to all nuclei in [def_exp_mass](#)).
- [ame_mass](#) [def_exp_mass](#)
The default experimental nuclear mass object for [def_dist](#).

4.16.2 Field Documentation

4.16.2.1 [gsl_mmin_simp](#)<int,multi_funct<int> > [def_mmin](#)

The value of [def_mmin::ntrial](#) is automatically multiplied by 10 in the constructor because the minimization frequently requires more trials than the default.

Definition at line 77 of file mass_fit.h.

The documentation for this class was generated from the following file:

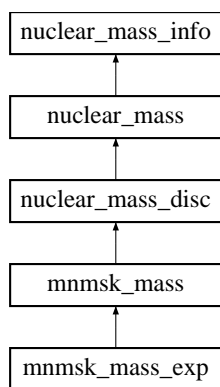
- mass_fit.h

4.17 mnmsk_mass Class Reference

Mass formula from Moller, Nix, Myers and Swiatecki.

```
#include <nuclear_mass.h>
```

Inheritance diagram for [mnmsk_mass](#):



4.17.1 Detailed Description

The data containing an object of type `moller_mass_entry` for 8979 nuclei is automatically loaded by the constructor. If the file (`nucmass/mnmsk.o2`) is not found, then `is_loaded()` will return false and all calls to `get_ZN()` will return an object with $N=Z=0$.

There are several entries in the original **table** which are blank because they are in some way not known, measured, or computable. To distinguish these values from zero, blank entries have been replaced by the number 1.0×10^{99} . For convenience, this value is returned by `blank()`.

Definition at line 874 of file `nuclear_mass.h`.

Public Member Functions

- virtual bool `is_included` (int Z, int N)
*Return false if the mass formula does not include specified **nucleus**.*
- virtual double `mass_excess` (int Z, int N)
Given Z and N, return the mass excess in MeV.
- `mnmsk_mass_entry` `get_ZN` (int l_Z, int l_N)
Get the entry for the specified proton and neutron number.
- bool `is_loaded` ()
*Verify that the constructor properly loaded the **table**.*
- double `blank` ()
The value which corresponds to a blank entry.
- double `neither` ()
Neither beta+ or beta- is possible.
- double `beta_stable` ()
The value which corresponds to a blank entry.
- double `beta_plus_and_minus` ()
Both beta+ and beta- are possible.
- double `greater_100` ()
The value is greater than 100.
- double `very_large` ()
The value is greater than 10^{20} .
- virtual const char * `type` ()
Return the type, "mnmsk_mass".

Data Fields

- int `n`
*The number of **table** entries.*
- `mnmsk_mass_entry` * `mass`
*The array containing the **table**.*

Protected Attributes

- bool `loaded`
*True if the **table** was successfully loaded.*
- int `last`
*The last **table** index for caching.*

4.17.2 Member Function Documentation

4.17.2.1 mnmsk_mass_entry get_ZN (int *I_Z*, int *I_N*)

This method searches the **table** using a cached binary search algorithm. It is assumed that the **table** is sorted first by proton number and then by neutron number.

The documentation for this class was generated from the following file:

- `nuclear_mass.h`

4.18 mnmsk_mass_entry Struct Reference

Mass formula entry structure for Moller, et al.

```
#include <nuclear_mass.h>
```

4.18.1 Detailed Description

Definition at line 740 of file `nuclear_mass.h`.

Data Fields

- int `N`
Neutron number.
- int `Z`
Proton number.
- int `A`
Atomic number.
- double `Emic`
The ground-state microscopic energy.
- double `Mth`
The theoretical mass excess (in MeV).
- double `Mexp`
The experimental mass excess (in MeV).
- double `sigmaexp`
Experimental mass excess error.
- double `EmicFL`
The ground-state microscopic energy in the FRLDM.
- double `MthFL`
The theoretical mass excess in the FRLDM.
- std::string `spinp`
Spin and parity of odd proton.
- std::string `spinn`
Spin and parity of odd neutron.
- double `gapp`
Lipkin-Nogami proton gap.
- double `gapn`

Lipkin-Nogami neutron gap.

- double [be](#)
Total binding energy.
- double [S1n](#)
One neutron separation energy.
- double [S2n](#)
Two neutron separation energy.
- double [PA](#)
Percentage of daughters generated in beta decay after beta-delayed neutron emission.
- double [PAm1](#)
Desc.
- double [PAm2](#)
Desc.
- double [Qbeta](#)
Energy released in beta-decay.
- double [Tbeta](#)
Half-life w.r.t. GT beta-decay.
- double [S1p](#)
One proton separation energy.
- double [S2p](#)
Two proton separation energy.
- double [Qalpha](#)
Energy released in alpha-decay.
- double [Talpha](#)
Half-life w.r.t. alpha-decay.

Ground state deformations (perturbed-spheroid parameterization)

- double [eps2](#)
Quadrupole.
- double [eps3](#)
Octupole.
- double [eps4](#)
Hexadecapole.
- double [eps6](#)
Hexacontatetrapole.
- double [eps6sym](#)
Hexacontatetrapole without mass asymmetry.

Ground state deformations in the spherical-harmonics expansion

- double [beta2](#)
Quadrupole.
- double [beta3](#)
Octupole.
- double [beta4](#)
Hexadecapole.
- double [beta6](#)
Hexacontatetrapole.

The documentation for this struct was generated from the following file:

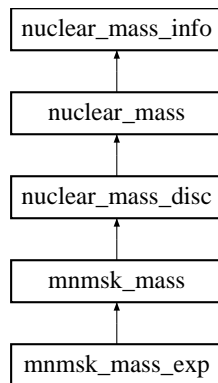
- `nuclear_mass.h`

4.19 mnmsk_mass_exp Class Reference

The experimental values from Moller, Nix, Myers and Swiatecki.

```
#include <nuclear_mass.h>
```

Inheritance diagram for mnmsk_mass_exp::



4.19.1 Detailed Description

Definition at line 953 of file `nuclear_mass.h`.

Public Member Functions

- virtual bool `is_included` (int Z, int N)
Return false if the mass formula does not include specified *nucleus*.
- virtual double `mass_excess` (int Z, int N)
Given Z and N, return the mass excess in MeV.

The documentation for this class was generated from the following file:

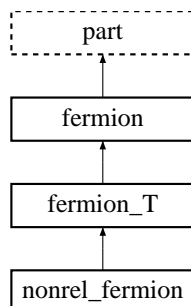
- `nuclear_mass.h`

4.20 nonrel_fermion Class Reference

Nonrelativistic `fermion` class.

```
#include <nonrel_fermion.h>
```

Inheritance diagram for nonrel_fermion::



4.20.1 Detailed Description

The rest mass energy density is given by $n*m$ not $n*ms$. Note that the effective mass here is the Landau mass, not the Dirac mass.

Pressure is computed with

$$P = 2\varepsilon/3$$

and entropy density with

$$s = \frac{5\varepsilon}{3T} - \frac{n\mu}{T}$$

These relations can be verified with an integration by parts. See, e.g. [Callen](#) pg. 403 or [Landau](#) pg. 164.

The functions `fermion::pair_density()` and `pair_mu()` have not been implemented.

Todo

Check behaviour of `calc_density()` at zero density, and compare with that from `eff_fermion`, `rel_fermion`, and `classical`.

Todo

Make sure to test with `non-interacting` equal to true or false, and document whether or not it works with both `inc_rest_mass` equal to true or false

Idea for future

This could be improved by performing a Chebyshev approximation (for example) to invert the density integral so that we don't need to use a solver.

Definition at line 85 of file `nonrel_fermion.h`.

Public Member Functions

- `nonrel_fermion` (double `m`=0.0, double `g`=0.0)
Create a nonrelativistic fermion with mass 'm' and degeneracy 'g'.
- virtual int `calc_mu_zerot` ()
Zero temperature fermions.
- virtual int `calc_density_zerot` ()
Zero temperature fermions.
- virtual int `calc_mu` (const double `temper`)
Calculate properties as function of chemical potential.
- virtual int `calc_density` (const double `temper`)
Calculate properties as function of density.
- virtual int `pair_mu` (const double `temper`)
Calculate properties with antiparticles as function of chemical potential.
- virtual int `pair_density` (const double `temper`)
Calculate properties with antiparticles as function of density.
- virtual int `nu_from_n` (const double `temper`)
Calculate effective chemical potential from density.
- int `set_density_root` (**root**< double, **funct**< double > > &rp)
Set the solver for use in calculating the chemical potential from the density.
- virtual const char * `type` ()
Return string denoting type ("nonrel_fermion").

Data Fields

- **cern_mroot_root**< double, **funct**< double > > `def_density_root`
The default solver for `calc_density()`.

4.20.2 Member Function Documentation

4.20.2.1 virtual int calc_density (const double *temper*) [virtual]

If the density is zero, this function just sets [part::mu](#), [part::nu](#), [part::ed](#), [part::pr](#), and [part::en](#) to zero and returns without calling the error handler (even though at zero density and finite temperature, the chemical potentials formally are equal to $-\infty$).

Implements [fermion_T](#).

The documentation for this class was generated from the following file:

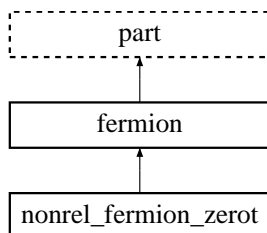
- nonrel_fermion.h

4.21 nonrel_fermion_zerot Class Reference

A zero temperature non-relativistic [fermion](#).

```
#include <nonrel_fermion.h>
```

Inheritance diagram for nonrel_fermion_zerot::



4.21.1 Detailed Description

Definition at line 164 of file nonrel_fermion.h.

Public Member Functions

- [nonrel_fermion_zerot](#) (double *m*=0.0, double *g*=0.0)
Create a nonrelativistic fermion with mass 'm' and degeneracy 'g'.
- virtual int [calc_mu_zerot](#) ()
Zero temperature fermions.
- virtual int [calc_density_zerot](#) ()
Zero temperature fermions.
- virtual const char * [type](#) ()
Return string denoting type ("nonrel_fermion_zerot").

The documentation for this class was generated from the following file:

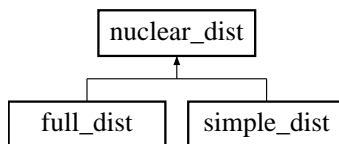
- nonrel_fermion.h

4.22 nuclear_dist Class Reference

A distribution of nuclei [abstract base].

```
#include <nuclear_dist.h>
```

Inheritance diagram for nuclear_dist::



4.22.1 Detailed Description

The virtual base class for a collection of objects of type [nucleus](#) . See [simple_dist](#) and [full_dist](#) for implementations of this base class. Definition at line 41 of file `nuclear_dist.h`.

Data Structures

- class [iterator](#)
An [iterator](#) for the nuclear distribution.

Public Member Functions

- virtual [iterator](#) `begin` ()=0
The beginning of the distribution.
- virtual [iterator](#) `end` ()=0
The end of the distribution.
- virtual `size_t` `size` ()=0
The number of nuclei in the distribution.

The documentation for this class was generated from the following file:

- `nuclear_dist.h`

4.23 nuclear_dist::iterator Class Reference

An [iterator](#) for the nuclear distribution.

```
#include <nuclear_dist.h>
```

4.23.1 Detailed Description

The standard usage of this [iterator](#) is something of the form:

```
mnmsk_mass mth;
simple_dist sd(5,6,10,12,&mth);
for(nuclear_dist::iterator ndi=sd.begin();ndi!=sd.end();ndi++) {
    // do something here for each nucleus
}
```

which would create a list consisting of three isotopes (A=10, 11, and 12) of boron and three isotopes carbon for a total of six nuclei.

Definition at line 71 of file `nuclear_dist.h`.

Public Member Functions

- `iterator (nuclear_dist *ndpp, nucleus *npp)`
Create an *iterator* from the given distribution using the *nucleus* specified in `npp`.
- `iterator operator++ ()`
Proceed to the next *nucleus*.
- `iterator operator++ (int unused)`
Proceed to the next *nucleus*.
- `nucleus * operator → () const`
Dereference the *iterator*.

Protected Attributes

- `nucleus * np`
A pointer to the current *nucleus*.
- `nuclear_dist * ndp`
A pointer to the distribution.

Friends

- `bool operator== (const nuclear_dist::iterator &i1, const nuclear_dist::iterator &i2)`
Give access to the `==` operator.
- `bool operator!= (const nuclear_dist::iterator &i1, const nuclear_dist::iterator &i2)`
Give access to the `!=` operator.

The documentation for this class was generated from the following file:

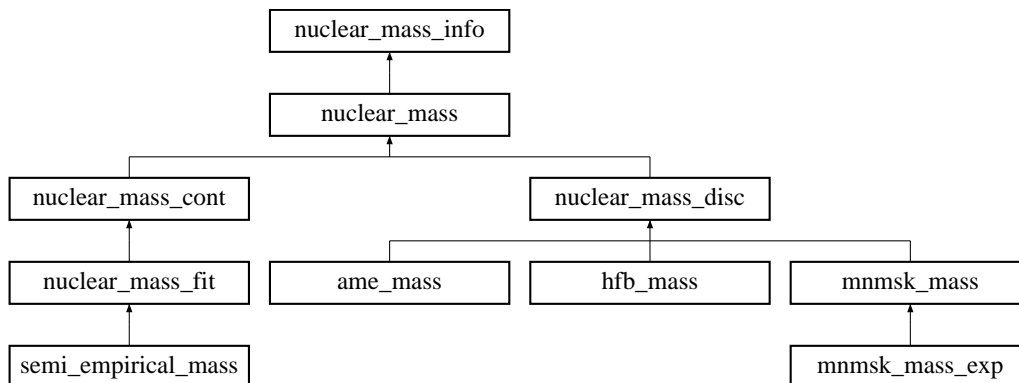
- `nuclear_dist.h`

4.24 nuclear_mass Class Reference

Nuclear mass formula base [abstract base].

```
#include <nuclear_mass.h>
```

Inheritance diagram for `nuclear_mass::`



4.24.1 Detailed Description

This base class provides some default functionality for the nuclear mass formulas. For typical usage, use `ame_mass`, `mnmsk_mass`, `mnmsk_mass_exp`, or `semi_empirical_mass`.

Elements 112-118 are named "Uub", "Uut", "Uuq", "Uup", "Uuh", "Uus", and "Uuo", respectively.

The binding energy is defined by

$$BE = Zm_H + Nm_n - m_{\text{nuclide}}$$

where m_{nuclide} is the mass of the [nucleus](#) including the mass of the electrons. The mass excess is defined by

$$m_{\text{excess}} = m_{\text{nuclide}} - Am_u$$

For example, for U^{238} , the binding energy is 1801.695 MeV, the mass excess is 47.30366 MeV, and m_{nuclide} is 221742.9 MeV. This is consistent with the above, as m_H is 938.7830 MeV, m_n is 939.5650 MeV, and m_u is 931.494 MeV.

Some mass formulas are undefined for sufficiently exotic nuclei. You can use the function [is_included\(\)](#) to find if a particular nucleus is included or not.

Warning:

The treatment of the electron binding energy contribution is not necessarily consistent at present.

Some common reaction Q-values and separation energies:

$Q(\beta^-) = M(A, Z) - M(A, Z + 1)$: Beta-decay energy

$Q(2\beta^-) = M(A, Z) - M(A, Z + 2)$: Double beta-decay energy

$Q(4\beta^-) = M(A, Z) - M(A, Z + 4)$: Four beta-decay energy

$Q(\alpha) = M(A, Z) - M(A - 4, Z - 2) - M(\text{He}^4)$: Alpha-decay energy

$Q(\beta - n) = M(A, Z) - M(A - 1, Z + 1) - M(n)$: Beta-delayed neutron emission decay energy

$Q(d, \alpha) = M(A, Z) - M(A - 2, Z - 1) - M(\text{He}^4) - M(\text{H}^2)$: (d, α) reaction energy

$Q(\text{EC}) = M(A, Z) - M(A, Z - 1)$: Electron capture decay energy

$Q(\text{ECp}) = M(A, Z) - M(A - 1, Z - 2)$: Electron capture with delayed proton emission decay energy

$Q(n, \alpha) = M(A, Z) - M(A - 3, Z - 2) - M(\text{He}^4) + M(n)$: (n, α) reaction energy

$Q(p, \alpha) = M(A, Z) - M(A - 3, Z - 1) - M(\text{He}^4) + M(p)$: (p, α) reaction energy

$S(n) = -M(A, Z) + M(A - 1, Z) + M(n)$: Neutron separation energy

$S(p) = -M(A, Z) + M(A - 1, Z - 1) + M(\text{H}^1)$: Proton separation energy

$S(2n) = -M(A, Z) + M(A - 2, Z) + 2M(n)$: Two neutron separation energy

$S(2p) = -M(A, Z) + M(A - 2, Z - 2) + 2M(\text{H}^1)$: Two proton separation energy

Idea for future

Make the treatment of the electron binding energy contribution more consistent.

Idea for future

It might be useful to consider a fudge factor to ensure no problems with finite precision arithmetic when converting double to int.

Definition at line 290 of file nuclear_mass.h.

Public Member Functions

- virtual const char * [type](#) ()
Return the type, "nuclear_mass".
- virtual bool [is_included](#) (int Z, int N)

Return false if the mass formula does not include specified [nucleus](#).

- `int get_nucleus (int Z, int N, nucleus &n)`
Fill `n` with the information from [nucleus](#) with the given neutron and proton number.
- `virtual double mass_excess (int Z, int N)=0`
Given `Z` and `N`, return the mass excess in MeV.
- `virtual double mass_excess_d (double Z, double N)=0`
Given `Z` and `N`, return the mass excess in MeV.
- `virtual double binding_energy (int Z, int N)`
Return the binding energy in MeV.
- `virtual double binding_energy_d (double Z, double N)`
Return the binding energy in MeV.
- `virtual double total_mass (int Z, int N)`
Return the total mass of the [nucleus](#) (without the electrons) in MeV.
- `virtual double total_mass_d (double Z, double N)`
Return the total mass of the [nucleus](#) (without the electrons) in MeV.

4.24.2 Member Function Documentation

4.24.2.1 virtual double [binding_energy](#) (int Z, int N) [inline, virtual]

The binding energy is defined to be negative for bound nuclei, thus the binding energy per baryon of Pb-208 is about $-8 \times 208 = -1664$ MeV.

Definition at line 351 of file `nuclear_mass.h`.

4.24.2.2 virtual double [binding_energy_d](#) (double Z, double N) [inline, virtual]

The binding energy is defined to be negative for bound nuclei, thus the binding energy per baryon of Pb-208 is about $-8 \times 208 = -1664$ MeV.

Definition at line 365 of file `nuclear_mass.h`.

4.24.2.3 int [get_nucleus](#) (int Z, int N, [nucleus](#) &n) [inline]

All masses are given in fm^{-1} . The total mass (withouth the electrons) is put in `part::m` and `part::ms`, the binding energy is placed in `nucleus::be`, the mass excess in `nucleus::mex` and the degeneracy (`part::g`) is arbitrarily set to 1 for even A nuclei and 2 for odd A nuclei.

Warning:

The spin degeneracy is not handled particularly intelligently. This function simply assumes 0 spin for even A and spin 1/2 for odd A.

Definition at line 325 of file `nuclear_mass.h`.

4.24.2.4 virtual bool [is_included](#) (int Z, int N) [inline, virtual]

Note:

By default, this returns false, so that children must overload this function if they do not provide masses for arbitrary nuclei.

Reimplemented in `ame_mass`, `mnmsk_mass`, `mnmsk_mass_exp`, and `hfb_mass`.

Definition at line 308 of file `nuclear_mass.h`.

The documentation for this class was generated from the following file:

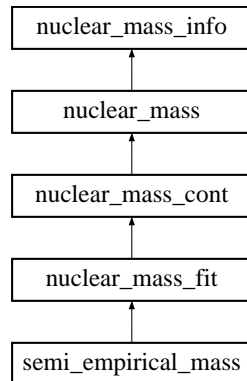
- `nuclear_mass.h`

4.25 nuclear_mass_cont Class Reference

Continuous nuclear mass formula [abstract base].

```
#include <nuclear_mass.h>
```

Inheritance diagram for nuclear_mass_cont::



4.25.1 Detailed Description

Definition at line 421 of file nuclear_mass.h.

Public Member Functions

- virtual double [mass_excess](#) (int Z, int N)
Given Z and N, return the mass excess in MeV.
- virtual double [mass_excess_d](#) (double Z, double N)=0
Given Z and N, return the mass excess in MeV.

The documentation for this class was generated from the following file:

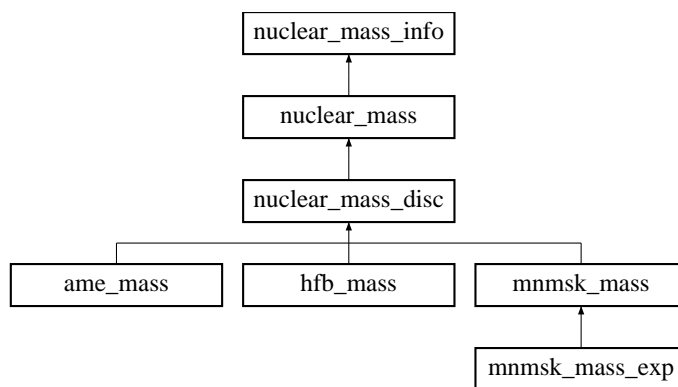
- nuclear_mass.h

4.26 nuclear_mass_disc Class Reference

Discrete nuclear mass formula [abstract base].

```
#include <nuclear_mass.h>
```

Inheritance diagram for nuclear_mass_disc::



4.26.1 Detailed Description

This uses simple linear interpolation to obtain masses of nuclei with non-integer Z and N which may be particularly sensitive to the form of the pairing.

Definition at line 398 of file nuclear_mass.h.

Public Member Functions

- virtual double [mass_excess](#) (int Z , int N)=0
Given Z and N , return the mass excess in MeV.
- virtual double [mass_excess_d](#) (double Z , double N)
Given Z and N , return the mass excess in MeV.

The documentation for this class was generated from the following file:

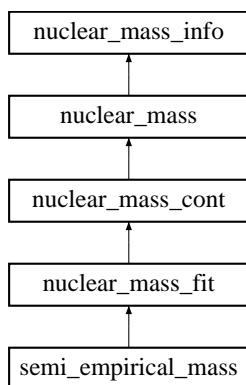
- nuclear_mass.h

4.27 nuclear_mass_fit Class Reference

Fittable mass formula.

```
#include <nuclear_mass.h>
```

Inheritance diagram for nuclear_mass_fit::



4.27.1 Detailed Description

Nuclear mass formulas which are descendants of this class can be fit to experiment using [mass_fit](#).

Definition at line 441 of file nuclear_mass.h.

Public Member Functions

- virtual const char * [type](#) ()
Return the type, "nuclear_mass_fit".
- virtual int [fit_fun](#) (size_t nv, const **ovector_base** &x)
Fix parameters from an array for fitting.
- virtual int [guess_fun](#) (size_t nv, **ovector_base** &x)
Fill array with guess from present values for fitting.

Data Fields

- size_t [nfit](#)
Number of fitting parameters.

The documentation for this class was generated from the following file:

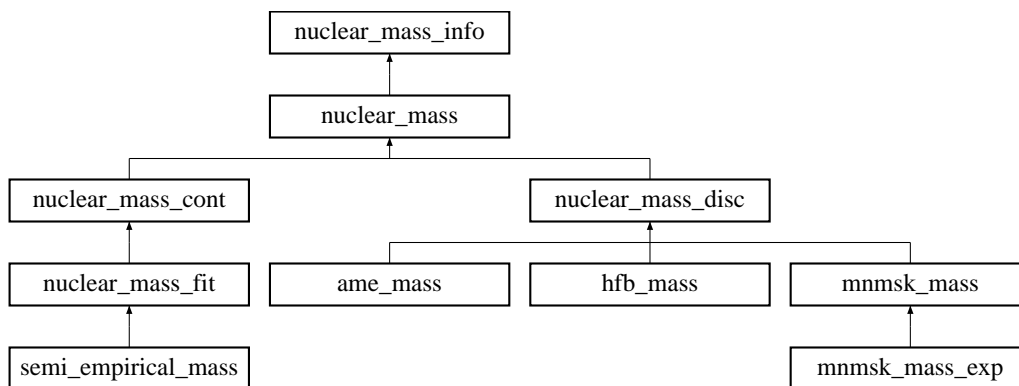
- nuclear_mass.h

4.28 nuclear_mass_info Class Reference

Nuclear mass info.

```
#include <nuclear_mass.h>
```

Inheritance diagram for nuclear_mass_info::



4.28.1 Detailed Description

Definition at line 39 of file nuclear_mass.h.

Data Structures

- struct [string_less_than](#)
String comparison operator for element_table.

Public Member Functions

- int `parse_elstring` (std::string *ela*, int &Z, int &N, int &A)
Parse a string representing an element.
- int `eltoZ` (std::string *el*)
Return Z given the element name.
- std::string `Ztoel` (size_t Z)
Return the element name given Z.
- std::string `tostring` (size_t Z, size_t N)
Return a string of the form "Pb208" for a given Z and N.

Protected Types

- typedef std::map< std::string, int, `string_less_than` >::iterator `table_it`
A convenient typedef for an iterator for element_table.

Protected Attributes

- std::map< std::string, int, `string_less_than` > `element_table`
A map containing the proton numbers organized by element name.
- std::string `element_list` [`nelements`]
The list of elements organized by proton number.

Static Protected Attributes

- static const int `nelements` = 119
The number of elements (proton number).

4.28.2 Member Function Documentation

4.28.2.1 int eltoZ (std::string *el*) [inline]

If the string parameter *el* is invalid, the error handler is called and the value -1 is returned.

Definition at line 143 of file nuclear_mass.h.

4.28.2.2 int parse_elstring (std::string *ela*, int &Z, int &N, int &A) [inline]

Accepts strings of one of the following forms:

- Pb208
- pb208
- Pb 208
- Pb-208
- pb 208
- pb-208 or one of the special strings n, p, d or t for the neutron, proton, deuteron, and triton, respectively.

Note:

At present, this allows nuclei which don't make sense because $A < Z$, such as Carbon-5.

Idea for future

Allow A to precede Z.

Idea for future

Right now, n4 is interpreted incorrectly as Nitrogen-4, rather than the tetraneutron.

Definition at line 66 of file nuclear_mass.h.

4.28.2.3 std::string toString (size_t Z, size_t N) [inline]

Note that if Z is zero, then and 'n' is used to indicate the a [nucleus](#) composed entirely of neutrons and if the argument Z is greater than 118, an empty string is returned (independ.

Definition at line 176 of file nuclear_mass.h.

4.28.2.4 std::string Ztoel (size_t Z) [inline]**Note:**

This function returns "n" indicating the neutron for Z=0, and if the argument Z is greater than 118, an empty string is returned after calling the error handler.

Definition at line 160 of file nuclear_mass.h.

The documentation for this class was generated from the following file:

- nuclear_mass.h

4.29 nuclear_mass_info::string_less_than Struct Reference

String comparison operator for element_table.

```
#include <nuclear_mass.h>
```

4.29.1 Detailed Description

Definition at line 189 of file nuclear_mass.h.

Public Member Functions

- bool **operator()** (const std::string s1, const std::string s2) const

The documentation for this struct was generated from the following file:

- nuclear_mass.h

4.30 nuclear_reaction Class Reference

A simple nuclear reaction specification.

```
#include <reaction_lib.h>
```

4.30.1 Detailed Description

Definition at line 41 of file reaction_lib.h.

Public Member Functions

- `std::string to_string ()`
Convert the reaction to a string for screen output.
- `int clear ()`
Clear the rate.
- `nuclear_reaction (const nuclear_reaction &nr)`
Copy constructor.
- `nuclear_reaction & operator= (const nuclear_reaction &nr)`
Copy constructor.
- `double rate (double T9)`
Compute the reaction rate from the temperature in units of $10^9 K$.

Data Fields

- `size_t chap`
Chapter.
- `std::string name [6]`
Names of the participating nuclei.
- `std::string ref`
Reference.
- `char type`
Type of rate (resonant/non-resonant/weak).
- `char rev`
Forward or reverse.
- `double Q`
Q value.
- `double a [7]`
Coefficients.
- `size_t Z [6]`
Proton number of participating nuclei.
- `size_t A [6]`
Mass number of participating nuclei.
- `size_t isomer [6]`
Isomer designation of participating nuclei.

The documentation for this class was generated from the following file:

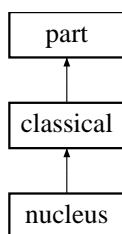
- reaction_lib.h

4.31 nucleus Class Reference

A simple `nucleus` class.

```
#include <nucleus.h>
```

Inheritance diagram for `nucleus::`



4.31.1 Detailed Description

The variable `part::m` is typically used for the mass of the `nucleus` with no electrons.

The binding energy of the `nucleus` (`be`) is typically defined as the mass of the `nucleus` (without the electrons) minus Z times the mass of the proton minus N times the mass of the neutron.

The mass excess (`be`) is defined as the mass of the `nucleus` including the electron contribution minus a times the mass of the atomic mass unit.

The variable `part::inc_rest_mass` is set to `false` by default, to insure that energies and chemical potentials do not include the rest mass. This is typically appropriate for nuclei.

Definition at line 51 of file `nucleus.h`.

Data Fields

- int `Z`
Proton number.
- int `N`
Neutron number.
- int `A`
Atomic number.
- double `mex`
Mass excess in fm^{-1} .
- double `be`
Binding energy in fm^{-1} (with a minus sign for bound nuclei).

The documentation for this class was generated from the following file:

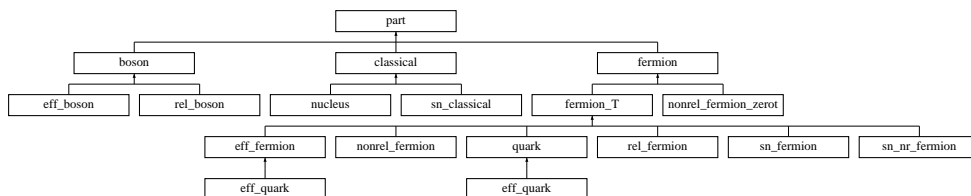
- `nucleus.h`

4.32 part Class Reference

Particle base class.

```
#include <part.h>
```

Inheritance diagram for `part::`



4.32.1 Detailed Description

Calculate the properties of particles from their chemical potential (`calc_mu()` and `pair_mu()`) or from the density (`calc_density()` and `pair_density()`).

When `non-interacting` is false, the thermodynamic integrals need both a value of "mu" and "nu". "nu" is an effective chemical potential which appears in the argument of the exponential of the Fermi-function.

Keep in mind, that the pair functions use `anti()`, which assumes that $\nu \rightarrow -\nu$ and $\mu \rightarrow -\mu$ for the anti-particles, which might not be true for interacting particles. When `non-interacting` is true, then "ms" is set equal to "m", and "nu" is set equal to "mu", everywhere.

The "density" functions use the value of nu (or mu when `non_interacting` is true) for an initial guess. Zero is very likely a bad guess, but these functions will not warn you about this.

Definition at line 94 of file `part.h`.

Public Member Functions

- `part` (double `m`=0.0, double `g`=0.0)
make a particle of mass `m` and degeneracy `g`.
- virtual int `init` (double `m`, double `g`)
Set the mass `m` and degeneracy `g`.
- virtual int `anti` (`part` &`ax`)
Make an anti-particle.
- virtual const char * `type` ()
Return string denoting type ("part").

Data Fields

- double `g`
degeneracy
- double `m`
mass
- double `n`
density
- double `ed`
energy density
- double `pr`
pressure
- double `mu`
chemical potential
- double `en`
entropy
- double `ms`
effective mass (Dirac unless otherwise specified)
- double `nu`
effective chemical potential
- bool `inc_rest_mass`
derivative of energy with respect to effective mass
- bool `non_interacting`
True if the particle is non-interacting (default true).

The documentation for this class was generated from the following file:

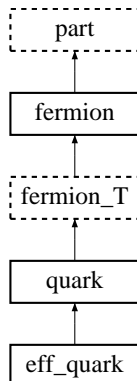
- `part.h`

4.33 quark Class Reference

Quark base [abstract base].

```
#include <quark.h>
```

Inheritance diagram for quark::



4.33.1 Detailed Description

Definition at line 43 of file quark.h.

Public Member Functions

- `quark` (double mass=0.0, double dof=0.0)
Create a `quark` with mass `m` and degeneracy `g`.
- virtual int `calc_mu` (const double temper)=0
Calculate properties as function of chemical potential.
- virtual int `calc_density` (const double temper)=0
Calculate properties as function of density.
- virtual int `pair_mu` (const double temper)=0
Calculate properties with antiparticles as function of chemical potential.
- virtual int `pair_density` (const double temper)=0
Calculate properties with antiparticles as function of density.
- virtual const char * `type` ()
Return string denoting type ("quark").

Data Fields

- double `B`
Contribution to the bag constant.
- double `qq`
Quark condensate.

The documentation for this class was generated from the following file:

- quark.h

4.34 reaction_lib Class Reference

Simple reaction library.

```
#include <reaction_lib.h>
```

4.34.1 Detailed Description

Units:

- Chapters 1,2,3, and 11: 1/s
- Chapters 4,5,6, and 7: $\text{cm}^3/\text{g}/\text{s}$
- Chapter 8 and 9: $\text{cm}^6/\text{g}^2/\text{s}$
- Chapter 10: $\text{cm}^9/\text{g}^3/\text{s}$

Chapters:

- 1: $\text{nuc1} \rightarrow \text{nuc2}$
- 2: $\text{nuc1} \rightarrow \text{nuc2} + \text{nuc3}$
- 3: $\text{nuc1} \rightarrow \text{nuc2} + \text{nuc3} + \text{nuc4}$
- 4: $\text{nuc1} + \text{nuc2} \rightarrow \text{nuc3}$
- 5: $\text{nuc1} + \text{nuc2} \rightarrow \text{nuc3} + \text{nuc4}$
- 6: $\text{nuc1} + \text{nuc2} \rightarrow \text{nuc3} + \text{nuc4} + \text{nuc5}$
- 7: $\text{nuc1} + \text{nuc2} \rightarrow \text{nuc3} + \text{nuc4} + \text{nuc5} + \text{nuc6}$
- 8: $\text{nuc1} + \text{nuc2} + \text{nuc3} \rightarrow \text{nuc4}$
- 9: $\text{nuc1} + \text{nuc2} + \text{nuc3} \rightarrow \text{nuc4} + \text{nuc5}$
- 10: $\text{nuc1} + \text{nuc2} + \text{nuc3} + \text{nuc4} \rightarrow \text{nuc5} + \text{nuc6}$
- 11: $\text{nuc1} \rightarrow \text{nuc2} + \text{nuc3} + \text{nuc4} + \text{nuc5}$

Original FORTRAN format:

```
FORMAT (i1, 4x, 6a5, 8x, a4, a1, a1, 3x, 1p e12.5)
FORMAT (4e13.6)
FORMAT (3e13.6)
```

Definition at line 202 of file reaction_lib.h.

Public Member Functions

- int [read_file_reaclib2](#) (std::string fname)
Read from a file in the REACLIB2 format.
- int [find_in_chap](#) (std::vector< [nuclear_reaction](#) > &nrl, size_t chap, std::string nuc1, std::string nuc2="", std::string nuc3="", std::string nuc4="", std::string nuc5="", std::string nuc6="")
Find a set of nuclear reactions in a specified chapter.

Data Fields

- `std::vector< nuclear_reaction > lib`
The library.

Protected Member Functions

- `bool matches (size_t ul, size_t ri)`
Test if entry ul in the arrays matches the library reaction.

Protected Attributes

- `int fN [6]`
Storage for the find function.
- `int fZ [6]`
- `int fA [6]`
- `size_t fi`

4.34.2 Member Function Documentation

4.34.2.1 `int read_file_reacli2 (std::string fname)`

Note:

This function does not check that the chapter numbers are correct for the subsequent reaction.

The documentation for this class was generated from the following file:

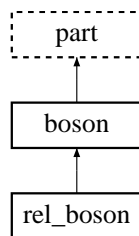
- `reaction_lib.h`

4.35 rel_boson Class Reference

Equation of state for a relativistic [boson](#).

```
#include <rel_boson.h>
```

Inheritance diagram for `rel_boson`:



4.35.1 Detailed Description

[Todo](#)

Testing not completely finished.

Definition at line 48 of file `rel_boson.h`.

Public Member Functions

- `rel_boson` (double `m`=0.0, double `g`=0.0)
Create a `boson` with mass `m` and degeneracy `g`.
- virtual int `calc_mu` (const double `temper`)
Calculate properties as function of chemical potential.
- virtual int `calc_density` (const double `temper`)
Calculate properties as function of density.
- virtual int `pair_mu` (const double `temper`)
Calculate properties with antiparticles as function of chemical potential.
- virtual int `pair_density` (const double `temper`)
Calculate properties with antiparticles as function of density.
- virtual int `nu_from_n` (const double `temper`)
Calculate effective chemical potential from density.
- int `set_inte` (**inte**< const double, **funct**< const double > > &l_nit, **inte**< const double, **funct**< const double > > &l_dit)
*Set **inte** object.*
- int `set_density_root` (**root**< const double, **funct**< const double > > &rp)
Set the solver for use in calculating the chemical potential from the density.
- virtual const char * `type` ()
Return string denoting type ("rel_boson").

Data Fields

- int `mroot_err`
The error value from `mroot`.
- int `inte_err`
The error value from `inte`.
- **cern_mroot_root**< const double, **funct**< const double > > `def_density_root`
The default solver for `calc_density()`.
- **gsl_inte_qagiu**< const double, **funct**< const double > > `def_nit`
Default nondegenerate integrator.
- **gsl_inte_qag**< const double, **funct**< const double > > `def_dit`
Default degenerate integrator.

The documentation for this class was generated from the following file:

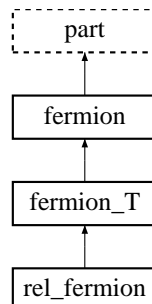
- `rel_boson.h`

4.36 rel_fermion Class Reference

Equation of state for a relativistic `fermion`.

```
#include <rel_fermion.h>
```

Inheritance diagram for `rel_fermion`:



4.36.1 Detailed Description

Method:

This implements an equation of state for a relativistic [fermion](#) using direct integration. Define the degeneracy parameter

$$\psi = (\nu - m^*)/T$$

where ν is the effective chemical potential and m^* is the effective mass. For ψ greater than [deg_limit](#) (degenerate regime), a finite interval integrator is used and for ψ less than [deg_limit](#) (non-degenerate regime), an integrator over the interval from $[0, \infty)$ is used. The upper limit on the degenerate integration is given by

$$\sqrt{(uT + \nu)^2 - m^{*,2}}$$

where u is [rel_fermion::upper_limit_fac](#).

In the degenerate regime, the entropy integrand

$$-k^2 [n \log n + (1 - n) \log (1 - n)]$$

can lose precision when $(E - \nu)/T$ is negative and sufficiently large in absolute magnitude. Thus when $(E - \nu)/T < S$ where S is stored in [deg_entropy_fac](#) (default is -30), the integrand is written as

$$-k^2 (E/T - \nu/T) e^{E/T - \nu/T}.$$

If $(E - \nu)/T < S$ is less than -1 times [exp_limit](#) (e.g. less than -200), then the entropy integrand is assumed to be zero.

The integrands in the non-degenerate regime are written in a dimensionless form, by defining $u \equiv p/T$, $y \equiv \nu/T$, and $mx \equiv m^*/T$. The density integrand is

$$(mx + u) \sqrt{u^2 + 2mxu} \left(\frac{e^y}{e^{mx+u} + e^y} \right),$$

the energy integrand is

$$(mx + u)^2 \sqrt{u^2 + 2mxu} \left(\frac{e^y}{e^{mx+u} + e^y} \right),$$

and the entropy integrand is

$$(mx + u) \sqrt{u^2 + 2mxu} (t_1 + t_2),$$

where

$$\begin{aligned} t_1 &= \log(1 + e^{y-mx-u}) / (1 + e^{y-mx-u}) \\ t_2 &= \log(1 + e^{mx+u-y}) / (1 + e^{mx+u-y}). \end{aligned}$$

The default integrators are [gsl_inte_qag](#) (for degenerate particles) and [gsl_inte_qagiu](#) (for non-degenerate particles).

Accuracy:

The default settings for for this class give an accuracy of at least 1 [part](#) in 10^6 (and frequently better than this).

When the integrators provide numerical uncertainties, these uncertainties are stored in [unc](#). In the case of [calc_density\(\)](#) and [pair_density\(\)](#), the uncertainty from the numerical accuracy of the solver is not included. (There is also a relatively small inaccuracy due to the mathematical evaluation of the integrands which is not included in [unc](#).)

One way to improve the accuracy of the computation is just to decrease the tolerances on the default integration objects. This can be done, using, for example

```
rel_fermion rf(1.0, 2.0);
rf.def_dit.tolx/=1.0e2;
rf.def_dit.tolf/=1.0e2;
rf.def_nit.tolx/=1.0e2;
rf.def_nit.tolf/=1.0e2;
```

which decreases the both the relative and absolute tolerances for both the degenerate and non-degenerate integrators. If one is using either the [calc_density\(\)](#) or [pair_density\(\)](#) functions, one may also have to improve the accuracy of the solver which determines the chemical potential from the density. For the default solver, this could be done with

```
rf.def_density_root.tolx/=1.0e2;
rf.def_density_root.tolf/=1.0e2;
```

Of course if these tolerances are too small, the calculation may fail.

Todos:

Note:

This does not work with `inc_rest_mass=false` (3/30/09: There is a significant amount of testing code to ensure that it does work with `inc_rest_mass=false`, so this may have been fixed already.)

Idea for future

Improve the asymptotics of the non-degenerate integrands.

Idea for future

It appears this doesn't compute the uncertainty in the chemical potential or density with [calc_density\(\)](#). This could be fixed.

Definition at line 164 of file `rel_fermion.h`.

Public Member Functions

- [rel_fermion](#) (double `m`=0.0, double `g`=0.0)
Create a [fermion](#) with mass `m` and degeneracy `g`.
- virtual int [calc_mu](#) (double `temper`)
Calculate properties as function of chemical potential.
- virtual int [calc_density](#) (double `temper`)
Calculate properties as function of density.
- virtual int [pair_mu](#) (double `temper`)
Calculate properties with antiparticles as function of chemical potential.
- virtual int [pair_density](#) (double `temper`)
Calculate properties with antiparticles as function of density.
- virtual int [nu_from_n](#) (double `temper`)
Calculate effective chemical potential from density.
- int [set_inte](#) (**inte**< double, **funct**< double > > &non_it, **inte**< double, **funct**< double > > °_it)
Set integrators.
- int [set_density_root](#) (**root**< double, **funct**< double > > &rp)
Set the solver for use in calculating the chemical potential from the density.
- virtual const char * [type](#) ()
Return string denoting type ("rel_fermion").

Data Fields

- [fermion_unc](#)
Storage for the uncertainty.
 - bool [guess_from_nu](#)
If true, use the present value of the chemical potential as a guess for the new chemical potential.
 - **cern_mroot_root**< double, **funct**< double > > [def_density_root](#)
The default solver for [calc_density\(\)](#).
 - **gsl_inte_qag**< double, **funct**< double > > [def_dit](#)
-

The default integrator for degenerate fermions.

- **gsl_inte_qagiu**< double, **funct**< double > > **def_nit**

The default integrator for non-degenerate fermions.

Numerical parameters

- double **deg_limit**
The critical degeneracy at which to switch integration techniques (default 2).
- double **exp_limit**
The limit for exponentials to ensure integrals are finite (default 200).
- double **upper_limit_fac**
The factor for the degenerate upper limits (default 20).
- double **deg_entropy_fac**
A factor for the degenerate entropy integration (default 30).

Protected Member Functions

- double **density_fun** (double u, double &pa)
The integrand for the density for non-degenerate fermions.
- double **energy_fun** (double u, double &pa)
The integrand for the energy density for non-degenerate fermions.
- double **entropy_fun** (double u, double &pa)
The integrand for the entropy density for non-degenerate fermions.
- double **deg_density_fun** (double u, double &pa)
The integrand for the density for degenerate fermions.
- double **deg_energy_fun** (double u, double &pa)
The integrand for the energy density for degenerate fermions.
- double **deg_entropy_fun** (double u, double &pa)
The integrand for the entropy density for degenerate fermions.
- int **solve_fun** (double x, double &yy, double &pa)
Solve for the chemical potential given the density.
- int **pair_fun** (double x, double &yy, double &pa)
Solve for the chemical potential given the density with antiparticles.

Protected Attributes

- **inte**< double, **funct**< double > > * **nit**
The non-degenerate integrator.
- **inte**< double, **funct**< double > > * **dit**
The degenerate integrator.
- **root**< double, **funct**< double > > * **density_root**
The solver for [calc_density\(\)](#).

The documentation for this class was generated from the following file:

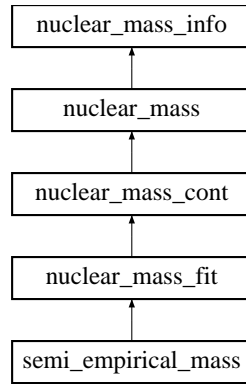
- rel_fermion.h

4.37 semi_empirical_mass Class Reference

Semi-empirical mass formula.

```
#include <nuclear_mass.h>
```

Inheritance diagram for semi_empirical_mass::



4.37.1 Detailed Description

A simple semi-empirical mass formula of the form

$$E/A = B + S_s \frac{1}{A^{1/3}} + E_c \frac{Z^2}{A^{4/3}} + S_v \left(1 - \frac{2Z}{A}\right)^2 + E_{\text{pair}}(Z, N)$$

where

$$E_{\text{pair}}(Z, N) = \frac{E_{\text{pair}}}{A^{3/2}} \times \begin{cases} -1 & \text{N and Z even} \\ +1 & \text{N and Z odd} \\ 0 & \text{otherwise} \end{cases}$$

Note:

The default parameters are arbitrary, and are not determined from a fit.

There is an example of the usage of this class given in [Nuclear mass fit example](#).

Definition at line 491 of file nuclear_mass.h.

Public Member Functions

- virtual const char * [type](#) ()
Return the type, "semi_empirical_mass".
- virtual double [mass_excess_d](#) (double Z, double N)
Given Z and N, return the mass excess in MeV.
- virtual int [fit_fun](#) (size_t nv, const [ovector_base](#) &x)
Fix parameters from an array for fitting.
- virtual int [guess_fun](#) (size_t nv, [ovector_base](#) &x)
Fill array with guess from present values for fitting.

Data Fields

- double [B](#)
Binding energy (negative and in MeV, default -16).
- double [Sv](#)
Symmetry energy (in MeV, default 23.7).
- double [Ss](#)
Surface energy (in MeV, default 18).
- double [Ec](#)
Coulomb energy (in MeV, default 0.7).
- double [Epair](#)

Pairing energy (MeV, default 13.0).

The documentation for this class was generated from the following file:

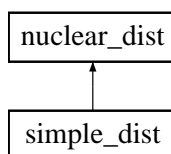
- nuclear_mass.h

4.38 simple_dist Class Reference

A simple nuclear distribution given a range in A and Z.

```
#include <nuclear_dist.h>
```

Inheritance diagram for simple_dist::



4.38.1 Detailed Description

The iterator for this distribution begins with the [nucleus](#) with the lowest Z and A, and increases A before incrementing Z and beginning again with the lowest A for that value of Z. In other words, it proceeds through all the isotopes of an element first, and then proceeds to the next element.

For example, to create a collection of isotopes of Carbon, Nitrogen and Oxygen using the most recent (2003) Atomic Mass Evaluation, and then output the nuclei in the collection

```
ame_mass ame;
simple_dist fd(6,8,2,30,&ame);
for(nuclear_dist::iterator ndi=fd.begin();ndi!=fd.end();ndi++) {
    cout << ndi->Z << " " << ndi->A << " " << ndi->m << endl;
}
```

Idea for future

Make the vector constructor into a template so it accepts any type. Do the same for [set_dist\(\)](#).

Definition at line 163 of file nuclear_dist.h.

Public Member Functions

- [simple_dist](#) ()
Create an empty distribution.
- [simple_dist](#) (int minZ, int maxZ, int minA[], int maxA[], [nuclear_mass](#) &nm)
Create a distribution from ranges in A specified for each Z.
- [simple_dist](#) (int minZ, int maxZ, int minA, int maxA, [nuclear_mass](#) &nm)
Create a square distribution in A and Z.
- virtual [iterator begin](#) ()
The beginning of the distribution.
- virtual [iterator end](#) ()
The end of the distribution.
- virtual [size_t size](#) ()
The number of nuclei in the distribution.

- `int set_dist (int minZ, int maxZ, int minA[], int maxA[], nuclear_mass &nm)`
Set the distribution from ranges in A specified for each Z.
- `int set_dist (int minZ, int maxZ, int minA, int maxA, nuclear_mass &nm)`
Set a square distribution in A and Z.

4.38.2 Constructor & Destructor Documentation

4.38.2.1 `simple_dist (int minZ, int maxZ, int minA[], int maxA[], nuclear_mass & nm)`

The length of the arrays minA and maxA should be exactly $\text{maxZ} - \text{minZ} + 1$.

4.38.3 Member Function Documentation

4.38.3.1 `int set_dist (int minZ, int maxZ, int minA[], int maxA[], nuclear_mass & nm)`

The length of the arrays minA and maxA should be exactly $\text{maxZ} - \text{minZ} + 1$.

The documentation for this class was generated from the following file:

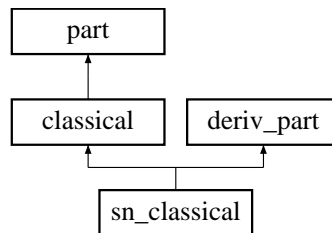
- `nuclear_dist.h`

4.39 `sn_classical` Class Reference

Equation of state for a [classical](#) particle with derivatives.

```
#include <sn_classical.h>
```

Inheritance diagram for `sn_classical`:



4.39.1 Detailed Description

[Todo](#)

This does not work with `inc_rest_mass=true`

Definition at line 42 of file `sn_classical.h`.

Public Member Functions

- `sn_classical (double m=0.0, double g=0.0)`
Create a [fermion](#) with mass `m` and degeneracy `g`.
- virtual `int calc_mu (const double temper)`
Calculate properties as function of chemical potential.
- virtual `int calc_density (const double temper)`
Calculate properties as function of density.
- virtual `const char * type ()`
Return string denoting type ("`sn_classical`").

The documentation for this class was generated from the following file:

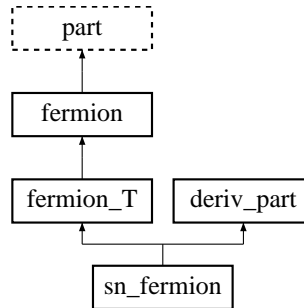
- sn_classical.h

4.40 sn_fermion Class Reference

Equation of state for a relativistic [fermion](#).

```
#include <sn_fermion.h>
```

Inheritance diagram for sn_fermion::



4.40.1 Detailed Description

Note:

This class does not work with inc_rest_mass=true.

This implements an equation of state for a relativistic [fermion](#) using direct integration. After subtracting the rest mass from the chemical potentials, the distribution function is

$$\left\{ 1 + \exp[(\sqrt{k^2 + m^{*2}} - m - \nu)/T] \right\}^{-1}$$

where k is the momentum, ν is the effective chemical potential, m is the rest mass, and m^* is the effective mass. For later use, we define $E^* = \sqrt{k^2 + m^{*2}}$. The degeneracy parameter is

$$\psi = (\nu + (m - m^*))/T$$

For ψ greater than [deg_limit](#) (degenerate regime), a finite interval integrator is used and for ψ less than [deg_limit](#) (non-degenerate regime), an integrator over the interval from $[0, \infty)$ is used. The upper limit on the degenerate integration is given by the solution of

$$(\sqrt{k^2 + m^{*,2}} - m - \nu)/T = \text{flimit}$$

which is

$$\sqrt{(m + \mathcal{L})^2 - m^{*2}}$$

where $\mathcal{L} \equiv \text{flimit} \times T + \nu$.

In the non-degenerate regime, we make the substitution $u = k/T$ to ensure that the variable of integration scales properly.

Uncertainties are given in [unc](#).

Todo

This needs to be corrected to calculate $\sqrt{k^2 + m^{*2}} - m$ gracefully when $m^* \approx m \ll k$.

Todo

Call error handler if inc_rest_mass is true or update to properly treat the case when inc_rest_mass is true.

Evaluation of the derivatives

The relevant derivatives of the distribution function are

$$\begin{aligned}\frac{\partial f}{\partial T} &= f(1-f) \frac{E^* - m - \nu}{T^2} \\ \frac{\partial f}{\partial \nu} &= f(1-f) \frac{1}{T} \\ \frac{\partial f}{\partial k} &= -f(1-f) \frac{k}{E^* T} \\ \frac{\partial f}{\partial m^*} &= -f(1-f) \frac{m^*}{E^* T}\end{aligned}$$

We also need the derivative of the entropy integrand w.r.t. the distribution function, which is

$$\mathcal{S} \equiv f \ln f + (1-f) \ln(1-f) \quad \frac{\partial \mathcal{S}}{\partial f} = \ln \left(\frac{f}{1-f} \right) = \left(\frac{\nu - E^* + m}{T} \right)$$

where the entropy density is

$$s = -\frac{g}{2\pi^2} \int_0^\infty \mathcal{S} k^2 dk$$

The derivatives can be integrated directly (`method = direct`) or they may be converted to integrals over the distribution function through an integration by parts (`method = byparts`)

$$\int_a^b f(k) \frac{dg(k)}{dk} dk = f(k)g(k) \Big|_{k=a}^{k=b} - \int_a^b g(k) \frac{df(k)}{dk} dk$$

using the distribution function for $f(k)$ and 0 and ∞ as the limits, we have

$$\frac{g}{2\pi^2} \int_0^\infty \frac{dg(k)}{dk} f dk = \frac{g}{2\pi^2} \int_0^\infty g(k) f(1-f) \frac{k}{E^* T} dk$$

as long as $g(k)$ vanishes at $k = 0$. Rewriting,

$$\frac{g}{2\pi^2} \int_0^\infty h(k) f(1-f) dk = \frac{g}{2\pi^2} \int_0^\infty f \frac{T}{k} \left[h' E^* - \frac{h E^*}{k} + \frac{h k}{E^*} \right] dk$$

as long as $h(k)/k$ vanishes at $k = 0$.

Explicit forms

1) The derivative of the density wrt the chemical potential

$$\left(\frac{dn}{d\mu} \right)_T = \frac{g}{2\pi^2} \int_0^\infty \frac{k^2}{T} f(1-f) dk$$

Using $h(k) = k^2/T$ we get

$$\left(\frac{dn}{d\mu} \right)_T = \frac{g}{2\pi^2} \int_0^\infty \left(\frac{k^2 + E^{*2}}{E^*} \right) f dk$$

2) The derivative of the density wrt the temperature

$$\left(\frac{dn}{dT} \right)_\mu = \frac{g}{2\pi^2} \int_0^\infty \frac{k^2 (E^* - m - \nu)}{T^2} f(1-f) dk$$

Using $h(k) = k^2(E^* - \nu)/T^2$ we get

$$\left(\frac{dn}{dT}\right)_\mu = \frac{g}{2\pi^2} \int_0^\infty \frac{f}{T} \left[2k^2 + E^{*2} - E^* (\nu + m) - k^2 \left(\frac{\nu + m}{E^*} \right) \right] dk$$

3) The derivative of the entropy wrt the chemical potential

$$\left(\frac{ds}{d\mu}\right)_T = \frac{g}{2\pi^2} \int_0^\infty k^2 f(1-f) \frac{(E^* - m - \nu)}{T^2} dk$$

This verifies the Maxwell relation

$$\left(\frac{ds}{d\mu}\right)_T = \left(\frac{dn}{dT}\right)_\mu$$

4) The derivative of the entropy wrt the temperature

$$\left(\frac{ds}{dT}\right)_\mu = \frac{g}{2\pi^2} \int_0^\infty k^2 f(1-f) \frac{(E^* - m - \nu)^2}{T^3} dk$$

Using $h(k) = k^2(E^* - \nu)^2/T^3$

$$\left(\frac{ds}{dT}\right)_\mu = \frac{g}{2\pi^2} \int_0^\infty \frac{f(E^* - m - \nu)}{E^* T^2} [E^{*3} + 3E^* k^2 - (E^{*2} + k^2)(\nu + m)] dk$$

5) The derivative of the density wrt the effective mass

$$\left(\frac{dn}{dm^*}\right)_{T,\mu} = -\frac{g}{2\pi^2} \int_0^\infty \frac{k^2 m^*}{E^* T} f(1-f) dk$$

Using $h(k) = -(k^2 m^*)/(E^* T)$ we get

$$\left(\frac{dn}{dm^*}\right)_{T,\mu} = -\frac{g}{2\pi^2} \int_0^\infty m^* f dk$$

Note:

The dsdT integration may fail if the system is very degenerate. When method is byparts, the integral involves a large cancellation between the regions from $k \in (0, \text{ulimit}/2)$ and $k \in (\text{ulimit}/2, \text{ulimit})$. Switching to method=direct and setting the lower limit to llimit, may help, but recent testing on this gave negative values for dsdT. For very degenerate systems, an expansion may be better than trying to perform the integration.

Idea for future

It might be worth coding up direct differentiation, or differentiating the eff results, as these may succeed more generally.

Idea for future

This class will have difficulty with extremely degenerate or extremely non-degenerate systems. Fix this.

Idea for future

Create a more intelligent method for dealing with bad initial guesses for the chemical potential in `calc_density()`.

Definition at line 227 of file sn_fermion.h.

Method of computing derivatives

- int `method`
Method (default is `byparts`).
- static const int `direct` = 1
In the form containing $f(1-f)$.
- static const int `byparts` = 2
Integrate by parts.

Public Member Functions

- `sn_fermion` (double `m`=0.0, double `g`=0.0)
Create a `fermion` with mass `m` and degeneracy `g`.
- virtual int `calc_mu` (double `temper`)
Calculate properties as function of chemical potential.
- virtual int `calc_density` (double `temper`)
Calculate properties as function of density.
- virtual int `pair_mu` (double `temper`)
Calculate properties with antiparticles as function of chemical potential.
- virtual int `pair_density` (double `temper`)
Calculate properties with antiparticles as function of density.
- virtual int `nu_from_n` (double `temper`)
Calculate effective chemical potential from density.
- int `set_inte` (`inte`< double, `funct`< double > > &`unit`, `inte`< double, `funct`< double > > &`udit`)
Set `inte` objects.
- int `set_density_root` (`root`< double, `funct`< double > > &`rp`)
Set the solver for use in calculating the chemical potential from the density.
- virtual const char * `type` ()
Return string denoting type ("sn_fermion").

Data Fields

- double `deg_limit`
The critical degeneracy at which to switch integration techniques (default 2.0).
- double `flimit`
The limit for the Fermi functions (default 20.0).
- `fermion unc`
Storage for the most recently calculated uncertainties.
- `deriv_part dunc`
Storage for the most recently calculated uncertainties.
- `gsl_inte_qagiu`< double, `funct`< double > > `def_nit`
The default integrator for the non-degenerate regime.
- `gsl_inte_qag`< double, `funct`< double > > `def_dit`
The default integrator for the degenerate regime.
- `cern_mroot_root`< double, `funct`< double > > `def_density_root`
The default solver for `npen_density()` and `pair_density()`.

4.40.2 Member Function Documentation

4.40.2.1 int set_inte (inte< double, funct< double > > & unit, inte< double, funct< double > > & udit)

The first integrator is used for non-degenerate integration and should integrate from 0 to ∞ (like `gsl_inte_qagiu`). The second integrator is for the degenerate case, and should integrate between two finite values.

4.40.3 Field Documentation

4.40.3.1 double flimit

`sn_fermion` will ignore corrections smaller than about $\exp(-\text{flimit})$.

Definition at line 246 of file `sn_fermion.h`.

The documentation for this class was generated from the following file:

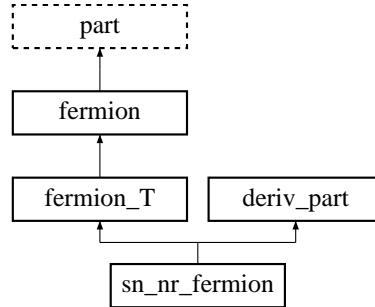
- `sn_fermion.h`

4.41 sn_nr_fermion Class Reference

Equation of state for a nonrelativistic [fermion](#).

```
#include <sn_nr_fermion.h>
```

Inheritance diagram for sn_nr_fermion::



4.41.1 Detailed Description

This does not include the rest mass energy in the chemical potential or the rest mass energy density in the energy density to alleviate numerical precision problems at low densities

This implements an equation of state for a nonrelativistic [fermion](#) using direct integration. After subtracting the rest mass from the chemical potentials, the distribution function is

$$\left\{ 1 + \exp \left[\left(\frac{k^2}{2m^*} - \nu \right) / T \right] \right\}^{-1}$$

where ν is the effective chemical potential, m is the rest mass, and m^* is the effective mass. For later use, we define $E^* = k^2/2/m^*$.

Uncertainties are given in [unc](#).

Evaluation of the derivatives

The relevant derivatives of the distribution function are

$$\frac{\partial f}{\partial T} = f(1-f) \frac{E^* - \nu}{T^2}$$

$$\frac{\partial f}{\partial \nu} = f(1-f) \frac{1}{T}$$

$$\frac{\partial f}{\partial k} = -f(1-f) \frac{k}{m^* T}$$

$$\frac{\partial f}{\partial m^*} = f(1-f) \frac{k^2}{2m^{*2} T}$$

We also need the derivative of the entropy integrand w.r.t. the distribution function, which is quite simple

$$\mathcal{S} \equiv f \ln f + (1-f) \ln(1-f) \quad \frac{\partial \mathcal{S}}{\partial f} = \ln \left(\frac{f}{1-f} \right) = \left(\frac{\nu - E^*}{T} \right)$$

where the entropy density is

$$s = -\frac{g}{2\pi^2} \int_0^\infty \mathcal{S} k^2 dk$$

The derivatives can be integrated directly or they may be converted to integrals over the distribution function through an integration by parts

$$\int_a^b f(k) \frac{dg(k)}{dk} dk = f(k)g(k)|_{k=a}^{k=b} - \int_a^b g(k) \frac{df(k)}{dk} dk$$

using the distribution function for $f(k)$ and 0 and ∞ as the limits, we have

$$\frac{g}{2\pi^2} \int_0^\infty \frac{dg(k)}{dk} f dk = \frac{g}{2\pi^2} \int_0^\infty g(k) f(1-f) \frac{k}{E^* T} dk$$

as long as $g(k)$ vanishes at $k = 0$. Rewriting,

$$\frac{g}{2\pi^2} \int_0^\infty h(k) f(1-f) dk = \frac{g}{2\pi^2} \int_0^\infty f \frac{T m^*}{k} \left[h' - \frac{h}{k} \right] dk$$

as long as $h(k)/k$ vanishes at $k = 0$.

Explicit forms

1) The derivative of the density wrt the chemical potential

$$\left(\frac{dn}{d\mu} \right)_T = \frac{g}{2\pi^2} \int_0^\infty \frac{k^2}{T} f(1-f) dk$$

Using $h(k) = k^2/T$ we get

$$\left(\frac{dn}{d\mu} \right)_T = \frac{g}{2\pi^2} \int_0^\infty m^* f dk$$

2) The derivative of the density wrt the temperature

$$\left(\frac{dn}{dT} \right)_\mu = \frac{g}{2\pi^2} \int_0^\infty \frac{k^2 (E^* - \nu)}{T^2} f(1-f) dk$$

Using $h(k) = k^2 (E^* - \nu)/T^2$ we get

$$\left(\frac{dn}{dT} \right)_\mu = \frac{g}{2\pi^2} \int_0^\infty \frac{f}{T} [m^* (E^* - \nu) - k^2] dk$$

3) The derivative of the entropy wrt the chemical potential

$$\left(\frac{ds}{d\mu} \right)_T = \frac{g}{2\pi^2} \int_0^\infty k^2 f(1-f) \frac{(E^* - \nu)}{T^2} dk$$

This verifies the Maxwell relation

$$\left(\frac{ds}{d\mu} \right)_T = \left(\frac{dn}{dT} \right)_\mu$$

4) The derivative of the entropy wrt the temperature

$$\left(\frac{ds}{dT} \right)_\mu = \frac{g}{2\pi^2} \int_0^\infty k^2 f(1-f) \frac{(E^* - \nu)^2}{T^3} dk$$

Using $h(k) = k^2 (E^* - \nu)^2/T^3$

$$\left(\frac{ds}{dT} \right)_\mu = \frac{g}{2\pi^2} \int_0^\infty f \frac{m^*}{T^2} \left[(E^* - \nu)^2 + \frac{2k^2}{m^*} (E^* - \nu) \right] dk$$

5) The derivative of the density wrt the effective mass

$$\left(\frac{dn}{dm^*} \right)_{T,\mu} = \frac{g}{2\pi^2} \int_0^\infty \frac{k^2}{2m^{*2}T} f(1-f) k^2 dk$$

Using $h(k) = k^4/(2m^{*2}T)$ we get

$$\left(\frac{dn}{dm^{*}}\right)_{T,\mu} = \frac{g}{2\pi^2} \int_0^\infty f \frac{3k^2}{2m^{*}} dk$$

New section

$u = k^2/2/m^{*}/T$ and $y = \mu/T$, so

$$kdk = m^{*}T du$$

or

$$dk = \frac{m^{*}T}{\sqrt{2m^{*}T}u} du = \sqrt{\frac{m^{*}T}{2u}} du$$

1) The derivative of the density wrt the chemical potential

$$\left(\frac{dn}{d\mu}\right)_T = \frac{gm^{*3/2}\sqrt{T}}{2^{3/2}\pi^2} \int_0^\infty u^{-1/2} f du$$

2) The derivative of the density wrt the temperature

$$\left(\frac{dn}{dT}\right)_\mu = \frac{gm^{*3/2}\sqrt{T}}{2^{3/2}\pi^2} \int_0^\infty f du \left[3u^{1/2} - yu^{-1/2}\right]$$

4) The derivative of the entropy wrt the temperature

$$\left(\frac{ds}{dT}\right)_\mu = \frac{gm^{*3/2}T^{1/2}}{2^{3/2}\pi^2} \int_0^\infty f \left[5u^{3/2} - 6yu^{1/2} + y^2u^{-1/2}\right] du$$

5) The derivative of the density wrt the effective mass

$$\left(\frac{dn}{dm^{*}}\right)_{T,\mu} = \frac{3gm^{*1/2}T^{3/2}}{2^{3/2}\pi^2} \int_0^\infty u^{1/2} f du$$

Definition at line 221 of file sn_nr_fermion.h.

Public Member Functions

- [sn_nr_fermion](#) (double **m**=0.0, double **g**=0.0)
Create a *fermion* with mass **m** and degeneracy **g**.
- virtual int [calc_mu](#) (const double temper)
Calculate properties as function of chemical potential.
- virtual int [calc_density](#) (const double temper)
Calculate properties as function of density.
- virtual int [pair_mu](#) (const double temper)
Calculate properties with antiparticles as function of chemical potential.
- virtual int [pair_density](#) (const double temper)
Calculate properties with antiparticles as function of density.
- virtual int [nu_from_n](#) (const double temper)
Calculate effective chemical potential from density.
- int [set_density_root](#) (**root**< const double, **funct**< const double > > &rp)
Set the solver for use in calculating the chemical potential from the density.
- virtual const char * [type](#) ()
Return string denoting type ("sn_nr_fermion").

Data Fields

- double `flimit`
The limit for the Fermi functions (default 20.0).
- `fermion unc`
Storage for the most recently calculated uncertainties.
- `deriv_part dunc`
Storage for the most recently calculated uncertainties.
- bool `guess_from_nu`
If true, use the present value of the chemical potential as a guess for the new chemical potential.
- `cern_mroot_root` < const double, `funct` < const double > > `def_density_root`
The default solver for `npen_density()` and `pair_density()`.

Protected Member Functions

- int `solve_fun` (double x, double &yy, const double &temper)
Function to compute chemical potential from density.
- int `pair_fun` (double x, double &yy, const double &temper)
Function to compute chemical potential from density when antiparticles are included.

Protected Attributes

- `root` < const double, `funct` < const double > > * `density_root`
Solver to compute chemical potential from density.

4.41.2 Field Documentation

4.41.2.1 double flimit

`sn_nr_fermion` will ignore corrections smaller than about $\exp(-\text{flimit})$.

Definition at line 235 of file `sn_nr_fermion.h`.

The documentation for this class was generated from the following file:

- `sn_nr_fermion.h`

4.42 thermo Class Reference

A class for the thermodynamical variables (energy density, pressure, entropy density).

```
#include <part.h>
```

4.42.1 Detailed Description

Definition at line 46 of file `part.h`.

Public Member Functions

- const char * `type` ()
Return string denoting type ("thermo").

Data Fields

- double [pr](#)
pressure
- double [ed](#)
energy density
- double [en](#)
entropy density

The documentation for this class was generated from the following file:

- [part.h](#)

5 File Documentation

5.1 [part.h](#) File Reference

File for definitions for [thermo](#) and [part](#).

```
#include <string>
#include <iostream>
#include <cmath>
#include <o2scl/constants.h>
#include <o2scl/inte.h>
#include <o2scl/collection.h>
#include <o2scl/funct.h>
#include <o2scl/mroot.h>
```

5.1.1 Detailed Description

Definition in file [part.h](#).

Data Structures

- class [thermo](#)
A class for the thermodynamical variables (energy density, pressure, entropy density).
- class [part](#)
Particle base class.

Functions

- [thermo operator+](#) (const [thermo](#) &left, const [thermo](#) &right)
Addition operator.
- [thermo operator-](#) (const [thermo](#) &left, const [thermo](#) &right)
Subtraction operator.
- [thermo operator+](#) (const [thermo](#) &left, const [part](#) &right)
Addition operator.
- [thermo operator-](#) (const [thermo](#) &left, const [part](#) &right)
Subtraction operator.

Index

ame_entry, 7
ame_entry03_io_type, 8
ame_entry95_io_type, 8
ame_mass, 9

binding_energy
 nuclear_mass, 40
binding_energy_d
 nuclear_mass, 40
boson, 11
 co, 12
 massless_calc_mu, 12

calc_density
 eff_fermion, 21
 eff_quark, 22
 nonrel_fermion, 36
calc_density_zerot
 fermion, 24
calc_mu
 eff_fermion, 21
 eff_quark, 22
calc_mu_zerot
 fermion, 24
classical, 12
co
 boson, 12

def_massless_root
 fermion_T, 26
def_mmin
 mass_fit, 30
deriv_part, 13

eff_boson, 16
 load_coefficients, 18
eff_fermion, 18
 calc_density, 21
 calc_mu, 21
 load_coefficients, 21
 pair_mu, 21
eff_quark, 21
 calc_density, 22
 calc_mu, 22
 pair_mu, 22
eltoZ
 nuclear_mass_info, 44
energy_density_zerot
 fermion, 24

fermion, 23
 calc_density_zerot, 24
 calc_mu_zerot, 24
 energy_density_zerot, 24

 kf_from_density, 24
 pressure_zerot, 24
fermion_T, 24
 def_massless_root, 26
 massless_pair_density, 26
flimit
 sn_fermion, 62
 sn_nr_fermion, 66
full_dist, 27
 set_dist, 27

get_nucleus
 nuclear_mass, 40
get_ZN
 hfb_mass, 29
 mnmsk_mass, 32

hfb_mass, 27
 get_ZN, 29
 hfb_mass, 28
 hfb_mass, 28
hfb_mass_entry, 29

is_included
 nuclear_mass, 40

kf_from_density
 fermion, 24

load_coefficients
 eff_boson, 18
 eff_fermion, 21

mass_fit, 29
 def_mmin, 30
massless_calc_mu
 boson, 12
massless_pair_density
 fermion_T, 26
mnmsk_mass, 30
 get_ZN, 32
mnmsk_mass_entry, 32
mnmsk_mass_exp, 34

nonrel_fermion, 34
 calc_density, 36
nonrel_fermion_zerot, 36
nuclear_dist, 36
nuclear_dist::iterator, 37
nuclear_mass, 38
 binding_energy, 40
 binding_energy_d, 40
 get_nucleus, 40
 is_included, 40

- nuclear_mass_cont, [41](#)
- nuclear_mass_disc, [41](#)
- nuclear_mass_fit, [42](#)
- nuclear_mass_info, [43](#)
 - eltoZ, [44](#)
 - parse_elstring, [44](#)
 - tostring, [45](#)
 - Ztoel, [45](#)
- nuclear_mass_info::string_less_than, [45](#)
- nuclear_reaction, [45](#)
- nucleus, [46](#)

- pair_mu
 - eff_fermion, [21](#)
 - eff_quark, [22](#)
- parse_elstring
 - nuclear_mass_info, [44](#)
- part, [47](#)
- part.h, [67](#)
- pressure_zerot
 - fermion, [24](#)

- quark, [49](#)

- reaction_lib, [50](#)
 - read_file_reaclib2, [51](#)
- read_file_reaclib2
 - reaction_lib, [51](#)
- rel_boson, [51](#)
- rel_fermion, [52](#)

- semi_empirical_mass, [55](#)
- set_dist
 - full_dist, [27](#)
 - simple_dist, [58](#)
- set_inte
 - sn_fermion, [62](#)
- simple_dist, [57](#)
 - set_dist, [58](#)
 - simple_dist, [58](#)
 - simple_dist, [58](#)
- sn_classical, [58](#)
- sn_fermion, [59](#)
 - flimit, [62](#)
 - set_inte, [62](#)
- sn_nr_fermion, [63](#)
 - flimit, [66](#)

- thermo, [66](#)
- tostring
 - nuclear_mass_info, [45](#)

- Ztoel
 - nuclear_mass_info, [45](#)
